A FULLY-DISCRETE-STATE KINETIC THEORY APPROACH TO MODELING VEHICULAR TRAFFIC*

LUISA FERMO† AND ANDREA TOSIN‡

Abstract. This paper presents a new mathematical model of vehicular traffic, based on the methods of the generalized kinetic theory, in which the space of microscopic states (position and speed) of the vehicles is genuinely discrete. While in the recent literature discrete-velocity kinetic models of car traffic have already been successfully proposed, this is, to our knowledge, the first attempt to account for all aspects of the physical granularity of car flow within the formalism of the aforesaid mathematical theory. Thanks to a rich but handy structure, the resulting model allows one to easily implement and simulate various realistic scenarios giving rise to characteristic traffic phenomena of practical interest (e.g., queue formation due to roadwork or to a traffic light). Moreover, it is analytically tractable under quite general assumptions, whereby fundamental properties of the solutions can be rigorously proved.

Key words. traffic granularity, generalized kinetic theory, vehicle interactions, stochastic games

AMS subject classifications. 34A34, 35L65, 76P05, 90B20

DOI. 10.1137/120897110

1. Introduction. Modeling vehicular traffic at the kinetic scale offers several advantages in terms of capturing some of the relevant aspects of the complexity of car flow. The kinetic approach is indeed suitable for an aggregate representation of the distribution of vehicles, not necessarily focused on single cars, while still allowing for a detailed characterization of the microscopic vehicle-to-vehicle dynamics, which are ultimately responsible for the large-scale behavior of the system. In this respect, kinetic models are competitive over microscopic ones; see, e.g., [13, 14, 19]. In fact, they require a lower number of equations, which makes them more amenable to computational and mathematical analysis as far as the study of average collective trends is concerned. On the other hand, it should be pointed out that any aggregate representation of traffic flow, including therefore classical kinetic ones, suffers from the problem that the number of vehicles distributed along the typical length scale of a road is in general not so high locally in order for statistical or continuous approaches to be fully justified. Gas dynamic kinetic models share such a conceptual difficulty with fluid dynamic macroscopic models, since in both cases cars are assumed to be continuously distributed in space in such a way that a description in terms of a density (possibly a probability density from a mesoscopic point of view) per unit length of the road is possible. In addition, macroscopic models usually require the prescription of a phenomenological constitutive relationship linking the car flux to the density, the socalled fundamental diagram, in order for the principle of conservation of cars along the road to be sufficient for determining the dynamics of the system [6, 16, 17]. Kinetic models replace instead such an empirical closure of the equations with a smaller-scale

^{*}Received by the editors October 30, 2012; accepted for publication (in revised form) May 6, 2013; published electronically July 16, 2013.

http://www.siam.org/journals/siap/73-4/89711.html

[†]Department of Mathematics and Computer Science, University of Cagliari, 09123 Cagliari, Italy (fermo@unica.it).

[‡]Istituto per le Applicazioni del Calcolo "M. Picone," Consiglio Nazionale delle Ricerche, 00185 Rome, Italy (a.tosin@iac.cnr.it). This author's work was partially funded by the European Commission under the 7th Framework Program (grant 257462 HYCON2, Network of Excellence) and by the Google Research Award "Multipopulation Models for Vehicular Traffic and Pedestrians," 2012–2013.

statistical modeling of microscopic interactions among cars. Actually, we mention for the sake of completeness that not all macroscopic models rely strictly on fundamental diagrams. Some get rid of them by envisaging closures of the balance equations oriented to a phenomenological interpretation of the driver behavior [1, 4]. For a more comprehensive overview of vehicular traffic models at all scales, the interested reader is referred to [2].

Classically, the kinetic representation of vehicular traffic along a one-way road is provided by a distribution function f over the mechanical microscopic state of the vehicles. The latter is identified by the scalar position $x \in D_x$ and speed $v \in D_v$, where $D_x, D_v \subseteq \mathbb{R}$ are the spatial and speed domains, respectively. While the former may either be a bounded interval, such as [0, L], L > 0 being the length of the road, or coincide with the whole real axis, the latter is normally of the form $D_v = [0, V_{\text{max}}]$, where $V_{\text{max}} > 0$ is the maximum speed allowed along the road, or possibly $D_v = [0, V'_{\text{max}}]$, with $V'_{\text{max}} \geq V_{\text{max}}$ corresponding to the maximum average speed attainable by a single vehicle in free flow conditions. The distribution function $f = f(t, x, v) : [0, T] \times D_x \times D_v \to \mathbb{R}^+$, where T > 0 is the final time (possibly $T = +\infty$), is then such that f(t, x, v) dx dv is the (infinitesimal) number of vehicles that at time t are located between t and t and travel with a speed between t and t and t and t and t are located between t and t and t and t and t are located between t and t and t and t are located between t and t and t and t are located between t and t and t and t are located between t and t and t and t and t are located between t and t and t and t are located between t and t are located between t and t are located between t and t and t are located between t and t are

According to the typical custom of gas-kinetic theories, in the above presentation the spatial position and speed of the vehicles are tacitly assumed to be continuously distributed over $D_x \times D_v$. However, as already remarked at the beginning of this introduction, this does not correctly reflect the physical reality of vehicular flow. Indeed, the number of vehicles along a road is normally not large enough for the continuity of the distribution function over the microscopic states to be an acceptable approximation. Vehicles do not continuously span the whole set of admissible speeds; likewise, one cannot expect that they are so spread in space that at every point x of the road there may be some. In other words, the actual distribution of vehicles in space, as well as that of their speeds, is strongly granular. It is reasonable to expect this fact to have a nontrivial impact on the resulting dynamics. It is hence worth taking into account in a mathematical model.

Recently discrete velocity models have been introduced [3, 7, 11]. The idea is to relax the hypothesis that the speed distribution is continuous by introducing in the domain D_v a lattice of discrete speeds $\{v_1, v_2, \ldots, v_n\}$, with $v_1 = 0$, $v_n = V_{\max}$, $v_i < v_{i+1}$ for all $i = 1, \ldots, n-1$. The physical system is then described by n distribution functions $f_j = f_j(t, x) : [0, T] \times D_x \to \mathbb{R}^+$, $j = 1, \ldots, n$, such that $f_j(t, x) = f(t, x, v_j)$ or, in the distributional sense, $f(t, x, v) = \sum_{j=1}^n f_j(t, x) \delta_{v_j}(v)$. In particular, $f_j(t, x)dx$ is the (infinitesimal) number of vehicles traveling at speed v_j that, at time t, occupy a position between x and x + dx. In this way, the granular character of the car flow is at least partially taken into account from the point of view of the speed distribution.

Starting from the discrete velocity kinetic framework developed in the papers cited above, in the present work we undertake the discretization also of the microscopic space variable, in order to accomplish the program of a fully-discrete-state kinetic theory of vehicular traffic. The basic idea is to partition the spatial domain D_x into disjoint cells (subintervals) I_i such that $D_x = \bigcup_i I_i$. Each cell identifies a portion of the road in which only the number of vehicles is known. No matter how rough this assumption may seem, it is nevertheless realistic in view of the spatial granularity of traffic, which does not allow one to make finer predictions on the spatial position of

vehicles within each cell. The microscopic state (x, v) of the vehicles belongs then to the discrete state space $\{I_i\}_i \times \{v_i\}_i$, which generates a new structure of the kinetic equations not immediately deducible from more standard frameworks. Particularly, it is worth anticipating that the resulting mathematical structure is not a cellular automaton [8, 9, 12] despite the discreteness of the space state. In fact, as will be clear in the next sections, vehicles are not regarded as fictive particles jumping from their current site to a neighboring one with prescribed probability. They actually flow through the cells with their true speed, according to a transport term duly implemented in the time-evolution equations of their distribution functions. Car-to-car interactions, responsible for speed variations, are in turn coded in the aforesaid equations, so that finally the evolution of the system is not seen as a stepwise algorithmic update of the lattice of microscopic states. The final form of the proposed model can be rather compared to that of a master equation, namely a set of ordinary differential equations describing the time evolution of the probability that a generic representative vehicle of the system will be in one of the discrete microscopic states at each (continuous) time instant.

In developing the program just outlined, this paper aims at covering the whole path from modeling to numerical simulations to well-posedness analysis of the resulting mathematical problems. In more detail, section 2 derives the general discretestate equations, resting on the methods of the generalized kinetic theory [2, sect. 6], to be used as a reference framework for specific models. This framework accounts, in particular, for stochastic dynamics of speed variations, which model the not strictly deterministic, nor purely mechanical, behavior of drivers. Section 3 presents a model obtained from the previous framework, which is then used in section 4 for producing exploratory numerical simulations of some case studies. These include the long-term trend of the system, particularly with reference to the fundamental diagrams predicted by the kinetic equations, and the dynamics of queue formation/depletion along a road triggered by two different causes: roadworks and traffic lights. Finally, section 5 is devoted to the analytical study of some relevant qualitative properties of the initial/boundary-value problems. Under quite general assumptions, to be possibly regarded as modeling guidelines, global existence of the solution, uniqueness, continuous dependence on the initial and boundary data, nonnegativity, and boundedness are proved, which provide the modeling framework with the necessary degree of mathematical robustness.

2. Discrete space kinetic framework. In this section, starting from the discrete velocity setting recalled in the introduction, we derive kinetic structures based on a fully discrete space of the mechanical microscopic states of the vehicles.

