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Corrugated shells: an algorithm for generating double-curvature geometric surfaces for structural analysis

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Abstract

Analysis of corrugated shell structures is an interesting problem in Structural Mechanics, which has many practical applications in Civil Engineering and Architecture. Thanks to corrugation, these structures have a remarkable feature: the wavy (undulated) shape in their edge provides significant enhancements in their structural behaviour, increasing the bending stiffness at the edge and allowing for a non-negligible reduction of its thickness. Moreover, looking at the non-linear behaviour, domes corrugation plays a relevant role in instability phenomena, such as the influence of imperfections and increasing resistance to snap-through.

A problem in the study of such kind of shells is the definition of mathematical and geometrical model and the construction of a suitable mesh to perform FE analyses. The aim of this paper is to find an automated way to generate a double-curvature geometric surface that can be used both in

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static and in non-linear stability analyses of such corrugated shell structures. A method to generate a NURBS surface, suitable for a parametric FE analysis from a geometrical model expressed in a parametric form, is proposed and applied to a shell inspired by the well-known dome designed by Pier Luigi Nervi in 1959 for the roof of the Palasport Flaminio in Rome.

Keywords: corrugated shells, shallow shells, domes, Palasport Flaminio, Pier Luigi Nervi

1. Introduction

The problem presented in this work concerns the structural analysis of corrugated shell structures. These structures have a remarkable characteristic: the wavy shape of their edge gives a significant improvement to their structural behaviour, increasing the bending stiffness at the edge, thus allowing the designer to reduce its thickness. A problem in studying this structural typology is how to deal with its complex geometry. This paper aims to find an automated way to generate a double-curvature geometric surface, given its mathematical description, which can be used both in static and non-linear stability analyses of corrugated shell structures.

For properly treating this topic, it is useful to recall a state-of-the-art where the problem studied in the present work can be framed. Some relevant examples in the field of corrugated shell structures can be found along the history of Civil Engineering. In East Anglia (England) from XVII until XIX century wavy fences were largely built as garden walls, known as *Crinckle crankle walls* (see Fig. 1). Some of them still exist in Suffolk and Hampshire. This kind of construction, which has been attributed to Dutch

18 engineers [1], presents a wavy shape that provides the wall with bending
 19 stiffness and improves its structural response to horizontal loads. As a con-
 20 sequence, bricklayers could build a slender wall made of a single line of bricks
 21 without the need for abutments or buttresses. The effect of corrugation, then,
 22 is to improve the mechanical behaviour of this structure. The same remark
 23 will also be true for shells, as it will be clear in the following. Specifically,
 24 one can observe corrugation in seashell structures arising as a result of opti-
 25 mization processes.

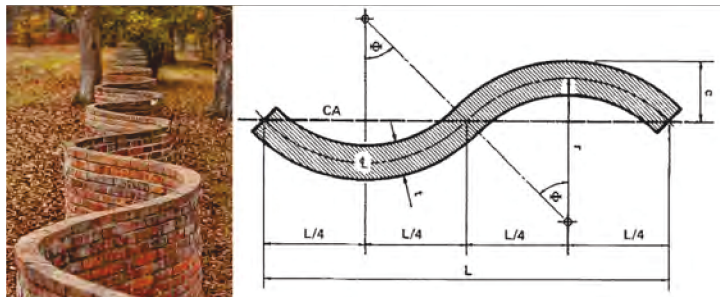


Figure 1: A Crinkle Crinkle wall in Suffolk (UK), left. A sketch of the same wall, right.

26 Indeed, corrugated shells can be found in nature: structures of such type can
 27 be obtained by a topologic optimization process, as those occurring in bone
 28 reconstruction (see, for example, [2–4]). A relevant example, also of interest
 29 for the topic of the present work, is represented by seashells. It is observed
 30 that the mussel minimizes the effort in building its dwelling changing from a
 31 smooth to a corrugated shape [5]. The outcome of the new smart shape is to
 32 increase the mechanical resistance with the same amount of material. As a
 33 consequence, the optimized structure is said to be shape-resistant. The study
 34 presented herein is motivated by similar considerations related to efficiency

35 and resistance criteria.

36 On the Civil Engineering side, the need for corrugated shells or plates
37 is also motivated by structural efficiency. The main differences between
38 the Civil Engineering case and other fields (such as the cited example of
39 seashells) consist in scale and, clearly, in the employed materials. Generally
40 in Civil Engineering and Architecture a standard material is reinforced con-
41 crete (shortly, RC), whose high-performance dissipation properties are also
42 known, as it was pointed out in [6–8]. In addition, special attention should
43 be given to the durability of this material, [9–11].

44 RC has always been an excellent material to be used in optimized [11],
45 [12] and customized-shaped building [13], also in the Italian school of Struc-
46 tural Engineers, which between the '50s and '70s was led by Pier Luigi Nervi
47 and Sergio Musmeci. An outstanding piece of Italian architecture, where a
48 corrugated shell made of reinforced concrete is used, is the roof of a gasoline
49 station in Sesto San Giovanni (Milano), designed by Aldo Favini in 1949,
50 which was unfortunately destroyed some years later. Then this type of con-
51 struction has been progressively fallen into disuse, due to the increase in the
52 cost of the formwork and scaffolding.

53 As it has been already remarked, some models can be suitably adapted to
54 apparently different structures. Referring to Fig. 2, one can observe different
55 objects where corrugation has a relevant role: potentially the spirit of the
56 present work is to develop an algorithm which is useful for all these cases,
57 independently on the scale or material. An extensive body of literature exists
58 and deals with the connection between form and structure, and these topics
59 are covered in foremost books like [14, 15]. A new gaze is provided from

60 the SIXXI project, whose purpose is to give a distinct point of view on the
61 Italian school of Structural Engineering, and is set out in [16]. Even though
62 primary source can be found in Nervi's book [17], some recent advancements
63 have been provided in [18, 19]. The inspiration for this work has been taken
64 from one of Nervi's works: the shallow shell designed for Rome Olympics
65 game in 1960 to cover the roof of the Palazzetto dello Sport or, shortly,
66 *Palasport Flaminio* (from the name of the district of Rome where it was
67 built and still stands nowadays).



Figure 2: Example of corrugated shells: (a) Nervi's Palasport Flaminio dome, (b) Hobermann deployable structure, (c) Favini's roof, (d) corrugated sea-shells.

68 Nervi's shell is a foremost piece of unique architecture and it also constitutes
69 an inexhaustible source for structural design, even in different fields. For
70 instance, it has been a source of influence for the Iris Dome retractable roof,
71 which was designed by Charles Hobermann [20]. This kind of corrugated
72 shapes, known as *umbrella-type surfaces*, can be studied from a mathematical

73 point of view as it has been done, in a more general framework, in [21].

74 One can employ a representation of 2-D surfaces in Cartesian coordinates
75 depending on a set of parameters defined in a closed interval. A comprehen-
76 sive guide for a wide variety of parametric equations can be found in [22].