As previously anticipated, partitioning the spatial domain D_x into a number of cells of finite size is a more realistic way of detecting the positions of the vehicles along the road. Indeed, it is consistent with the intrinsic granularity of the flow, which does not allow for a statistical description of their spatial distribution more accurate than a certain minimum level of detail. In addition, as we will see, it enables one to account easily for some effects due to the finite size of the vehicles even in a context in which the actual representation is not focused on each of them. To this purpose, a useful partition of the spatial domain D_x is into pairwise disjoint cells I_i , whose union is D_x . For a road $D_x = [0, L]$ of length L > 0, we set

$$D_x = \bigcup_{i=1}^m I_i, \qquad I_{i_1} \cap I_{i_2} = \emptyset \quad \forall i_1 \neq i_2,$$

 $m \in \mathbb{N}$ being the total number of cells needed for covering D_x , which depends on the size ℓ_i of each I_i . We will henceforth assume that cells have a constant size ℓ , chosen in a such a way that $L/\ell \in \mathbb{N}$; thus $m = L/\ell$.

Let $f_{ij} = f_{ij}(t)$ be the distribution function of vehicles that, at time t, are located in the ith cell with a speed in the jth class. Then the total number N_{ij} of vehicles in I_i with speed v_j is $N_{ij} = f_{ij}\ell$, while the total number N_i of vehicles in I_i is

$$N_i = \sum_{j=1}^{n} N_{ij} = \ell \sum_{j=1}^{n} f_{ij}.$$

Notice that $\rho_i := N_i/\ell$ is instead the local density of vehicles in the *i*th cell. By further summing over *i*, one gets the total number *N* of vehicles along the road:

(2.1)
$$N = \sum_{i=1}^{m} N_i = \ell \sum_{i=1}^{m} \sum_{j=1}^{n} f_{ij}.$$

Remark 2.1. If $N_{\text{max}} > 0$ is the maximum number of vehicles allowed in a single cell, then the maximum density is $\rho_{\text{max}} := N_{\text{max}}/\ell$.

From the f_{ij} 's, the kinetic distribution function f can be recovered as

(2.2)
$$f(t, x, v) = \sum_{i=1}^{m} \sum_{j=1}^{n} f_{ij}(t) \mathbb{1}_{I_i}(x) \delta_{v_j}(v),$$

 $\mathbb{1}_{I_i}$ being the characteristic function of the cell I_i ($\mathbb{1}_{I_i}(x) = 1$ if $x \in I_i$, $\mathbb{1}_{I_i}(x) = 0$ if $x \notin I_i$). In practice, f is an atomic distribution with respect to the variable v, like in the discrete velocity framework, and is piecewise constant with respect to the variable x. In particular, this latter characteristic implies that vehicles are thought of as uniformly distributed within each cell.

Usual macroscopic variables of traffic, such as the vehicle density ρ , flux q, and average speed u are obtained from (2.2) as distributional moments of f with respect to v:

(2.3)
$$\rho(t, x) = \sum_{i=1}^{m} \left(\sum_{j=1}^{n} f_{ij}(t) \right) \mathbb{1}_{I_{i}}(x),$$

$$q(t, x) = \sum_{i=1}^{m} \left(\sum_{j=1}^{n} v_{j} f_{ij}(t) \right) \mathbb{1}_{I_{i}}(x), \quad u(t, x) = \frac{q(t, x)}{\rho(t, x)}.$$

Notice that ρ can also be written as $\rho(t, x) = \sum_{i=1}^{m} \rho_i(t) \mathbb{1}_{I_i}(x)$, thus showing that the vehicle density on the whole road is a piecewise constant function in x built from the local cell densities. Analogously, after defining the flux at the ith cell as $q_i := \sum_{j=1}^{n} v_j f_{ij}$, we obtain $q(t, x) = \sum_{j=1}^{m} q_i(t) \mathbb{1}_{I_i}(x)$.

2.1. Evolution equations for the f_{ij} 's. A mathematical model of vehicular traffic in a fully discrete kinetic context is obtained by deriving suitable evolution equations for the distribution functions f_{ij} . To this end, we propose a method grounded on classical ideas of conservation laws, which, however, directly exploits the discrete space structure as a distinctive feature of the present setting.

Fix a cell I_i and a speed class j. The basic idea is, as usual, that during a time span $\Delta t > 0$ the number N_{ij} of vehicles in that cell with that velocity may vary because of

- vehicles leaving the cell I_i or new vehicles entering it at speed v_j from neighboring cells,
- \bullet interactions among vehicles within the cell I_i inducing speed changes.

On the basis of these two principles, the following preliminary balance can be written:

$$(2.4) N_{ij}(t + \Delta t) - N_{ij}(t) = -N_{ij}^{\text{out}}([t, t + \Delta t]) + N_{ij}^{\text{in}}([t, t + \Delta t]) + J_{ij}(t)\ell\Delta t + o(\Delta t)$$

with obvious meaning of the symbols N_{ij}^{in} , N_{ij}^{out} . The term $J_{ij}(t)$ on the right-hand side accounts for instantaneous acceleration/braking dynamics per unit space ℓ , which produce transitions among the various speed classes. For the moment we keep it generic, deferring to the next section a detailed discussion on its structure. Moreover, the term $o(\Delta t)$ indicates that we are neglecting higher order effects which may contribute to the variation of N_{ij} during the time Δt .

If the time span Δt is sufficiently small, we can assume that the number of vehicles leaving the ith cell depends essentially on the number of vehicles already present in I_i at time t and on the free room in the next cell I_{i+1} . To be definite, let $v_j \Delta t \leq \ell$ for all j, so that no vehicle can move forward for more than one cell in a single Δt . Then the following form of N_{ij}^{out} can be proposed:

$$N_{ij}^{\text{out}}([t, t + \Delta t]) = \frac{v_j \Delta t}{\ell} \Phi_{i,i+1} N_{ij}(t) + o(\Delta t),$$

where

- the ratio $v_j \Delta t/\ell$ accounts for the percent amount of vehicles traveling at speed v_j , which are sufficiently close to the boundary between the *i*th cell (whence they are coming) and the (i+1)th cell (whither they are going) for changing cell during the time span Δt ;
- the term $\Phi_{i,i+1}$, that we name flux limiter across the interface separating the cells I_i , I_{i+1} , accounts for the percentage of vehicles in I_i that can actually move to I_{i+1} according to the free space available in the latter.

Likewise, the following form of N_{ij}^{in} can be proposed:

$$N_{ij}^{\text{in}}([t, t + \Delta t]) = \frac{v_j \Delta t}{\ell} \Phi_{i-1,i} N_{i-1,j}(t) + o(\Delta t),$$

where now the space dynamics involve the (i-1)th and the *i*th cells.

Plugging these expressions into (2.4) and rewriting in terms of the f_{ij} 's yields

(2.5)

$$f_{ij}(t + \Delta t)\ell - f_{ij}(t)\ell = -\frac{v_j \Delta t}{\ell} \left(\Phi_{i,i+1} f_{ij}(t)\ell - \Phi_{i-1,i} f_{i-1,j}(t)\ell \right) + J_{ij}(t)\ell \Delta t + o(\Delta t),$$

whence dividing by $\ell \Delta t$, taking the limit $\Delta t \to 0^+$, and rearranging the terms produces

(2.6)
$$\frac{df_{ij}}{dt} + \frac{v_j}{\ell} (\Phi_{i,i+1} f_{ij} - \Phi_{i-1,i} f_{i-1,j}) = J_{ij}.$$

This equation is in principle well defined only in the internal cells of the domain D_x , namely for $2 \le i \le m-1$. Indeed, at the boundary cells i=1, i=m the flux limiters $\Phi_{0,1}$, $\Phi_{m,m+1}$ call for the two nonexistent external cells I_0 , I_{m+1} . In addition, for i=1 the equation also requires the value f_{0j} . In order to overcome these difficulties,

we consider that the condition $v_j \geq 0$ for all j implies a unidirectional (rightward) flux of vehicles, which thus requires a condition on the left boundary of the domain. Consequently, the values $\Phi_{0,1}$ and f_{0j} have to be provided as boundary conditions. In contrast, the flux limiter $\Phi_{m,m+1}$ has to be prescribed at the right boundary in order to specify how the natural outflow of the vehicles is possibly modified by external conditions. In standard situations, it will simply be $\Phi_{m,m+1} = 1$.

Remark 2.2. The formal derivation of (2.6) closely recalls the standard one of conservation laws: the time variation of the conserved quantity is related to some drift plus microscopic conservative interactions typical of kinetic theory. Nevertheless, it is worth pointing out that classical procedures, such as the straightforward use of the divergence theorem, for deducing conservation laws in differential form cannot be applied in the present context. Indeed, the finite size of the space cells is a structural feature of the problem at hand, which prevents one from taking the limit $\ell \to 0^+$.

Remark 2.3 (conservativeness of (2.6)). The principle of conservation of cars implies that the time variation of the total number N of vehicles along the road, defined by (2.1), has to balance with the inflow/outflow of vehicles from the boundaries of D_x . We can formally check that model (2.6) indeed satisfies this property, after recalling that the conservative interaction operator J_{ij} is required to fulfill

(2.7)
$$\sum_{j=1}^{n} J_{ij}(t) = 0 \quad \forall i = 1, \dots, m, \ t > 0.$$

Multiplying (2.6) by ℓ and summing both sides over i, j gives

$$\frac{dN}{dt} + \sum_{i=1}^{n} v_{j} \left[\sum_{i=1}^{m} \left(\Phi_{i,i+1} f_{ij} - \Phi_{i-1,i} f_{i-1,j} \right) \right] = 0,$$

which, summing telescopically in the square brackets, yields the desired result:

$$\frac{dN}{dt} = -\Phi_{m,m+1} \sum_{j=1}^{n} v_j f_{mj} + \Phi_{0,1} \sum_{j=1}^{n} v_j f_{0j}$$
$$= -\Phi_{m,m+1} q_m + \Phi_{0,1} q_0,$$

 q_0 , q_m being the incoming and outgoing fluxes of vehicles through the boundaries x = 0 and x = L, respectively.