77 From a Structural Mechanics point of view, the consequences of corruga-
78 tions in building structures have not yet been entirely investigated, perhaps
79 due to the intrinsic difficulties to manage the mathematical implications of
80 corrugation. Some theoretical background is given for static analyses in [23]
81 and considerations about the stability and multi-stability of open corrugated
82 shell are pointed out in [24]. In [25] it is possible to find a parametric analy-
83 sis devoted to understanding the role of corrugation in improving the seismic
84 resistance of vaults and domes.

85 In [26–29] some relevant results in the field of shells were set out. These
86 studies can be useful also for generalizing the results presented in this work [30–
87 34].

88 Frequently, in Architecture it is needed to design large-span roofs: to this
89 aim, the theory of shells provides the most effective approach, introducing
90 structural problems that need to be properly taken into account. A remark-
91 able problem consists in enhancing the structural resistance. This can be
92 made in different ways, as, for instance, by increasing the thickness of the
93 shell surface or by placing a ring-beam on the shell edge. In this context
94 a smarter solution (also from an architectural and aesthetics point of view)
95 consists in employing corrugated shell surfaces which allow in reducing the
96 shell thickness.

97 The above recalled literature is necessary to address the main aim of this

98 paper: to recognise the influence of corrugations in the mechanics of shells,
99 taking into account relevant non-linear effects affecting slender and shal-
100 low shells, whose edge is wavy-corrugated. Non-linear behaviour remarkably
101 affects the shell mechanical performances, such as snap-through mechanism
102 and buckling instability phenomena [35–42]. A successfully employed method
103 in dealing with this kind of problem consists in using a set of safety factors
104 to knock down the theoretical results: see for example the NASA aeronautics
105 recommendations [43].

106 The first step for performing a correct numerical analysis is to set a pro-
107 cedure which can produce geometrical objects replicating the mathematical
108 dome shape in a process suitable for structural analysis using doubly-curved
109 elements. In Section 2 the geometrical representation of a corrugated shell is
110 introduced in such a way that mathematical parametric equations are given.
111 In Section 3 an algorithm to represent a NURBS based surface is presented.
112 Numerical results are shown and discussed in Section 4 in order to investi-
113 gate the influence of shell corrugation. Finally, in Section 5, some concluding
114 remarks are presented.

115 **2. Wavy-edge shell parametric description**

116 In this section, a mathematical description of a wavy-edge surface inspired
117 to the Nervi’s Palasport Flaminio dome is proposed. Its equations depend on
118 several parameters which control the corrugation shape along the shell side.

119 The adopted spherical polar reference system is shown in Fig. 3, where r
120 is the radial distance from the pole, ϑ is the colatitude angle (the complement
121 to the latitude angle) and φ is the longitude angle. So, a generic point P ,

122 belonging to the 3-D space, is uniquely identified by its spherical coordinates
 123 (r, ϑ, φ) . A parametric representation of a wavy-edge spherical shell can be
 124 given introducing a parametrization of its radius. A surface could thus be
 125 described by using two parameters only, viz. ϑ and φ , where each pair $(\vartheta_i,$
 126 $\varphi_i)$ describes a point P_i which belongs to the surface.

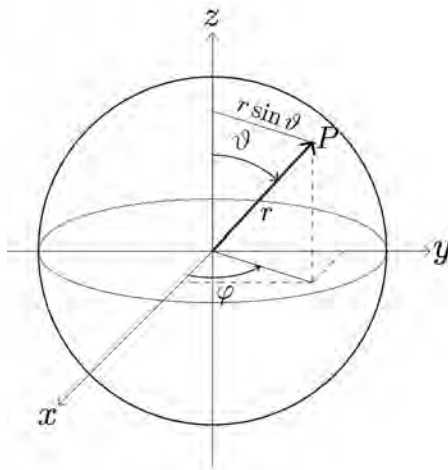


Figure 3: Spherical polar coordinate system.

127 Considering a perfect spherical shell, whose radius is R_0 , its *parametric equa-*
 128 *tions* are the classical ones

$$\left\{ \begin{array}{l} x = R_0 \sin \vartheta \cos \varphi \\ y = R_0 \sin \vartheta \sin \varphi \\ z = R_0 \cos \vartheta. \end{array} \right. \quad (1)$$

130 From Eq. (1), by squaring and summing up (term by term) both sides, pa-
 131 rameters ϑ and φ can be eliminated and the resulting *implicit* representation

132 of the spherical surface is obtained:

$$133 \quad x^2 + y^2 + z^2 - R_0^2 = 0. \quad (2)$$

134 Now, recalling that the radial distance r from the pole is given, in terms of
135 Cartesian coordinates, by:

$$r = \sqrt{x^2 + y^2 + z^2},$$

136 an *explicit* representation of the spherical surface results:

$$137 \quad r = R_0. \quad (3)$$

138 Looking at Eq. (3) it is apparent that in the case of a sphere the radial
139 distance of any point of the surface is independent of the spherical coordinates
140 ϑ, φ . This consideration suggests an easy way to construct a surface, shaped
141 as a portion of a hemispherical shell but exhibiting a corrugation on the
142 edge. Indeed such corrugated edge can be treated as a perturbation of the
143 constant radius R_0 , resulting in a wavy-edge. Then, for such surface the
144 radius $r = r(\vartheta, \varphi)$ may be represented as

$$145 \quad r = R_0 [1 + f(\vartheta)g(\varphi)]. \quad (4)$$

146 In Eq. (4) the perturbation is made up by two factors: the former $f(\vartheta)$,
147 depends only on colatitude ϑ and gives the shape of the perturbed meridian,
148 while the latter $g(\varphi)$ depends only on the longitude angle φ and modulates the
149 form of all parallels. In order to get a cyclic symmetry along each parallel line,
150 function $g(\varphi)$ must be periodic; a suitable choice to get a smooth repetition
151 by a whole number n of a basic wave pattern is then:

$$152 \quad g(\varphi) = \cos(n\varphi). \quad (5)$$

153 This ensures that an undulated wave is repeated n times along the surface
 154 edge, i.e. the period of function g is simply $2\pi/n$; in order to obtain that
 155 the fundamental (or *zero*) meridian $\varphi = 0$ is indeed perturbed with reference
 156 to the spherical shape, the cosine function has been preferred to its sine
 157 counterpart.

158 Function $f(\vartheta)$, which controls the perturbation of the radius along the merid-
 159 ian with reference to that of a perfect sphere, R_0 , can be chosen in several
 160 ways. A possible choice is:

$$161 \quad f(\vartheta) = aH(\vartheta - \vartheta_0) \left(\frac{\vartheta - \vartheta_0}{\vartheta_0} \right)^2. \quad (6)$$

162 In Eq. (6) a is a parameter controlling the amplitude of the perturbation, H
 163 is Heaviside step function (or unit step function), defined as:

$$H(\vartheta - \vartheta_0) = \begin{cases} 1, & \text{if } \vartheta \geq \vartheta_0 \\ 0, & \text{if } \vartheta < \vartheta_0, \end{cases}$$

164 whose role is to switch on the radius perturbation in correspondence of ϑ_0 ,
 165 namely the colatitude angle at which such perturbation originates. Finally
 166 the term $(\vartheta - \vartheta_0)^2/\vartheta_0^2$ has been introduced to produce a smooth variation
 167 of r along the meridian in a neighborhood of ϑ_0 . Despite the presence of
 168 Heaviside's step function, it comes out from Eq. (6) that the resulting radius
 169 $r(\vartheta, \varphi)$ is an almost everywhere continuous and differentiable function of its
 170 arguments.