2.2. Stochastic dynamics of speed transitions. The right-hand side of (2.6) is the phenomenological core of the kinetic equations, because it accounts for the time variation of the conserved quantities f_{ij} in terms of microscopic interactions among vehicles. The operator J_{ij} is called *interaction operator*. It describes modifications of the speed v_j of a vehicle traveling in the cell I_i due to the action exerted on it by other vehicles.

Vehicles are not purely mechanical particles, for the presence of drivers provides them with the ability of taking decisions actively, i.e., without necessarily the influence of external force fields. Consequently, the description delivered by J_{ij} cannot be purely deterministic: personal behaviors have to be taken into account, which are suitably modeled from a stochastic point of view. Specifically, interactions among vehicles are understood as *stochastic games* among pairs of players. The game strategy of each player (viz. vehicle) is represented by its speed, while the payoff of the game is the new speed class it shifts to after the interaction. Games are stochastic because payoffs are

known only in probability. In other words, the model considers that vehicles might not react always the same way when placed in the same conditions due to possible partly subjective decisions of the drivers.

These ideas are brought to a formal level by writing J_{ij} as a balance, in the space of microscopic states $\{I_i\}_{i=1}^m \times \{v_j\}_{j=1}^n$, of vehicles gaining and losing the state (I_i, v_j) in the unit time:

(2.8)
$$J_{ij} = \frac{\ell}{2} \left(\sum_{h,k=1}^{n} \eta_{hk}(i) A_{hk}^{j}(i) f_{ih} f_{ik} - f_{ij} \sum_{k=1}^{n} \eta_{jk}(i) f_{ik} \right).$$

This expression is formally obtained by integrating, over a single space cell I_i , the interaction operator deduced in previous works dealing with the discrete-velocity kinetic theory of vehicular traffic (see, e.g., [18]). In this sense, (2.8) is consistent with other models relying on similar theoretical backgrounds.

We note the following:

- Only binary interactions among vehicles are accounted for, which possibly produce a change of speed but not of position of the interacting pairs. In particular, it is customary to term the *candidate* the vehicle which is likely to modify its current speed v_h into the *test* speed v_j (the payoff of the game) after an interaction with a *field* vehicle with speed v_k .
- The quantitative description of the interactions is provided by the terms $\eta = \{\eta_{hk}(i)\}_{h,k=1,\dots,n}^{i=1,\dots,m}$ and $A = \{A_{hk}^j(i)\}_{h,j,k=1,\dots,n}^{i=1,\dots,m}$, called the interaction rate and the table of games, respectively. The former models the frequency of the interactions among candidate and field vehicles. The latter gives instead the probabilities that candidate vehicles get the test speeds after interacting with field vehicles. (Hence A_{hk}^j is, for $j=1,\dots,n$, the probability distribution of the candidate player's payoff conditioned to the strategies v_h , v_k by which the two involved players approach the game.) Notice that both $\eta_{hk}(i)$ and $A_{hk}^j(i)$ may vary from cell to cell. Since the table of games is a discrete probability distribution, the following conditions must hold:

(2.9)
$$0 \le A_{hk}^{j}(i) \le 1 \qquad \forall h, j, k = 1, ..., n, \quad \forall i = 1, ..., m,$$

$$\sum_{j=1}^{n} A_{hk}^{j}(i) = 1 \qquad \forall h, k = 1, ..., n, \quad \forall i = 1, ..., m.$$

In particular, from (2.9) follows the fulfillment of (2.7).

• The coefficient $\ell/2$ appearing in expression (2.8) weights the above interactions over the characteristic length of a single space cell.

Remark 2.4 (generalization to nonlocal interactions). The form (2.8) of the interaction operator tacitly implies that candidate and field vehicles stay within the same space cell. However, one can also assume that the candidate vehicle interacts with field vehicles located in further cells ahead. This amounts to introducing a set of interaction cells

$$\{I_i, I_{i+1}, \dots, I_{i+\mu(i)}\}$$

extending from the *i*th cell up to $\mu(i)$ cells in front. The number $\mu(i)$ varies in general with the index *i* so as to satisfy the constraint $i + \mu(i) \leq m$ (i.e., interaction cells cannot lie beyond the last cell I_m). Consequently, the structure (2.8) of J_{ij} modifies

as

$$(2.10) J_{ij} = \frac{\ell}{2} \left(\sum_{l=i}^{i+\mu(i)} \sum_{h,k=1}^{n} \eta_{hk}(l) A_{hk}^{j}(l) f_{ih} f_{lk} w_l - f_{ij} \sum_{l=i}^{i+\mu(i)} \sum_{k=1}^{n} \eta_{jk}(l) f_{lk} w_l \right),$$

where $\{w_l\}_{l=i}^{i+\mu(i)}$ are suitable weights accounting for different effects of the interactions based on the distance of the *l*th cell from the *i*th one. They have to comply with some minimal requirements:

(2.11)
$$w_l \ge 0 \quad \forall l = i, \dots, i + \mu(i), \qquad \sum_{l=i}^{i+\mu(i)} w_l = 1 \quad \forall i = 1, \dots, m.$$

Notice that, for $\mu \equiv 0$, (2.11) implies that the only remaining weight w_i is equal to 1, which ultimately enables one to recover (2.8) as a special case of (2.10).

2.3. Nondimensionalization. At this point, it is useful to rewrite (2.6) in dimensionless form. To this end, we choose ℓ and V_{max} as characteristic values of length and speed, respectively, and we define the following nondimensional variables and functions:

$$v_j^* := \frac{v_j}{V_{\max}}, \qquad t^* := \frac{V_{\max}}{\ell}t, \qquad f_{ij}^*(t^*) := \frac{\ell}{N_{\max}}f_{ij}\left(\frac{\ell}{V_{\max}}t^*\right).$$

Substituting these expressions into (2.6), (2.8), we get, after some algebraic manipulations,

(2.12)

$$\frac{df_{ij}^*}{dt^*} + v_j^* \left(\Phi_{i,i+1} f_{ij}^* - \Phi_{i-1,i} f_{i-1,j}^* \right) = \sum_{h=1}^n \eta_{hk}^*(i) A_{hk}^j(i) f_{ih}^* f_{ik}^* - f_{ij}^* \sum_{k=1}^n \eta_{jk}^*(i) f_{ik}^*,$$

where

$$\eta_{hk}^*(i) := \frac{\ell N_{\max}}{2V_{\max}} \eta_{hk}(i)$$

is the dimensionless interaction rate.

It is worth noting that the new speed lattice is now

$$v_1^* = 0 < \dots < v_i^* < v_{i+1}^* < \dots < v_n^* = 1,$$

while the nondimensional length of the space cells becomes unitary. In addition, one can check that the dimensionless local cell density reads

(2.13)
$$\rho_i^*(t^*) := \frac{\ell}{N_{\text{max}}} \rho_i \left(\frac{\ell}{V_{\text{max}}} t^* \right) = \frac{N_i \left(\frac{\ell}{V_{\text{max}}} t^* \right)}{N_{\text{max}}},$$

whence $0 \le \rho_i^* \le 1$. Consequently, ρ_i^* can be directly thought of as a probability density.

We will henceforth refer to (2.12), omitting the asterisks on the dimensionless quantities for brevity.

3. From the general framework to particular models. Starting from the mathematical structures given in section 2, specific models can be produced by detailing the flux limiters, the interaction rate, and the table of games. This amounts to analyzing the dynamics of microscopic interactions occurring among vehicles.

Flux limiters. We recall that the term $\Phi_{i,i+1}$ limits the number of vehicles that can actually travel across the cells on the basis of the occupancy of the destination cells.

A prototypical form of such a limiter can be derived via the following argument. The free space left in the cell I_{i+1} for vehicles coming from the cell I_i is $N_{\max} - N_{i+1}$. Assuming $N_i > N_{\max} - N_{i+1}$, i.e., that the number of vehicles in the *i*th cell is greater than such an available space, the percentage of vehicles that can actually move is

$$\Phi_{i,i+1} = \frac{N_{\text{max}} - N_{i+1}}{N_i}.$$

On the other hand, if $N_i \leq N_{\text{max}} - N_{i+1}$, then in the cell I_{i+1} there is, in principle, enough room for all vehicles coming from the cell I_i ; thus in this case $\Phi_{i,i+1} = 1$.

Rewriting the flux limiter in terms of the local cell density via (2.13), we find

(3.1)
$$\Phi_{i,i+1} = \begin{cases} \frac{1 - \rho_{i+1}}{\rho_i} & \text{if } \rho_i + \rho_{i+1} > 1, \\ 1 & \text{if } \rho_i + \rho_{i+1} \le 1. \end{cases}$$

The interpretation is that if the total density in the pair of cells (I_i, I_{i+1}) is greater than the (dimensionless) maximum density $\rho_{\text{max}} = 1$, then the microscopic granularity of traffic starts acting and all vehicles cannot freely flow through the cells. Otherwise, the flux limiter has actually no effect.

Interaction rate. The frequency of the interactions can be assumed to depend on the number of vehicles in each cell, in such a way that the larger this number, the higher the frequency. In view of the indistinguishability of the vehicles, a possible simple form implementing this idea is

(3.2)
$$\eta_{hk}(i) = \eta_0 \frac{N_i}{N_{\text{max}}} = \eta_0 \rho_i,$$

where $\eta_0 > 0$ is a basic interaction frequency. Notice that this specific form does not depend explicitly on the speed classes of the interacting pairs. From now on we will invariably consider interaction rates of this kind, and we will use accordingly the simplified notation $\eta(i)$.

Table of games. Speed transitions within the space cells depend, in general, on the evolving traffic conditions, such as the local car congestion and the quality of the environment (e.g., road or weather). In this paper, we account for the former via the concept of fictitious density explained below, which aims at introducing in the table of games the active ability of drivers to anticipate the actual evolution of traffic. Conversely, we identify the latter simply by a parameter $\alpha \in [0, 1]$, with $\alpha = 0$ standing for the worst environmental conditions and $\alpha = 1$ for the best ones.

The concept of fictitious density, originally introduced in [10], is here revisited in the frame of the discrete space structure and of the stochastic game approach to vehicle interactions. We define the fictitious car density $\tilde{\rho}_i$ felt by drivers in the cell I_i as the following weighted average of the actual car densities in the cells I_i , I_{i+1} :

(3.3)
$$\tilde{\rho}_i := \begin{cases} (1-\beta)\rho_i + \beta\rho_{i+1} & \text{for } i = 1, \dots, m-1, \\ \rho_m & \text{for } i = m, \end{cases}$$

where $\beta \in [0, 1]$ is a parameter related to the anticipation ability of drivers. From (3.3) it results that if $\rho_{i+1} \leq \rho_i$, then $\tilde{\rho}_i \leq \rho_i$; thus if the car density in the cell ahead is lower than that in the current cell, drivers may be motivated to behave as if the traffic congestion in their cell were lower than the actual value. This implies, for instance, a higher inclination to maintain the speed, or even to accelerate, despite possibly insufficient local free space, because they are anticipating the nearby evolution of traffic. Conversely, if $\rho_{i+1} > \rho_i$, then we have $\tilde{\rho}_i > \rho_i$; hence if the car density in the cell ahead is higher than that in the current cell, drivers may be motivated to adopt more precautionary behavior, for instance by decelerating in spite of the real local free space. Notice that $\tilde{\rho}_i = \rho_i$ if and only if $\beta = 0$ or $\rho_{i+1} = \rho_i$. In these cases, the driver's anticipation ability has no effect because of either a specifically modeled behavior ($\beta = 0$) or the instantaneous traffic evolution ($\rho_{i+1} = \rho_i$).

The table of games stemming from these ideas is described in the following. Three cases need to be dealt with, corresponding to candidate vehicles h traveling more slowly, or faster, or at the same speed as field vehicles k.

• v_h < v_k. In this case, the candidate vehicle may either decide to maintain its speed v_h or be motivated by the faster field vehicle to accelerate to the speed v_{h+1}. The better the environmental conditions (α), and the more the free room felt by drivers in the ith cell (1 − ρ̃_i), the higher the probability of accelerating. Such actions are further modulated by the actual free room available in the destination cell I_{i+1}. This effect is accounted for by the flux limiter Φ_{i,i+1}, specifically by the fact that the candidate vehicle can be forced to stop (corresponding to the speed transition v_h → v₁ = 0) as the next cell fills. A technical modification is necessary at the boundary h = 1 of the speed lattice, for if the candidate vehicle is already still, then the options of either maintaining the speed or being forced to stop by the flux limiter coincide. We can summarize this case as follows:

$$(3.4) v_h < v_k \begin{cases} h = 1 & \begin{cases} A_{1k}^1(i) = 1 - \alpha(1 - \tilde{\rho}_i)\Phi_{i,i+1}, \\ A_{1k}^2(i) = \alpha(1 - \tilde{\rho}_i)\Phi_{i,i+1}, \\ A_{1k}^j(i) = 0 & \text{otherwise}; \end{cases} \\ h > 1 & \begin{cases} A_{hk}^1(i) = 1 - \Phi_{i,i+1}, \\ A_{hk}^h(i) = [1 - \alpha(1 - \tilde{\rho}_i)]\Phi_{i,i+1}, \\ A_{hk}^{h+1}(i) = \alpha(1 - \tilde{\rho}_i)\Phi_{i,i+1}, \\ A_{hk}^j(i) = 0 & \text{otherwise}. \end{cases}$$

• $v_h > v_k$. In this case, the candidate vehicle may either maintain its speed v_h , if, e.g., it can overtake the leading field vehicle, or be forced to queue by decelerating to the speed v_k of the latter. Specifically, from the expression of the corresponding transition probability reported in (3.5) below, it can be noted that this latter event is fostered by bad environmental conditions as well as by reduced free room felt by drivers in the *i*th cell. In particular, if the candidate vehicle interacts with a stationary field vehicle (k = 1), the transition probabilities are deduced by merging, as before, the first two cases of the expression valid for k > 1. To summarize,

$$(3.5) v_h > v_k \begin{cases} k = 1 & \begin{cases} A_{h1}^1(i) = 1 - \alpha(1 - \tilde{\rho}_i) \Phi_{i,i+1}, \\ A_{h1}^h(i) = \alpha(1 - \tilde{\rho}_i) \Phi_{i,i+1}, \\ A_{h1}^j(i) = 0 & \text{otherwise}; \end{cases} \\ k > 1 & \begin{cases} A_{hk}^1(i) = 1 - \Phi_{i,i+1}, \\ A_{hk}^k(i) = [1 - \alpha(1 - \tilde{\rho}_i)] \Phi_{i,i+1}, \\ A_{hk}^h(i) = \alpha(1 - \tilde{\rho}_i) \Phi_{i,i+1}, \\ A_{hk}^j(i) = 0 & \text{otherwise}. \end{cases}$$

• $v_h = v_k$. Finally, in this case the candidate vehicle can either be motivated to accelerate to the speed v_{h+1} or be induced to decelerate to the speed v_{h-1} by the field vehicle; alternatively, it may also decide to maintain its speed v_h . The corresponding transition probabilities are built from arguments analogous to those discussed in the previous points, thus invoking again the concepts of environmental conditions, free room felt by drivers in the *i*th cell, and modulation effect from the flux limiter. At the boundaries of the speed lattice some technical modifications are needed, for some output events may coincide while others may not apply (e.g., deceleration and acceleration are not possible in the lowest and highest speed classes, respectively). Summarizing,

$$\begin{cases} h = 1 & \begin{cases} A_{11}^{1}(i) = 1 - \alpha(1 - \tilde{\rho}_{i})\Phi_{i,i+1}, \\ A_{11}^{2}(i) = \alpha(1 - \tilde{\rho}_{i})\Phi_{i,i+1}, \\ A_{11}^{2}(i) = 0 \text{ otherwise;} \end{cases} \\ h = 2 & \begin{cases} A_{12}^{1}(i) = 1 - \Phi_{i,i+1} + (1 - \alpha)\tilde{\rho}_{i}\Phi_{i,i+1}, \\ A_{22}^{2}(i) = 1 - \Phi_{i,i+1} + (1 - \alpha)\tilde{\rho}_{i}\Phi_{i,i+1}, \\ A_{22}^{3}(i) = \alpha(1 - \tilde{\rho}_{i})\Phi_{i,i+1}, \\ A_{22}^{3}(i) = 0 \text{ otherwise;} \end{cases} \\ \begin{cases} A_{1h}^{1}(i) = 1 - \Phi_{i,i+1}, \\ A_{hh}^{1}(i) = (1 - \alpha)\tilde{\rho}_{i}\Phi_{i,i+1}, \\ A_{hh}^{1}(i) = (1 - \alpha)\tilde{\rho}_{i}\Phi_{i,i+1}, \\ A_{hh}^{1}(i) = \alpha(1 - \tilde{\rho}_{i})\Phi_{i,i+1}, \\ A_{hh}^{2}(i) = 0 \text{ otherwise;} \end{cases} \\ \begin{cases} A_{1n}^{1}(i) = 1 - \Phi_{i,i+1}, \\ A_{nn}^{2}(i) = (1 - \alpha)\tilde{\rho}_{i}\Phi_{i,i+1}, \end{cases} \end{cases}$$

- 4. Computational analysis. The validity of the model presented in the previous sections, and indirectly also that of the methodological approach which generated it, can be assessed through exploratory numerical simulations addressing typical phenomena of vehicular traffic.
- **4.1. The spatially homogeneous problem.** A first common case study refers to car flow in uniform space conditions, i.e., when one assumes that cars are homogeneously distributed along the road. As a matter of fact, this ideal scenario is the tacit assumption behind the empirical study of the so-called *fundamental diagrams* out of experimental data measured for real flows. Such diagrams relate the density of cars to either their average speed or their flux, thus providing synthetic insights into the gross

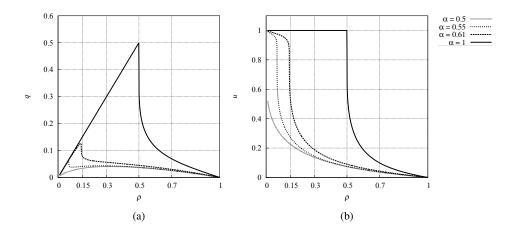


Fig. 4.1. Fundamental diagrams of (a) the average flux and (b) the average speed computed from the stationary solutions of the spatially homogeneous model (4.1) for various values of the parameter α indicated in the key.

phenomenology of car flow expected in stationary conditions. The representation at the mesoscopic (kinetic) scale is particularly suited to simulate, by mathematical models, this kind of information. In fact, the distribution function allows one to obtain the required statistics (cf. (2.3)) genuinely from the modeled microscopic interactions, rather than relying on them for modeling directly the average macroscopic dynamics. Specifically, considering the basic assumption of spatial homogeneity, the equations to be solved are

(4.1)
$$\frac{df_j}{dt} = \eta \left(\sum_{h,k=1}^n A_{hk}^j f_h f_k - f_j \rho \right), \qquad j = 1, \dots, n,$$

which are derived from (2.12) with $f_{ij} = f_{i+1,j}$ for i = 1, ..., m-1 and all j, along with the assumption that the interaction rate does not depend on h, k; cf. (3.2). In view of (3.1), the flux limiters also turn out to be independent of i, which makes the transport term on the left-hand side of (2.12) vanish. Ultimately, the index i of the spatial cell has been removed because it is unnecessary in the present context.