171 If a smoother shape is desired, the unit step function H can be replaced
 172 by a continuously differentiable function approximating it, like, for instance,

173 the hyperbolic tangent; consequently, in this case, $f(\vartheta)$ can be expressed by:

$$174 \quad f(\vartheta) = \frac{a}{2} [1 + \tanh(b(\vartheta - \vartheta_0))], \quad (7)$$

175 where a is again a parameter controlling the amplitude of the perturbation,
 176 while b is a second parameter which, when increases, makes steeper the graph
 177 of the function and allows approximating, with the desired accuracy, a step
 178 function with a continuous one. Besides, one should notice that the value of
 179 the derivative at $\vartheta = \vartheta_0$ differs for the two proposed parametrizations: for
 180 Eq. (6) such value is zero, whereas for Eq. (7) this is not the case.

181 A mathematical representation of the corrugated surface is then given
 182 by updating the previously mentioned equations of a hemispherical shell,
 183 using $r(\vartheta, \varphi)$ defined by Eq. (4) instead of the constant radius R_0 . As a
 184 consequence, the parametric equations of the corrugated surface become:

$$185 \quad \begin{cases} x = r(\vartheta, \varphi) \sin \vartheta \cos \varphi \\ y = r(\vartheta, \varphi) \sin \vartheta \sin \varphi \\ z = r(\vartheta, \varphi) \cos \vartheta. \end{cases} \quad (8)$$

186 The difference between the two possible choices which were presented above is
 187 shown in Fig. 4. For the case described by Eq. (6), the following parameters
 188 have been adopted: $\vartheta_0 = \pi/6$, $a = \vartheta_0^2$; for that represented by Eq. (7)
 189 $\vartheta_0 = \pi/6$, $a = 1/50$, $b = 50$. In both cases $g(\varphi)$ has been defined as in
 190 Eq. (5) where a value $n = 36$ has been assumed; for comparison purposes
 191 the opening of the dome has been fixed in both cases to the value $\vartheta_f = \pi/5$.
 192 A magnified portion of the corrugated edge is shown for both cases in Fig. 5.

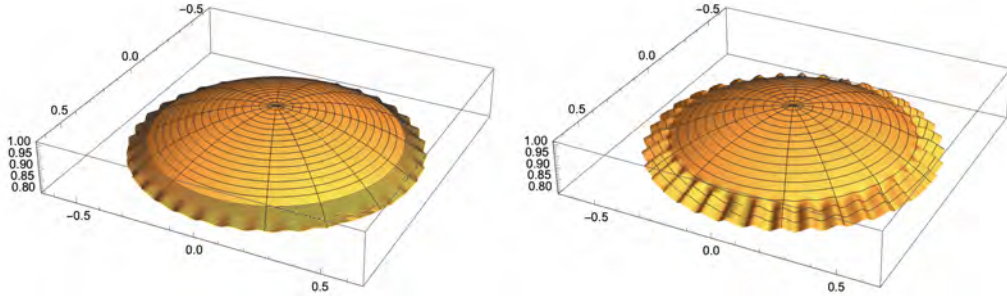


Figure 4: Corrugated surface produced by two possible choices of the perturbation function $f(\vartheta)$: unit step function, Eq. (6) (left) and hyperbolic tangent, Eq. (7) (right). In both cases the same opening of the dome ϑ_f and unperturbed radius R_0 have been assumed.

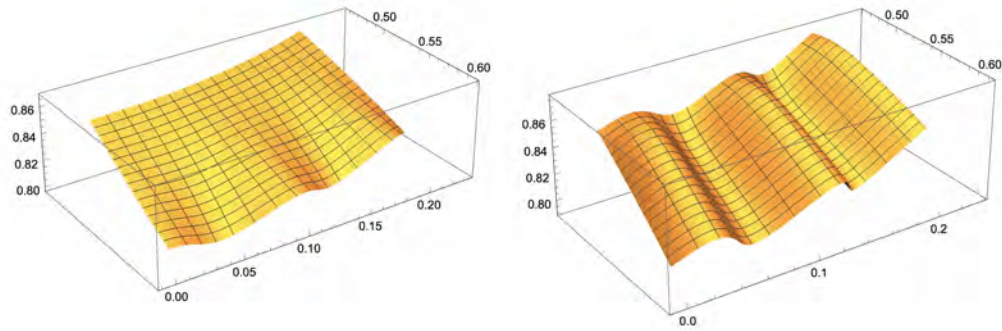


Figure 5: Magnified portion of the corrugated edge for the two cases presented in Eq. (6) (left) and Eq. (7) (right).

193 3. Generating a suitable geometry for FE computations

194 Starting from the above introduced parametric description of the corru-
 195 gated shell surface, it is now possible to generate a geometry which is suitable
 196 for the subsequent either linear or non-linear analyses. Of course to this aim,

197 geometry formulation must be accurate. Indeed, in non-linear analyses any
198 imperfection would result in a sudden reduction of the critical load. It is com-
199 mon to introduce slight imperfections (related to the geometry) to trigger an
200 equilibrium path bifurcation in large-displacement or buckling analyses [44–
201 48]. Now, a standard procedure to create a geometric model adopts usually
202 a flat-faceted surface generated by CAD software. This does not guarantee
203 that geometrical accuracy can be achieved in subsequent computations. A
204 better option consists in using computational tools such as Non Uniform Ra-
205 tional Basis Spline functions (henceforth, NURBS) to model the surface. To
206 conceive a geometric object, the following steps must be followed. As a basic
207 criterion, given the cyclic symmetry of the surface, only a piece of surface
208 must be generated, for instance (in the present case) one of the slices lying
209 between two subsequent supports has been drawn. In Nervi’s dome, there
210 are 36 supports and each such slice spans exactly 10° . The procedure for
211 generating the geometry, which is described in Fig. 6, can be summarized as
212 follows:

- 213 i. A code has been developed in a geometric modelling software, whose
214 aim is to use the parametric equation of the surface to numerically
215 compute a satisfactory set of coordinate pairs (ϑ_i, φ_i) . The dimension
216 set depends on the specified number of points along the colatitude and
217 longitude direction. A cross-reference algorithm is employed to create
218 a pair (ϑ_i, φ_i) , representing a single point belonging to the surface.
219 Therefore, a numerical algorithm (available on the software library),
220 employing NURBS, is applied to the points set **which should better**
221 **approximate the shape of the geometric object**. The accuracy of the

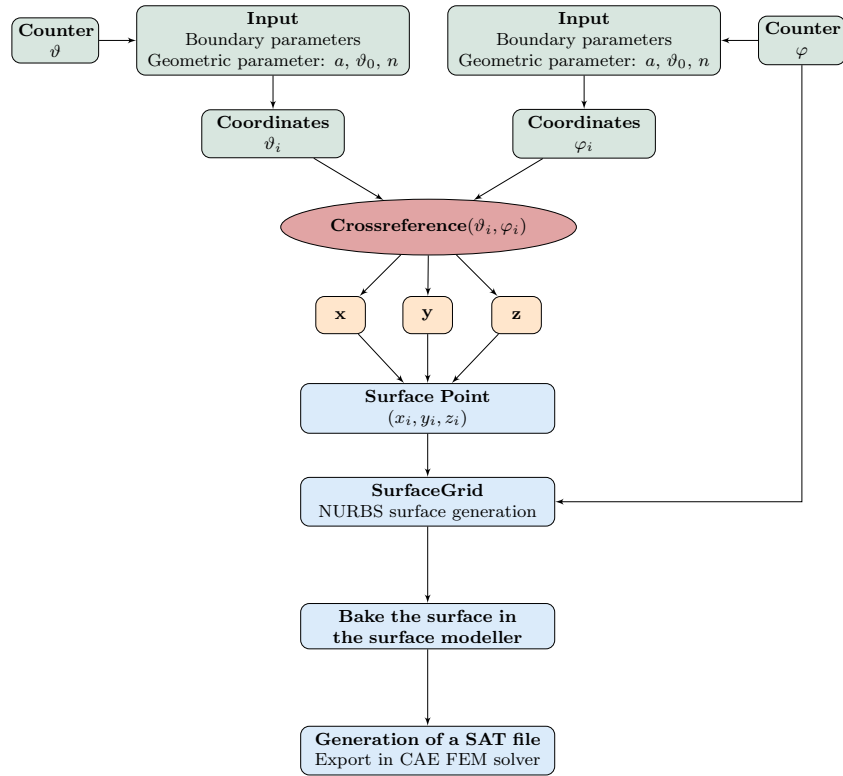


Figure 6: Block-diagram of surface generating algorithm.

222 representation depends on the initial user choice. Nevertheless, it is
 223 clear that the accuracy in generating the surface depends on the chosen
 224 number of points;

225 ii. The algorithm outcome is then graphically visualized in the surface
 226 modeler. To produce the obtained NURBS surface, it is necessary to
 227 bake it into the graphical interface;

228 iii. The surface is exported into a convenient file exchange format, such
 229 as the ACIS format. This file includes complete pieces of information

230 about the geometry and can be imported into an advanced FEM solver;
231 iv. The ACIS file is imported into the CAE FEM solver and constitutes a
232 single portion of the surface, see Fig. 7. Inside the pre-processor, the
233 portion is suitably reflected and then a circular pattern is implemented
234 around the z -axis to generate the whole dome. All these slices need
235 then to be merged into a single object.

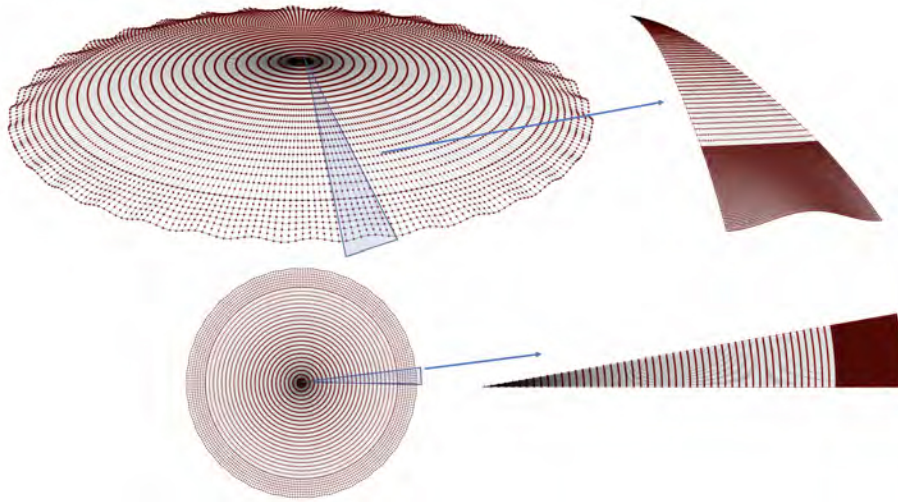


Figure 7: The automatically generated set of points for the model of Nervi's dome of Palasport Flaminio.

236 The obtained part can then be meshed inside the FE pre-processor. It is
237 requested that even the mesh closely reproduces the shape without deforming
238 the geometry. As a consequence, the chosen type of finite element should be
239 able to reproduce a conveniently small portion of double-curvature surface
240 such as an eight-noded shell element with six or five degrees of freedom, and
241 where that is not possible (e.g. in the dome apex) to six-noded triangular

242 elements with five degrees of freedom. In addition, this kind of elements is
243 suitable to perform buckling analysis. However, to obtain a regular mesh,
244 a sweep algorithm has been used for getting a mesh made of quadrilateral
245 elements, whose edges are aligned with meridians and parallels. The resulting
246 mesh shows however a main drawback: indeed, the elements which are close
247 to the vertex turn out to be severely distorted. In order to overcome such
248 problem, a minute partition is applied near the apex. Hence, in this region
249 a single six-noded triangular element is employed for each slice.

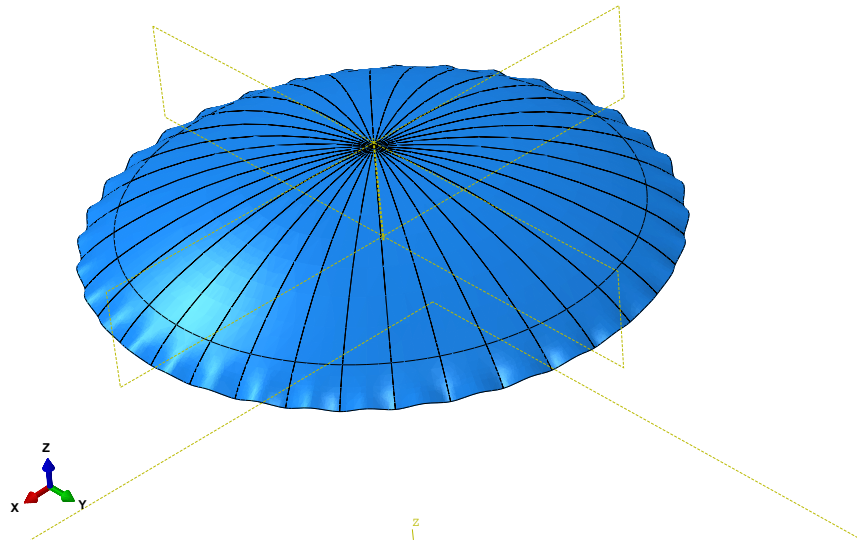


Figure 8: Resulting assembly of the slices of the model of Palasport Flaminio dome.