Owing to (2.9), equation (4.1) is such that the car density $\rho = \sum_{j=1}^{n} f_j$ is conserved in time:

$$\frac{d\rho}{dt} = \eta \left(\sum_{h,k=1}^{n} f_h f_k - \rho \sum_{j=1}^{n} f_j \right) = 0.$$

Therefore one can fix $\rho \in [0, 1]$ at the initial time and then, by means of (4.1), look for the asymptotic statistical distribution (if any) of microscopic speeds corresponding to the very same car density. Following this procedure, we obtained the fundamental diagrams of Figure 4.1 numerically for various values of the parameter α representing the environmental conditions in the table of games (cf. (3.4)–(3.6)). It is worth noticing that the other parameters η_0 (cf. (3.2)) and β (cf. (3.3)) instead play no role in shaping the asymptotic distributions, thus also the fundamental diagrams. Indeed, the former can be hidden in the time scale, whereas the latter drops because in spatially homogeneous conditions the fictitious car density coincides with the actual one.

The graphs in Figure 4.1 were computed by duly adapting to the present case the formulas reported in (2.3). They demonstrate that the model is able to catch qualitatively the well-known phase transition from free to congested traffic flow [15]. We recall that free flow is characterized by a positive gradient of the flux-density diagram and is separated from congested flow by the so-called density at capacity, say ρ_c , at which the highest traffic flux (road capacity) is attained. The trend of the average flux q, illustrated in Figure 4.1(a), is indeed basically linear and increasing for low density; then it becomes markedly nonlinear when ρ overcomes the critical value at which q is maximum, i.e., when the road progressively clogs up. Notice that for $\alpha = 1$, i.e., in the ideal situation of optimal environmental conditions, the average flux is exactly linear in the free flow phase, which persists until $\rho_c = 0.5$. In parallel, the trend of the average speed u, illustrated in Figure 4.1(b), is initially almost constant and close to the maximum possible value u = 1 (actually, for $\alpha = 1$ the model predicts that it is exactly constant to 1 until $\rho_c = 0.5$); then it drops steeply to zero when the density enters the congested flow range.

Remark 4.1. Fundamental diagrams obtained from the kinetic model can also be used for identifying admissible values of the parameter α , which otherwise might be hardly measured directly. To this purpose we relate α to the value of the density at capacity ρ_c , which can instead be detected quite reliably for a given road. For instance, experimental measurements on the Venice-Mestre highway in Italy [5] indicate that a realistic (dimensionless) range of the density at capacity there is $\rho_c \in [0, 0.15]$. According to the fundamental diagram in Figure 4.1(a), $\alpha = 0.55$ and $\alpha = 0.61$ are therefore two admissible values for that road, whereas $\alpha = 0.5$ and $\alpha = 1$ are not because they both imply a density at capacity greater than 0.15. In general, different roads will produce different empirical evaluations of ρ_c ; however, the tuning procedure of α based on the fundamental diagram remains conceptually the same.

Figure 4.1 indicates also that in no case is the (dimensionless) density at capacity expected to be greater than $\rho_c = 0.5$ (which is the value found for $\alpha = 1$). This can be explained by considering that for $\rho > 0.5$ cars are, on average, closer to one another than their characteristic size (i.e., the average distance between two successive cars is lower than the typical car size), which might prevent any of them from being unperturbed by cars ahead.

4.2. The spatially inhomogeneous problem. A second case study concerns the evolution in time of the car density along a road. Nonuniform space conditions are now explicitly considered, the interest being also in the transient trend of the traffic flow. A qualitative validation of the model can be performed upon assessing its ability to reproduce spatial flow patterns observed in real conditions. Here we focus specifically on the formation of queues, presenting two prototypical scenarios which may be of practical interest for traffic management. All relevant parameters used in the next simulations are summarized in Tables 4.1 and 4.2.

The first scenario refers to queue formation caused by a worsening of the driving conditions, for instance a narrowing of the roadway due to roadworks. Using the full equation (2.12), this can be simulated by assuming that the parameter α in the table of games varies from cell to cell, $\alpha = \alpha_i$, i = 1, ..., m. In particular, inspired by the reasoning proposed in Remark 4.1 for the Venice-Mestre highway, we consider a gradual decrease of α from 0.61 at the beginning of the road to 0.55 at its end. Figure 4.2 clearly shows the formation and backward propagation of a queue for such a variable α with respect to the basically unperturbed flow of vehicles produced when α is kept constant to the higher value ($\alpha = 0.61$) all along the road.

 $\begin{tabular}{ll} Table 4.1 \\ Model \ parameters \ for \ the \ roadworks \ problem. \end{tabular}$

Parameter	Value	Description
m	10	Number of space cells
n	6	Number of speed classes
η_0	1	Coefficient of the interaction rate; see (3.2)
α	0.61	Constant environmental conditions
$lpha_i$	$\begin{cases} 0.61 & \text{if } i \le 5\\ \frac{1}{10} \left(7 - \frac{3}{20}i\right) & \text{if } 6 \le i \le 9\\ 0.55 & \text{if } i = 10 \end{cases}$	Cell-to-cell variable environmental conditions
β	0	Coefficient of the fictitious density; see (3.3)

Table 4.2 Model parameters for the traffic light problem.

Parameter	Value	Description
m	10	Number of space cells
n	6	Number of speed classes
η_0	1	Coefficient of the interaction rate; see (3.2)
α	0.55	Environmental conditions
β	1	Coefficient of the fictitious density; see (3.3)

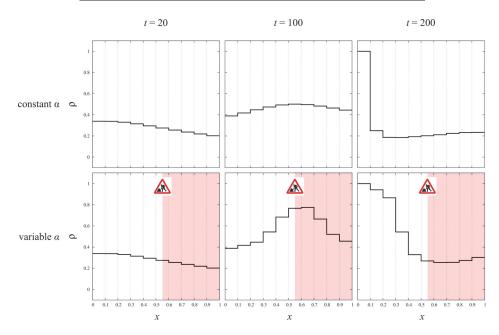


Fig. 4.2. Comparison, at the same computational times, of the traffic flow under constant and variable road conditions due to roadworks. The road is initially empty, while a constant density ρ_0 of incoming vehicles, uniformly distributed over the speed classes, is imposed at the left boundary at all times. The flux limiters $\Phi_{0,1}$ and $\Phi_{10,11}$ at the left and right boundaries are set to $\frac{1-\rho_1}{\rho_0}$ and to 1, respectively.

The second scenario refers instead to the alternate formation and depletion of a queue behind a traffic light, which, as Figure 4.3 demonstrates, are well reproduced by the model. In more detail, the action of the traffic light on the flow of vehicles is easily simulated, in our framework, by identifying the traffic light with one of the flux limiters ($\Phi_{5,6}$ in the case of Figure 4.3), to which a periodic time evolution is

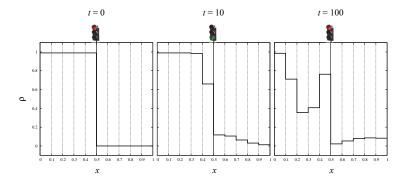


Fig. 4.3. Evolution of a queue during several red-green cycles of a traffic light placed at x = 0.5. No cars enter the domain from the left boundary for all t > 0. The initial condition is the preformed queue illustrated in the left snapshot. The traffic light period is of 20 computational time units.

assigned, mimicking the successive red-green cycles of the light. Specifically, during the red phase the selected limiter is set invariably to zero, since no car is allowed to flow beyond the traffic light regardless of the actual free space in the cell ahead. Conversely, during the green phase the selected limiter switches to the expression given by (3.1); hence the flow of cars is again possible according to the usual rules. The initial condition for this problem is, as shown by the first snapshot of Figure 4.3, a fully formed queue (i.e., $\rho_i = 1$) standing at a red light with empty road ahead. When, at t > 0, the light turns to green, cars in the cell $I_5 = [0.4, 0.5]$ can start moving only if they are able to see the empty space in front of them, namely if they react to a fictitious density $\tilde{\rho}_5$ lower than the real one $\rho_5 = 1$. Indeed, by carefully inspecting the table of games reported in (3.4)–(3.6), we find that standing vehicles (i.e., vehicles in the speed class $v_1 = 0$) in a full cell (i.e., one with $\rho_i = 1$) have no chance to shift to the velocity class $v_2 > 0$ if the transition probabilities depend on the actual cell density. Therefore, the concept of fictitious density as introduced in (3.3) is important for conferring realism to the model.