250 4. Geometrical and mechanical data

251 The analysis of a shell inspired by RC Palasport Flaminio dome designed
252 by Nervi has been accomplished. In this section the geometrical and mechan-

Table 1: Geometric properties of the Palasport Flaminio dome

Radius	R_0	51.039 m
Roof span	L	60.000 m
Shell thickness	t	0.200 m
Opening angle	ϑ_f	$\pi/5$ rad
Angle where perturbation starts	ϑ_0	$\pi/6$ rad
Wave number	n	36

253 ical assumptions adopted for such RC shell are summarized. Original design
 254 blueprints are available in the MAXXI (the Art Museum of the XXI century)
 255 archives in Rome. Besides, the span and the opening angle are given in [18],
 256 as reliable average measurements. All these data are presented in Table 1.

257 **In particular, a sketch of the geometry is provided in Fig. 9.**

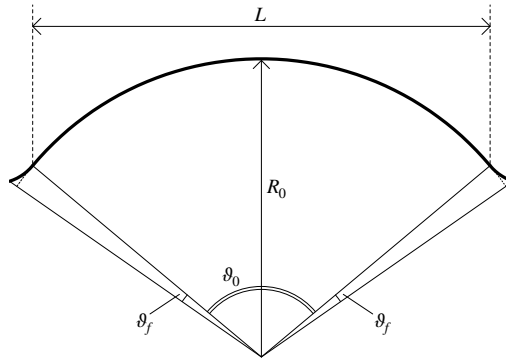


Figure 9: **Schematic representation of the shell geometry.**

258 In [49] a different opening angle has been considered, which, however,
259 does not seem to agree with the architectural blueprints reported in [18].
260 Instead mechanical characterization of the material has been simplified as
261 follows:

- 262 i. Only the linear part of the stress-strain curve is taken into account. So,
263 plasticity and other kind of non-linear behaviour have been disregarded;
- 264 ii. RC is always assumed to be uncracked;
- 265 iii. Time dependent effects on material properties are completely neglected,
266 and Young's modulus E is assumed to be equal to 30 GPa;
- 267 iv. Poisson's coefficient ν is similarly assumed to be constant and equal to
268 0.2;
- 269 v. Concrete strength class is required to be C20/C25 (Eurocode classi-
270 fication), whose density is 2500 kg/m³. Even though Nervi's dome
271 is a RC shell with stiffeners, which follow a certain spherical path, it
272 is reasonable to assume that shell thickness is constant and equal to
273 0.2 m.

274 For a more accurate survey about mechanical properties of concrete, partially
275 based on tested specimens and survey, even though they are referred to a
276 nearby canopy which was realised by the same Nervi (Stadium Flaminio) a
277 reader may rely on [50].

278 In the mentioned reference, particular emphasis has been placed on the
279 differences between *ferrocemento*, the building material adopted by Nervi,
280 and the standard RC.

281 **5. Analysis**

282 It is now possible to carry out a statical analysis of the structure by using
283 computational methods. For the sake of simplicity, only the shape produced
284 by Eqs. (4), (5) and (6) has been considered.

285 The load-case consists of a uniform external normal pressure q_0 applied
286 inwards to the whole surface, whose magnitude is 5 kPa. **Since the extension**
287 **of the shell is such that its rise over span ratio makes it rather shallow,**
288 **it is possible to approximate the self-weight load condition with a uniform**
289 **pressure. This approximation allows tackling the problem of stability of**
290 **edge-corrugated shells in an easier framework.** To understand the effect of
291 edge-corrugation, a comparison between a shell with the same geometrical
292 measure but without corrugation is shown.

293 The structure has been envisioned to behave according to the classical
294 membrane theory. In agreement with this theory, supports should restrain
295 motion only along the tangent direction. According to that, the dome edge re-
296 quires to be simply supported, so that edge rotation and out-of-plane surface
297 extension/shrinking are allowed. Therefore, the designer put all his efforts
298 to guarantee that no edge disturbances occur. As a matter of fact, the pillar
299 inclination follows the tangent to the boundary surface. Nevertheless, the
300 edge is fully restrained on the support and free between the supports. If
301 other constraints are applied, the membrane state will be supplemented by
302 bending and twisting effects. As a consequence, such supplementary effects
303 should be taken into account.

304 In terms of the membrane stress resultants N_ϑ , N_φ and $N_{\vartheta\varphi} = N_{\varphi\vartheta}$, which
305 are, respectively, the in-plane normal components of stress directed along the

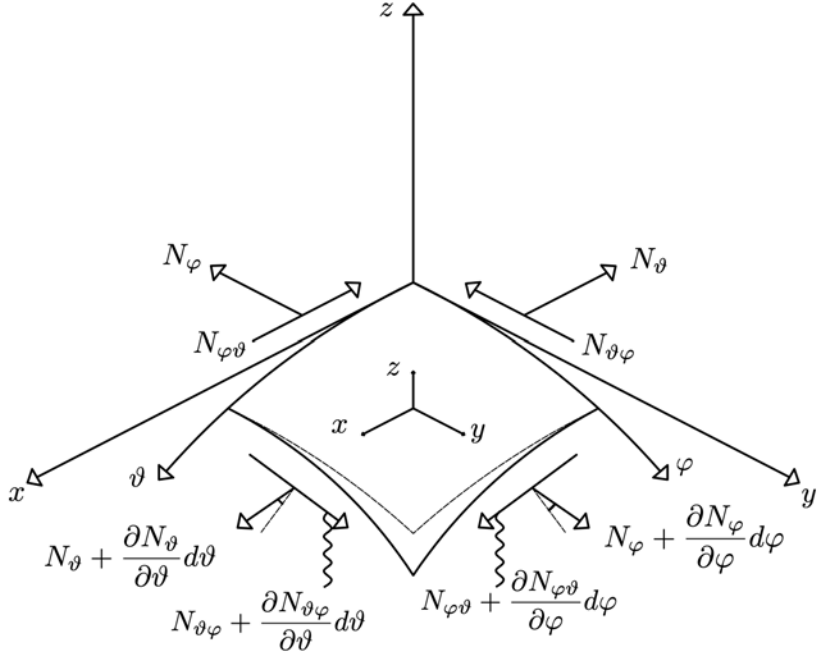


Figure 10: Equilibrium of a differential element of a shell with membrane stress-resultants.

306 meridian line and the parallel line, and the in plane shear stress components
 307 (see Fig. 10), the fundamental differential equations of equilibrium in the
 308 membrane theory of shells, written with reference to the lines of curvature
 309 on the middle-surface are given by [51] and read:

$$\begin{aligned}
 \frac{\partial}{\partial \vartheta}(BN_{\vartheta}) - \frac{\partial B}{\partial \vartheta}N_{\varphi} + \frac{1}{A} \frac{\partial}{\partial \varphi}(A^2N_{\vartheta\varphi}) + ABX &= 0 \\
 \frac{\partial}{\partial \varphi}(AN_{\varphi}) - \frac{\partial A}{\partial \varphi}N_{\vartheta} + \frac{1}{B} \frac{\partial}{\partial \vartheta}(B^2N_{\vartheta\varphi}) + ABY &= 0 \\
 \frac{N_{\vartheta}}{R_1} + \frac{N_{\varphi}}{R_2} - Z &= 0.
 \end{aligned} \tag{9}$$

310 In Eq. (9) A and B are the coefficients of the first fundamental form, which

311 gives the squared length of a line element as:

$$ds^2 = A^2 d\vartheta^2 + B^2 d\varphi^2;$$

312 R_1 and R_2 are, respectively, the curvature radii along the meridian (ϑ -line)
313 and perpendicularly to it; X and Y are external surface loads (i.e. loads per
314 unit area) acting towards increasing values of ϑ and φ , while Z denotes the
315 intensity of surface load per unit area acting along the outward normal.

316 So, for a perfectly spherical shell, whose radius is R_0 , acted upon by
317 a constant external, inward-directed pressure q_0 , the previous equations do
318 simplify because $R_1 = R_0$, $R_2 = R_0$, $A = R_0$, $B = R_0 \sin \vartheta$; moreover $X = 0$,
319 $Y = 0$, $Z = -q_0$. Finally for the axial symmetry, it is everywhere $N_{\vartheta\varphi} = 0$
320 and $N_\vartheta = N_\vartheta(\vartheta)$, $N_\varphi = N_\varphi(\vartheta)$, i.e. they do depend only on colatitude.

321 The reference solution for a membrane state produced by a uniform pres-
322 sure load q_0 acting on a hemispherical shell supported along the equator is
323 well-known and can be found, if attention is restricted to some of the major
324 sources only, in [52–54]. Indeed it results:

$$325 \quad N_\vartheta = N_\varphi = -q_0 \frac{R_0}{2} \quad (10)$$

326 It should be emphasized that N_ϑ and N_φ , which are given by Eq. (10) and
327 will be used in the sequel as a measure of variance from a perfect membrane
328 state, represent the resultant of the corresponding local stress components
329 σ_ϑ , σ_φ , once they are integrated along the thickness of the shell, t . Important,
330 both theoretical and experimental references about RC shells behaviour can
331 be found in [55–57]. On the other hand, a theoretical solution in terms
332 of Fourier series for a hemispherical shells on discrete supports was set out
333 in [52, 56].

334 The numerical solutions for the edge-corrugated spherical shells produce local
335 surface stresses, whose general expressions can be computed in terms of the
336 section resultants N_ϑ and N_φ as

$$\sigma_\vartheta = \frac{N_\vartheta}{t}, \quad \sigma_\varphi = \frac{N_\varphi}{t}.$$

337 In particular, σ_ϑ is the normal stress acting along the ϑ direction (i.e. that
338 tangent to the meridian) and σ_φ is the normal stress along the φ direction,
339 namely tangent to the parallel.

340 It should be noticed that, even for the corrugated shell, the above men-
341 tioned directions are principal direction of stress for the considered applied
342 load, namely uniform external pressure. All numerical results related to stress
343 are presented in dimensionless form: stress values are indeed divided by the
344 applied external pressure $q_0 = 5$ kPa, while the angular position is given in
345 the dimensionless form ϑ/ϑ_f , where ϑ_f is the colatitude value corresponding
346 to the position of the edge, which is assumed to be the same in all considered
347 cases.

348 The stresses σ_ϑ and σ_φ along the shell are displayed in the contour plot
349 on the left of Fig. 11 and Fig. 12.

350 As it is expected, the solution exhibits cyclic symmetry. A comparison with
351 a non-corrugated shell is displayed on the right of Fig. 11 and of Fig. 12.

352 The stresses corresponding to two meridians, one passing through a support
353 and the other through the crest of a wave will be considered in detail; the
354 output is concisely shown in the graphs that follow.

355 In particular, Fig. 13 and Fig. 14 show on the left σ_ϑ for the above mentioned
356 meridians, while the stress component σ_φ is shown, for the same meridians,

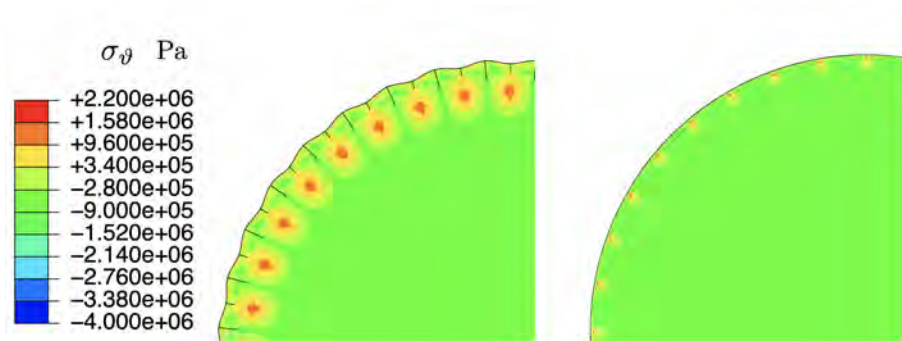


Figure 11: Contour plot of the stress component σ_ϑ for the corrugated-edge dome (left) and for the spherical cap without corrugation (right). In both cases, 36 discrete supports have been considered.

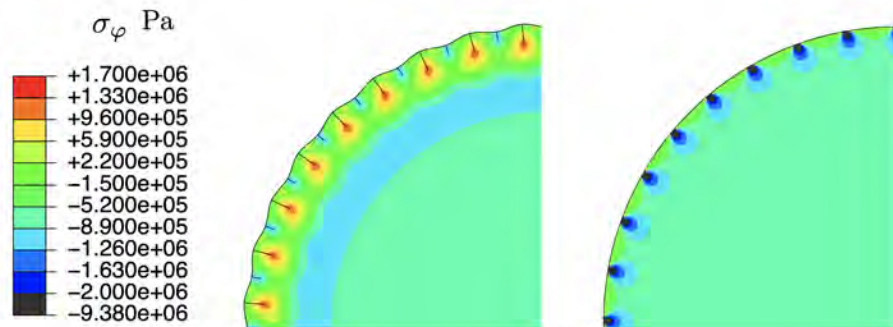


Figure 12: Contour plot of the stress component σ_φ for the corrugated-edge dome (left) and for the spherical cap without corrugation (right). In both cases, 36 discrete supports have been considered.