5. Basic qualitative analysis. The aim of this section is to study the wellposedness of the following class of initial/boundary-value problems, generated by the mathematical framework presented in the previous sections:

(5.1)
$$\begin{cases} \frac{df_{ij}}{dt} + v_j \left(\Phi_{i,i+1}[\mathbf{f}] f_{ij} - \Phi_{i-1,i}[\mathbf{f}] f_{i-1,j} \right) = \eta(i)[\mathbf{f}] \left(\sum_{h,k=1}^n A_{hk}^j(i)[\mathbf{f}] f_{ih} f_{ik} - f_{ij} \rho_i \right), \\ f_{ij}(0) = f_{ij}^0, \\ f_{0j}(t) = \bar{f}_j(t), \\ \Phi_{0,1} = \bar{\Phi}(t), \end{cases}$$

which we specifically consider in the mild formulation

(5.2)
$$\begin{cases} f_{ij}(t) = f_{ij}^{0} + \int_{0}^{t} \{v_{j} \left(\Phi_{i-1,i}[\mathbf{f}](s) f_{i-1,j}(s) - \Phi_{i,i+1}[\mathbf{f}](s) f_{ij}(s)\right) \\ + G_{ij}[\mathbf{f}, \mathbf{f}](s) - f_{ij}(s) L_{ij}[\mathbf{f}](s)\} ds, \\ f_{0j}(t) = \bar{f}_{j}(t), \\ \Phi_{0,1} = \bar{\Phi}(t), \end{cases}$$

where

(5.3)
$$G_{ij}[\mathbf{f}, \mathbf{f}] = \eta(i)[\mathbf{f}] \sum_{h,k=1}^{n} A_{hk}^{j}(i)[\mathbf{f}] f_{ih} f_{ik}, \qquad L_{ij}[\mathbf{f}] = \eta(i)[\mathbf{f}] \rho_{i}$$

are the so-called gain and loss operators, respectively (cf. (2.8)). Notice that, unlike the previous sections, here we have emphasized the dependence of the flux limiters, the interaction rate, and the table of games on the set of distribution functions $\mathbf{f} = \{f_{ij}\}_{i,j}$. This notation will indeed be useful in the next calculations.

We organize the qualitative analysis of problem (5.1) in two parts: first we prove the uniqueness of the solution and its continuous dependence on the initial and boundary data; next we state existence. As a preliminary to this, we introduce the functional framework in which we will consider problem (5.1) and the assumptions needed for achieving the proofs.

We denote by $X_T = C([0, T]; \mathbb{R}^{mn})$ the Banach space of vector-valued continuous functions $\mathbf{u} = \mathbf{u}(t) : [0, T] \to \mathbb{R}^{mn}$, $\mathbf{u}(t) = (u_{11}(t), \dots, u_{1n}(t), \dots, u_{m1}(t), \dots, u_{mn}(t))$, equipped with the uniform norm

$$\|\mathbf{u}\|_{\infty} := \sup_{t \in [0, T]} \|\mathbf{u}(t)\|_{1} = \sup_{t \in [0, T]} \sum_{i=1}^{m} \sum_{j=1}^{n} |u_{ij}(t)|,$$

where $\|\cdot\|_1$ is the classical 1-norm in \mathbb{R}^{mn} . We also introduce its subset

$$\mathcal{B} = \left\{ \mathbf{u} \in X_T : 0 \le u_{ij}(t) \le 1, \ \sum_{j=1}^n u_{ij}(t) \le 1 \ \forall i, j, \ \forall t \in [0, T] \right\},\,$$

which is readily seen to be closed. We point out that, by virtue of the equivalence of all norms in finite dimension, a different *p*-norm may be considered in \mathbb{R}^{mn} . However, the 1-norm takes advantage more directly of the interpretation in terms of number density: if $\mathbf{u} \in \mathcal{B}$, then $\|\mathbf{u}(t)\|_1$ is the density of the distribution \mathbf{u} at time t (cf. (2.1)).

A solution to our initial/boundary-value problem is defined as follows.

DEFINITION 5.1. A function $\mathbf{f} = \mathbf{f}(t) : [0, T] \to \mathbb{R}^{mn}$ is said to be a mild solution to problem (5.1) if $\mathbf{f} \in \mathcal{B}$ and \mathbf{f} satisfies (5.2).

Notice that $\mathbf{f} \in \mathcal{B}$ guarantees that \mathbf{f} fulfills the basic requirements of a physically acceptable distribution function: nonnegativity and boundedness in each space/speed class in the whole time interval of existence.

The following box summarizes the hypotheses, to be possibly regarded as modeling guidelines, under which we will achieve the proofs of the announced results.

Assumptions for problem (5.1).

- 1. The initial and boundary data f_{ij}^0 , \bar{f}_j belong to \mathcal{B} .
- 2. The flux limiter $\Phi_{i,i+1}[\mathbf{f}]$ is defined by a function $\Phi: [0, 1] \times [0, 1] \to [0, 1]$ acting on the densities of the two adjacent cells I_1 , I_{i+1} such that $\Phi_{i,i+1}[\mathbf{f}] = \Phi(\rho_i, \rho_{i+1})$. The function Φ has the following properties:
 - (i) $0 \le \Phi(u, v) \le 1 \ \forall u, v \in [0, 1];$
 - (ii) $\Phi(u, v)u < 1 v \ \forall (u, v) \in [0, 1] \times [0, 1]$ such that u + v > 1;
 - (iii) $|\Phi(u_2, v_2) \Phi(u_1, v_1)| \le \operatorname{Lip}(\Phi)(|u_2 u_1| + |v_2 v_1|), \forall (u_1, v_1), (u_2, v_2) \in [0, 1] \times [0, 1].$

- 3. The boundary datum $\bar{\Phi}$ is as in Assumption 2; i.e., $\bar{\Phi} = \Phi(\bar{\rho}, \rho_1)$, where $\bar{\rho} = \sum_{j=1}^{n} \bar{f}_j$ is the prescribed density of incoming cars.
- 4. The interaction rate $\eta(i)$ is
 - (a) nonnegative and uniformly bounded; i.e., there exists a constant $\bar{\eta} > 0$ such that

$$0 \le \eta(i)[\mathbf{u}] \le \bar{\eta}, \quad \forall i = 1, \dots m, \ \forall \, \mathbf{u} \in \mathcal{B};$$

(b) Lipschitz continuous with respect to the distribution in \mathcal{B} ; i.e.,

$$|\eta(i)[\mathbf{u}](t) - \eta(i)[\mathbf{v}](t)| \le \operatorname{Lip}(\eta(i)) ||\mathbf{u}(t) - \mathbf{v}(t)||_1$$

 $\forall t \in [0, T], \ \forall \mathbf{u}, \mathbf{v} \in \mathcal{B}.$

5. The elements $A_{hk}^{j}(i)$ of the table of games satisfy (2.9) and the following Lipschitz continuity condition:

$$\left| A_{hk}^{j}(i)[\mathbf{u}](t) - A_{hk}^{j}(i)[\mathbf{v}](t) \right| \leq \operatorname{Lip}(A_{hk}^{j}(i)) \|\mathbf{u}(t) - \mathbf{v}(t)\|_{1}$$
$$\forall t \in [0, T], \ \forall \mathbf{u}, \mathbf{v} \in \mathcal{B}.$$

We henceforth assume that Assumptions 1–4(a) are always satisfied; indeed they are needed in all of the next results. In contrast, we will explicitly mention Assumptions 4(b) and 5 only when they will actually be used.

5.1. Uniqueness and continuous dependence on the data. The following theorem states that the mild solution to problem (5.1) is unique in \mathcal{B} and that it depends continuously on the initial and boundary data.

THEOREM 5.2 (uniqueness and continuous dependence). Under Assumptions 4(b) and 5, let (\mathbf{f}^0 , $\bar{\mathbf{f}}$, $\bar{\Phi}^f$), (\mathbf{g}^0 , $\bar{\mathbf{g}}$, $\bar{\Phi}^g$) be two sets of initial and boundary data and \mathbf{f} , $\mathbf{g} \in \mathcal{B}$ two corresponding mild solutions to problem (5.1). Then there exists $\mathcal{C} > 0$ such that

$$\|\mathbf{f} - \mathbf{g}\|_{\infty} \le C \left[\|\mathbf{f}^0 - \mathbf{g}^0\|_1 + \int_0^T \left(|\bar{\Phi}^f(t) - \bar{\Phi}^g(t)| + \|\bar{\mathbf{f}}(t) - \bar{\mathbf{g}}(t)\|_1 \right) dt \right].$$

In particular, there is at most one solution corresponding to a given set of initial and boundary data.

Proof. Subtracting term by term the mild equations satisfied by \mathbf{f} , \mathbf{g} (cf. (5.2)) and then summing over i = 1, ..., m and j = 1, ..., n gives

$$\|\mathbf{f}(t) - \mathbf{g}(t)\|_{1} \leq \|\mathbf{f}^{0} - \mathbf{g}^{0}\|_{1}$$

$$+ \sum_{i=1}^{m} \sum_{j=1}^{n} \left[\int_{0}^{t} |v_{j}(\Phi_{i-1,i}[\mathbf{f}](s)f_{i-1,j}(s) - \Phi_{i-1,i}[\mathbf{g}](s)g_{i-1,j}(s) + \Phi_{i,i+1}[\mathbf{g}](s)g_{ij}(s) - \Phi_{i,i+1}[\mathbf{f}](s)f_{ij}(s)) | ds + \int_{0}^{t} (|G_{ij}[\mathbf{f}, \mathbf{f}](s) - G_{ij}[\mathbf{g}, \mathbf{g}](s)| + |g_{ij}(s)L_{ij}[\mathbf{g}](s) - f_{ij}(s)L_{ij}[\mathbf{f}](s)| ds \right].$$

By exploiting the Lipschitz continuity of the various terms and extracting from the sums the boundary terms corresponding to i = 1, we then discover

(5.4)
$$\sum_{i=1}^{m} \sum_{j=1}^{n} |\Phi_{i-1,i}[\mathbf{f}](s) f_{i-1,j}(s) - \Phi_{i-1,i}[\mathbf{g}](s) g_{i-1,j}(s) + \Phi_{i,i+1}[\mathbf{g}](s) g_{ij}(s) - \Phi_{i,i+1}[\mathbf{f}](s) f_{ij}(s)|$$

$$\leq 2(2 \operatorname{Lip}(\Phi) + 1) ||\mathbf{f}(s) - \mathbf{g}(s)||_{1} + |\bar{\Phi}^{f}(s) - \bar{\Phi}^{g}(s)| + ||\bar{\mathbf{f}}(s) - \bar{\mathbf{g}}(s)||_{1},$$

(5.5)
$$\sum_{i=1}^{m} \sum_{j=1}^{n} |G_{ij}[\mathbf{f}, \mathbf{f}](s) - G_{ij}[\mathbf{g}, \mathbf{g}](s)|$$

$$\leq \left[2\bar{\eta}n + \sum_{i=1}^{m} \left(n^{2} \operatorname{Lip}(\eta(i)) + \bar{\eta} \sum_{h, k, j=1}^{n} \operatorname{Lip}(A_{hk}^{j}(i)) \right) \right] \|\mathbf{f}(s) - \mathbf{g}(s)\|_{1},$$

$$\sum_{i=1}^{m} \sum_{j=1}^{n} |g_{ij}(s)L_{ij}[\mathbf{g}](s) - f_{ij}(s)L_{ij}[\mathbf{f}](s)| \le \left(2\bar{\eta}n + n^2 \sum_{i=1}^{m} \operatorname{Lip}(\eta(i))\right) \|\mathbf{f}(s) - \mathbf{g}(s)\|_1,$$

whence

$$\begin{aligned} \|\mathbf{f}(t) - \mathbf{g}(t)\|_{1} &\leq \|\mathbf{f}^{0} - \mathbf{g}^{0}\|_{1} \\ &+ \mathcal{C} \int_{0}^{t} \left(\|\mathbf{f}(s) - \mathbf{g}(s)\|_{1} + \left| \bar{\Phi}^{f}(s) - \bar{\Phi}^{g}(s) \right| + \|\bar{\mathbf{f}}(s) - \bar{\mathbf{g}}(s)\|_{1} \right) ds, \end{aligned}$$

where C is a suitable constant deduced from all those appearing in (5.4)–(5.6). Gronwall's inequality now implies

$$\|\mathbf{f}(t) - \mathbf{g}(t)\|_{1} \le e^{\mathcal{C}t} \left[\|\mathbf{f}^{0} - \mathbf{g}^{0}\|_{1} + \int_{0}^{t} \left(|\bar{\Phi}^{f}(s) - \bar{\Phi}^{g}(s)| + \|\bar{\mathbf{f}}(s) - \bar{\mathbf{g}}(s)\|_{1} \right) ds \right],$$

whence, taking the supremum over $t \in [0, T]$ on both sides, we obtain the continuous dependence estimate. Uniqueness follows straightforwardly from it by letting $\mathbf{f}^0 = \mathbf{g}^0$, $\bar{\mathbf{f}} = \bar{\mathbf{g}}$, and $\bar{\Phi}^f = \bar{\Phi}^g$.

5.2. Existence. In this section we prove the existence of the mild solution to problem (5.1). The idea of the proof is to consider at first the model at discrete time instants $t_n^k = n\Delta t_k$, where the index n labels the discrete time and the index k is a mesh parameter denoting the level of refinement of the time grid. In particular, the time step Δt_k is chosen such that it tends to zero for $k \to \infty$. After establishing some technical properties of the iterates of such a discrete-in-time model (cf. Lemma 5.3), we pass from discrete to continuous time by interpolating the iterates, and we prove that in the limit $k \to \infty$ this interpolation converges to a distribution function $f \in \mathcal{B}$ (cf. Lemma 5.4). Finally, we show that f is precisely a mild solution to problem (5.1) (cf. Theorem 5.5).

Step 1. We consider problem (5.1) at discrete time instants (cf. (2.5)):

$$f_{ij}^{\mathsf{n}+1,\mathsf{k}} = f_{ij}^{\mathsf{n},\mathsf{k}} - \Delta t_{\mathsf{k}} v_{j} (\Phi_{i,i+1}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] f_{ij}^{\mathsf{n},\mathsf{k}} - \Phi_{i-1,i}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] f_{i-1,j}^{\mathsf{n},\mathsf{k}}) + \Delta t_{\mathsf{k}} (G_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}, \mathbf{f}^{\mathsf{n},\mathsf{k}}] - f_{ij}^{\mathsf{n},\mathsf{k}} L_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}]).$$
(5.7)

The following lemma gives some properties of the iterates $f_{ij}^{n,k}$.

LEMMA 5.3. If Δt_k is sufficiently small, then $\mathbf{f}^{\mathsf{n},\mathsf{k}} \in \mathcal{B}$ for all n , $\mathsf{k} \geq 0$.

Proof. The proof is by induction on n, assuming that $\mathbf{f}^{n,k} \in \mathcal{B}$ and that the initial and boundary data have been chosen consistent with the properties of the distribution functions in \mathcal{B} .

(i) Nonnegativity of the iterates. Taking into account the nonnegativity of the terms containing $\Phi_{i-1,i}[\mathbf{f}^{\mathsf{n},\mathsf{k}}]$ and $G_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}},\mathbf{f}^{\mathsf{n},\mathsf{k}}]$, we have

$$f_{ij}^{\mathsf{n}+1,\mathsf{k}} \geq f_{ij}^{\mathsf{n},\mathsf{k}}[1 - \Delta t_{\mathsf{k}}(v_j \Phi_{i,i+1}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] + L_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}])] \geq f_{ij}^{\mathsf{n},\mathsf{k}}[1 - \Delta t_{\mathsf{k}}(1 + L_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}])],$$

with v_j , $\Phi_{i,i+1}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] \leq 1$. Moreover, $L_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] \leq \bar{\eta}\rho_i^{\mathsf{n},\mathsf{k}} \leq \bar{\eta}$; hence if $\Delta t_\mathsf{k} < 1/(1+\bar{\eta})$, we deduce

$$f_{ij}^{\mathsf{n}+1,\mathsf{k}} \ge f_{ij}^{\mathsf{n},\mathsf{k}} [1 - \Delta t_{\mathsf{k}} (1+\bar{\eta})] \ge 0.$$

(ii) Boundedness of the iterates. Considering that $v_i \leq 1$, we have

$$f_{ij}^{\mathsf{n}+1,\mathsf{k}} \leq f_{ij}^{\mathsf{n},\mathsf{k}} + \Delta t_\mathsf{k}(\Phi_{i-1,i}[\mathbf{f}^{\mathsf{n},\mathsf{k}}]f_{i-1,j}^{\mathsf{n},\mathsf{k}} + G_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}},\,\mathbf{f}^{\mathsf{n},\mathsf{k}}] - f_{ij}^{\mathsf{n},\mathsf{k}}L_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}]).$$

But

$$G_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}},\,\mathbf{f}^{\mathsf{n},\mathsf{k}}] - f_{ij}^{\mathsf{n},\mathsf{k}} L_{ij}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] \le \bar{\eta} \left[1 - \left(f_{ij}^{\mathsf{n},\mathsf{k}} \right)^2 \right],$$

and thus we conclude

(5.8)
$$f_{ij}^{\mathsf{n}+1,\mathsf{k}} \le f_{ij}^{\mathsf{n},\mathsf{k}} + \Delta t_{\mathsf{k}} \left\{ \Phi_{i-1,i}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] f_{i-1,j}^{\mathsf{n},\mathsf{k}} + \bar{\eta} \left[1 - \left(f_{ij}^{\mathsf{n},\mathsf{k}} \right)^2 \right] \right\}.$$

Now, if $f_{ij}^{\mathsf{n},\mathsf{k}} = 1$, then from Assumption 2(ii) it follows that $f_{ij}^{\mathsf{n}+1,\mathsf{k}} \leq 1$, as desired. Conversely, if $f_{ij}^{\mathsf{n},\mathsf{k}} < 1$, then according to Assumption 2(i)–(ii) we have

(5.9)
$$\Phi_{i-1,i}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] f_{i-1,j}^{\mathsf{n},\mathsf{k}} \leq \Phi_{i-1,i}[\mathbf{f}^{\mathsf{n},\mathsf{k}}] \rho_{i-1}^{\mathsf{n},\mathsf{k}} \\ \leq \left\{ \begin{array}{ll} 1 - \rho_i^{\mathsf{n},\mathsf{k}} & \text{if } \rho_i^{\mathsf{n},\mathsf{k}} + \rho_{i-1}^{\mathsf{n},\mathsf{k}} > 1\\ \rho_{i-1}^{\mathsf{n},\mathsf{k}} \leq 1 - \rho_i^{\mathsf{n},\mathsf{k}} & \text{if } \rho_i^{\mathsf{n},\mathsf{k}} + \rho_{i-1}^{\mathsf{n},\mathsf{k}} \leq 1 \end{array} \right\} \leq 1 - \rho_i^{\mathsf{n},\mathsf{k}}.$$

Hence we can continue the estimate (5.8) as

$$f_{ij}^{\mathsf{n}+1,\mathsf{k}} \leq f_{ij}^{\mathsf{n},\mathsf{k}} + \Delta t_{\mathsf{k}} \left\{ 1 - \rho_i^{\mathsf{n},\mathsf{k}} + \bar{\eta} \left[1 - \left(f_{ij}^{\mathsf{n},\mathsf{k}} \right)^2 \right] \right\} \leq f_{ij}^{\mathsf{n},\mathsf{k}} + \Delta t_{\mathsf{k}} (1 + 2\bar{\eta}) (1 - f_{ij}^{\mathsf{n},\mathsf{k}}),$$

whence we see again that $f_{ij}^{\mathsf{n}+1,\mathsf{k}} \leq 1$ if $\Delta t_{\mathsf{k}} < 1/(1+2\bar{\eta})$.