357 on the right of the above mentioned Figures.
 358 To understand the reason why in the corrugated shell there is a reduction
 359 of stresses, one can usefully look at the bending moment diagrams. Let
 360 M_ϑ be the section moment (which is dimensionally expressed as the ratio
 361 moment/thickness, thus being homogeneous to a force) along the ϑ direction
 362 and M_φ the section moment along the φ direction. Fig. 15 shows the contour

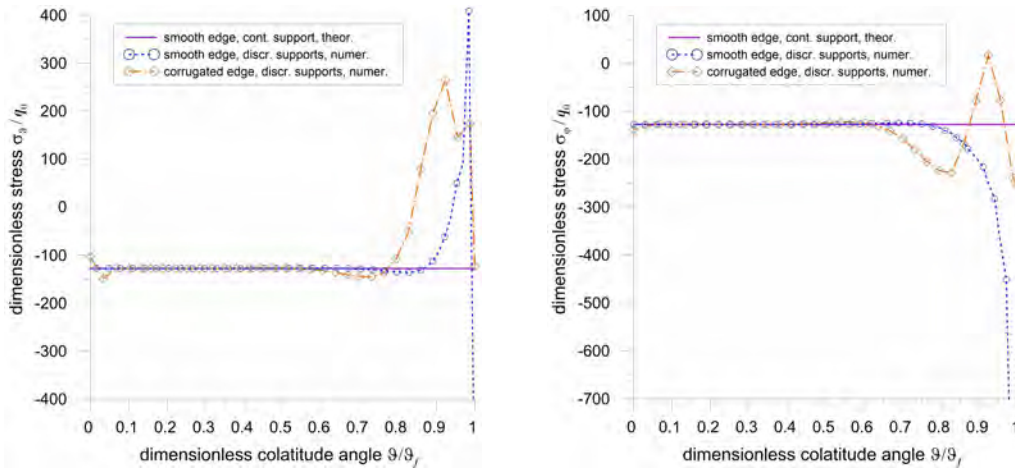


Figure 13: Stress components σ_θ and σ_φ along a meridian passing through a support.

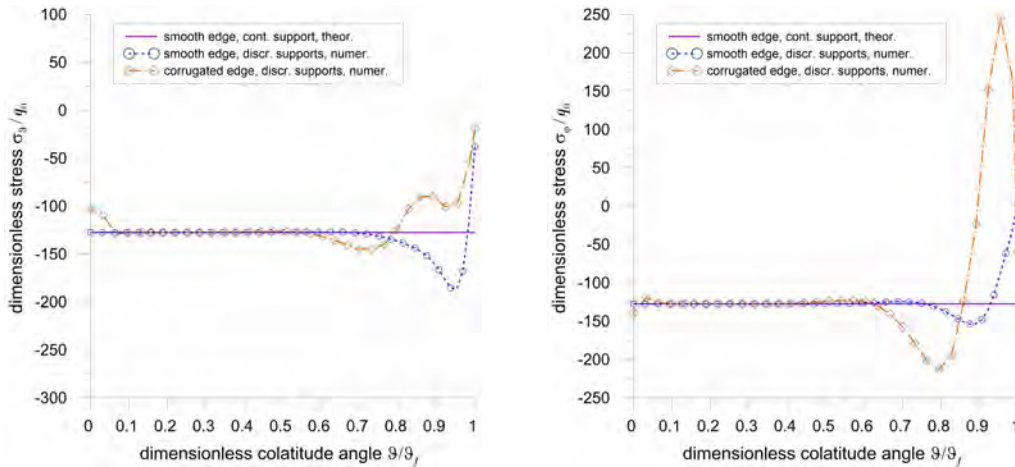


Figure 14: Stress components σ_θ and σ_φ along a meridian located between two supports, i.e. corresponding to the crest of the wave.

363 plot of the section moment M_θ for an edge-corrugated (left) and for a non
 364 corrugated shell (right).

365 Similarly, Fig. 16 shows the contour plot of the section moment M_φ for an
 366 edge-corrugated (left) and for a non corrugated shell (right).

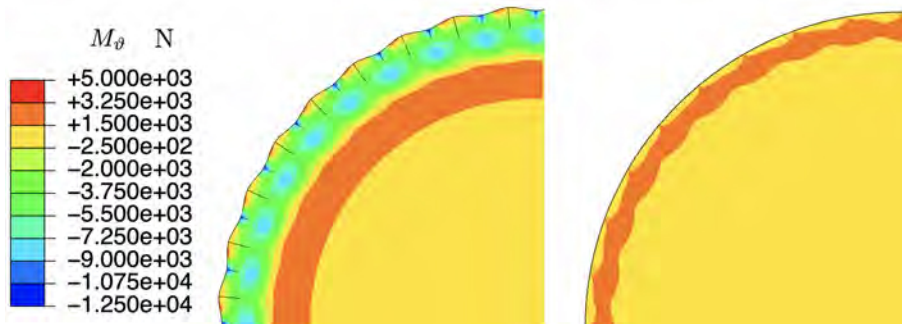


Figure 15: Contour plot of the section moment, M_θ for the corrugated-edge dome (left) and for the spherical cap without corrugation (right). In both cases 36 discrete supports have been considered.

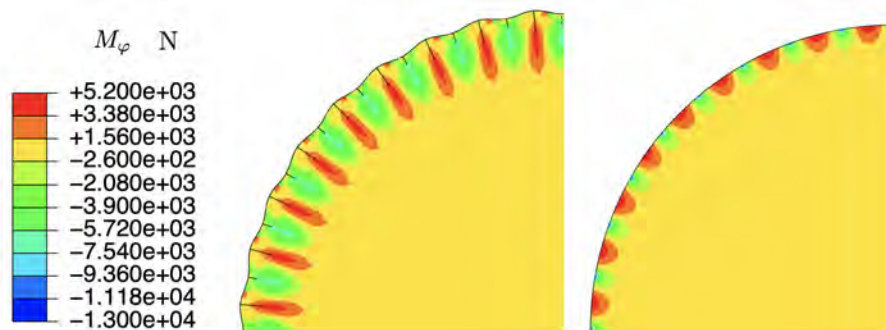


Figure 16: Contour plot of the section moment M_ϕ for the corrugated-edge dome (left) and for the spherical cap without corrugation (right). In both cases 36 discrete supports have been considered.

367 The section moments have been plotted for two different meridians, one pass-
 368 ing through a support, the other through the crest of the wave. The output
 369 is shown in Fig. 17 and 18 in dimensionless form, by dividing the relevant
 370 values by the constant $M_0 = 5000$ kN.

371 M_θ (left) and M_ϕ (right) are plotted in Fig. 17 with reference to the meridian
 372 passing through one of the supports; instead the same section moments, in

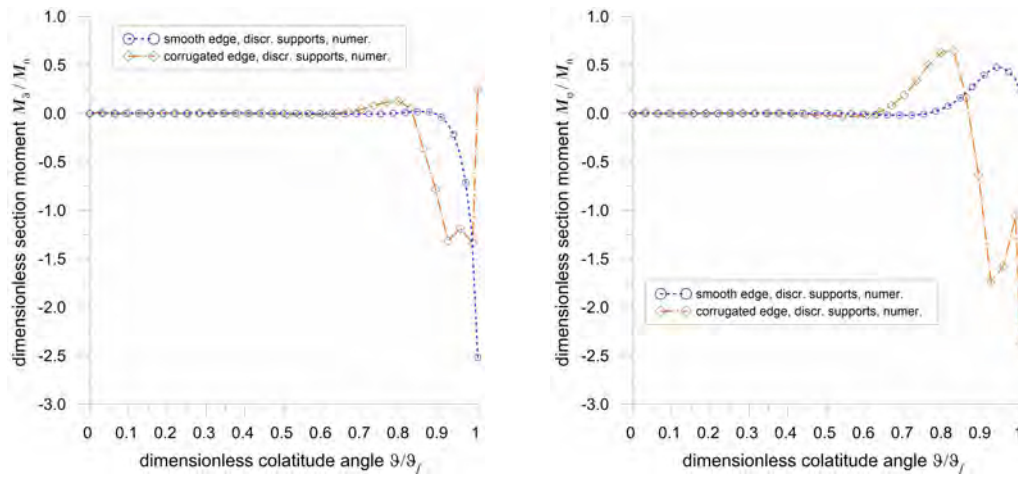


Figure 17: Section moments M_θ and M_φ along a meridian passing through a support.

373 the same order, are plotted in Fig. 18 for a meridian which passes through
 374 the crest of the corrugation.

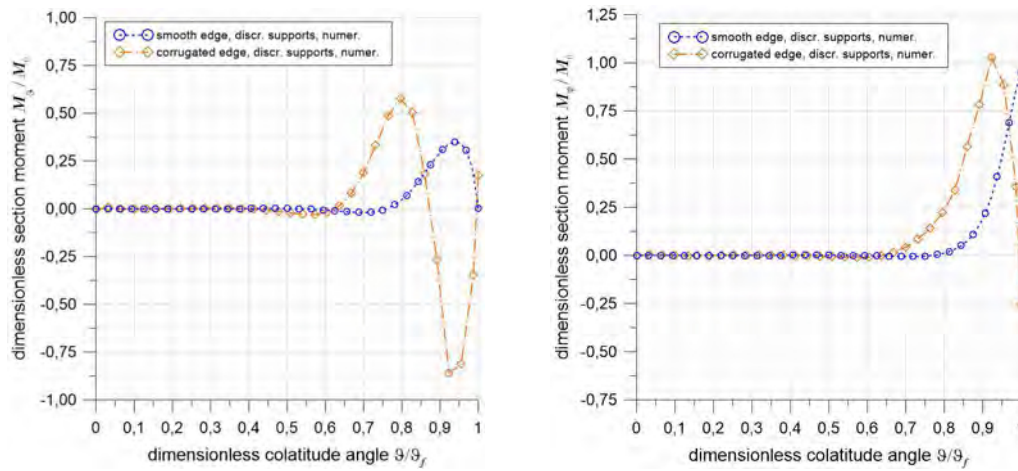


Figure 18: Section moments M_θ (left) and M_φ (right) along a meridian located between two supports, i.e. corresponding to the crest of the wave.

375 It can be pointed out the **significant** decrease, along the meridian direction,
 376 of the bending moment in the corrugated shell, in comparison with the non-

377 corrugated one. That is the most relevant issue. Indeed corrugation, due to
378 its shape, allows for a significant reduction of the bending moment: close to
379 the edge M_{ϑ} exhibits a noteworthy reduction involving even an inversion of
380 its signum. This improvement can be clearly perceived as a consequence of
381 the stiffness enhancement, which not only depends on the thickness of the
382 shell but also on its shape, i.e. on the presence of a wavy edge. Therefore,
383 edge corrugation enhances the mechanical performance of the shell, without
384 increasing its thickness. On the other hand, M_{φ} is not similarly subjected
385 to a significant decrease since the surface is essentially warped in the φ di-
386 rection. Nonetheless, from the point of view of the designer, this difference
387 of behaviour in terms of bending moments has little significance, since for
388 RC shells the standard design rules suggest to adopt symmetric steel bar
389 reinforcement along both ϑ and φ directions.

390 6. Conclusion

391 The parametric equations for a shell whose edge is corrugated have been
392 proposed and a suitable FE simulation procedure, which accurately handles
393 the doubly curved geometry, has been presented.

394 It has been shown through numerical simulation what was intuitively
395 clear to P.L. Nervi: a corrugation along the edge enhances the structural
396 performance of a shell. Furthermore, corrugation considerably decreases the
397 bending moment produced by discrete supports along the shell edge. Indeed,
398 although the shell is designed to behave as a perfect membrane, it can be
399 affected by significant bending on its edge. A reliable procedure has been
400 introduced to study the influence of corrugation into the above-mentioned

401 structures and to evaluate their stress distribution.

402 The original architectural function related to P.L. Nervi's work, which
403 represents the inspiration for the present paper, has been developed in the
404 sense of Civil and Structural Engineering: corrugated shell structures can be
405 both simply used as a canopy endowed also with some aesthetics, and can
406 be introduced in buildings for their high-performance mechanical properties,
407 like in the case of roofs for special and nuclear waste containers. In this case,
408 a unique shallow shell element must cover the vessel.

409 Other applications can be found in the field of automatization of building
410 processes. The procedure, which has been presented in this paper, may be ap-
411 plied to different shapes, such as free-form shells and concrete printed struc-
412 tures. Some improvements could be achieved in the LIDAR (Laser Imaging
413 Detection And Ranging) field to accurately identify the influence of small de-
414 viation in the structural behaviour; a comparison between a theoretical shape
415 and *in situ* surveys could be done, see for instance [58]. A main issue, which
416 is related to the topic discussed in the present work, will be developed in
417 the following investigations and involves the role of corrugation in instability
418 phenomena, such as snap-through. At this stage, it was intended indeed to
419 highlight the statical intuition envisioned by P.L. Nervi. In this preliminary
420 paper the effect of edge-corrugation is assessed only through linear elastic
421 stress analysis. Under this point of view, corrugation is not an imperfection
422 but is devised in such a way as to improve the structural performance. In
423 a forthcoming paper, buckling will be explicitly addressed and, for this pur-
424 pose, it will be necessary to take into consideration imperfections in order
425 to trigger the instability phenomena [59–61]. Indeed, the effect of increasing

426 bending stiffness in the spherical shell with respect to buckling problems has
427 been already addressed in the literature [62]. It constitutes however a very
428 challenging problem, whose solution has not been achieved yet.

429 The methods developed in this paper for Civil Engineering and Architec-
430 ture applications can be simply generalized to be used in different scientific
431 milieux. One can, for example, use them in the field of Bioengineering for
432 designing new smart contact lenses. This will be the topic of a future devel-
433 opment. Moreover, the presented study can be completed and enhanced by
434 taking into account results originally developed for the so-called generalized
435 theories: many interesting phenomena may arise if one introduces higher
436 gradient models like in [63–67].

437 Numerics play a fundamental role in the solution of problems like that
438 presented in this article: different numerical methods can help in studying
439 different phenomena and they could be implemented also in the study of
440 corrugated shells. One might refer to [68–73] for a detailed discussion.

441 **7. Declaration of competing interest**

442 The authors declare that they have no known competing financial inter-
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