(iii) Boundedness of the sum of the iterates. In order to prove that $\sum_{j=1}^{n} f_{ij}^{n,k} \leq 1$, we note that from (2.9) we have $\sum_{j=1}^{n} (G_{ij}[\mathbf{f}^{n,k}, \mathbf{f}^{n,k}] - f_{ij}^{n,k} L_{ij}[\mathbf{f}^{n,k}]) = 0$. Consequently, by (5.9) we obtain the thesis, since

$$\sum_{i=1}^{n} f_{ij}^{n+1,k} \le \rho_i^{n,k} + \Delta t_k \Phi_{i-1,i}[\mathbf{f}^{n,k}] \rho_{i-1}^{n,k} \le \rho_i^{n,k} + \Delta t_k (1 - \rho_i^{n,k}) \le 1,$$

provided that $\Delta t_{\mathsf{k}} \leq 1$.

Finally, we note that all properties proved so far hold simultaneously if $\Delta t_{\mathsf{k}} < 1/(1+2\bar{\eta})$. \square

Step 2. Now we pass by interpolation from discrete to continuous time. To this end, we introduce the following function $\hat{\mathbf{f}} = \hat{\mathbf{f}}^{\mathsf{k}}(t) : [0, T] \to \mathbb{R}^{mn}$ interpolating piecewise linearly the iterates $f_{ij}^{\mathsf{n},\mathsf{k}}$:

(5.10)
$$\hat{f}_{ij}^{\mathbf{k}}(t) = \sum_{\mathbf{n}=1}^{N_{\mathbf{k}}} \left[\left(1 - \frac{t - t_{\mathbf{n}-1}^{\mathbf{k}}}{\Delta t_{\mathbf{k}}} \right) f_{ij}^{\mathbf{n}-1,\mathbf{k}} + \frac{t - t_{\mathbf{n}-1}^{\mathbf{k}}}{\Delta t_{\mathbf{k}}} f_{ij}^{\mathbf{n},\mathbf{k}} \right] \mathbb{1}_{[t_{\mathbf{n}-1}^{\mathbf{k}}, t_{\mathbf{n}}^{\mathbf{k}}]}(t).$$

Here N_k is the total number of time steps in the kth mesh. We assume that N_k and Δt_k are chosen in such a way that $N_k \Delta t_k = T$ independently of the refinement parameter k.

For the function $\hat{\mathbf{f}}^{k}$ we can prove the following result.

Lemma 5.4. There exists $\mathbf{f} \in \mathcal{B}$ such that

$$\lim_{k \to \infty} \|\hat{\mathbf{f}}^k - \mathbf{f}\|_{\infty} = 0.$$

Proof. We first claim that $\hat{\mathbf{f}}^{\mathsf{k}} \in \mathcal{B}$. In fact, for all $t \in [0, T]$ there exists $\bar{\mathsf{n}} \leq N_{\mathsf{k}}$ such that $t \in [t_{\bar{\mathsf{n}}-1}^{\mathsf{k}}, t_{\bar{\mathsf{n}}}^{\mathsf{k}}]$, and thus

$$\hat{f}_{ij}^{\mathbf{k}}(t) = \left(1 - \frac{t - t_{\bar{\mathbf{n}} - 1}^{\mathbf{k}}}{\Delta t_{\mathbf{k}}}\right) f_{ij}^{\bar{\mathbf{n}} - 1, \mathbf{k}} + \frac{t - t_{\bar{\mathbf{n}} - 1}^{\mathbf{k}}}{\Delta t_{\mathbf{k}}} f_{ij}^{\bar{\mathbf{n}}, \mathbf{k}},$$

and according to Lemma 5.3 this value complies with all constraints that the functions in \mathcal{B} are required to satisfy pointwise in time. Thus the claim follows from the arbitrariness of t.

Next, our goal is to show that the sequence $\{\hat{\mathbf{f}}^k\}_{k\geq 0}$ is relatively compact in X_T . Indeed, in such a case we can conclude that (up to subsequences) it converges uniformly to some $\mathbf{f} \in X_T$. Moreover, since $\{\hat{\mathbf{f}}^k\}_{k\geq 0} \subset \mathcal{B}$ and \mathcal{B} is closed, we also have that $\mathbf{f} \in \mathcal{B}$. We prove the relative compactness by means of the Ascoli–Arzelà compactness criterion:

(i) $\{\hat{\mathbf{f}}^k\}_{k\geq 0}$ is uniformly bounded. Indeed, since $\hat{\mathbf{f}}^k\in\mathcal{B}$ we can write

$$\sup_{{\bf k} \geq 0} \|\hat{\bf f}^{\bf k}\|_{\infty} = \sup_{{\bf k} \geq 0} \sup_{t \in [0,\,T]} \sum_{i=1}^m \sum_{j=1}^n \left| \hat{f}^{\bf k}_{ij}(t) \right| \leq m < \infty.$$

(ii) $\{\hat{\mathbf{f}}^k\}_{k\geq 0}$ is equicontinuous. In fact, let $t_1\in[t_{\bar{m}-1}^k,\,t_{\bar{m}}^k]$ and $t_2\in[t_{\bar{n}-1}^k,\,t_{\bar{n}}^k]$ with $m\leq n$ (i.e., $t_1\leq t_2$); then we can write

$$\hat{f}_{ij}^{\mathbf{k}}(t_{2}) - \hat{f}_{ij}^{\mathbf{k}}(t_{1})
= \hat{f}_{ij}^{\mathbf{k}}(t_{2}) + \sum_{\ell=\bar{\mathbf{m}}}^{\bar{\mathbf{n}}-1} \hat{f}_{ij}^{\mathbf{k}}(t_{\ell}^{\mathbf{k}}) - \sum_{\ell=\bar{\mathbf{m}}}^{\bar{\mathbf{n}}-1} \hat{f}_{ij}^{\mathbf{k}}(t_{\ell}^{\mathbf{k}}) - \hat{f}_{ij}^{\mathbf{k}}(t_{1})
= [\hat{f}_{ij}^{\mathbf{k}}(t_{2}) - \hat{f}_{ij}^{\mathbf{k}}(t_{\bar{\mathbf{n}}-1}^{\mathbf{k}})] + [\hat{f}_{ij}^{\mathbf{k}}(t_{\bar{\mathbf{m}}}^{\mathbf{k}}) - \hat{f}_{ij}^{\mathbf{k}}(t_{1})] + \sum_{\ell=\bar{\mathbf{m}}}^{\bar{\mathbf{n}}-2} (f_{ij}^{\ell+1,\mathbf{k}} - f_{ij}^{\ell,\mathbf{k}})
=: A_{1} + A_{2} + A_{3}.$$

We begin by estimating the term A_3 . By (5.7) and applying the same argu-

ments as in the proof of Lemma 5.3, we have

$$\begin{split} \left| f_{ij}^{\ell+1,\mathsf{k}} - f_{ij}^{\ell,\mathsf{k}} \right| &= \Delta t_{\mathsf{k}} \left| v_{j}(\Phi_{i-1,i}[\mathbf{f}^{\ell,\mathsf{k}}] f_{i-1,j}^{\ell,\mathsf{k}} - \Phi_{i,i+1}[\mathbf{f}^{\ell,\mathsf{k}}] f_{i,j}^{\ell,\mathsf{k}}) \right. \\ &+ G_{ij}[\mathbf{f}^{\ell,\mathsf{k}},\, \mathbf{f}^{\ell,\mathsf{k}}] - f_{ij}^{\ell,\mathsf{k}} L_{ij}[\mathbf{f}^{\ell,\mathsf{k}}] \right| \\ &\leq 2(1+\bar{\eta}) \Delta t_{\mathsf{k}}; \end{split}$$

thus

(5.12)

$$|A_3| \le \sum_{\ell=\bar{\mathbf{m}}}^{\bar{\mathbf{n}}-2} \left| f_{ij}^{\ell+1,\mathsf{k}} - f_{ij}^{\ell,\mathsf{k}} \right| \le 2(\bar{\mathbf{n}} - 1 - \bar{\mathbf{m}})(1 + \bar{\eta}) \Delta t_{\mathsf{k}} = 2(1 + \bar{\eta})(t_{\bar{\mathbf{n}}-1}^{\mathsf{k}} - t_{\bar{\mathbf{m}}}^{\mathsf{k}}).$$

As for the term A_2 , we observe that, taking into account (5.10) and (5.12), we can write

$$|A_{2}| = \left| \hat{f}_{ij}^{\mathsf{k}}(t_{2}) - \hat{f}_{ij}^{\mathsf{k}}(t_{\bar{\mathsf{n}}-1}^{\mathsf{k}}) \right| = \left| \hat{f}_{ij}^{\mathsf{k}}(t_{2}) - f_{ij}^{\bar{\mathsf{n}}-1,\mathsf{k}} \right|$$

$$= \frac{t_{2} - t_{\bar{\mathsf{n}}-1}^{\mathsf{k}}}{\Delta t_{\mathsf{k}}} \left| f_{ij}^{\bar{\mathsf{n}},\mathsf{k}} - f_{ij}^{\bar{\mathsf{n}}-1,\mathsf{k}} \right| \le 2(1 + \bar{\eta})(t_{2} - t_{\bar{\mathsf{n}}-1}^{\mathsf{k}}).$$

Likewise,

$$|A_{1}| = \left| \hat{f}_{ij}^{\mathsf{k}}(t_{\bar{\mathsf{m}}}^{\mathsf{k}}) - \hat{f}_{ij}^{\mathsf{k}}(t_{1}) \right| = \left| f_{ij}^{\bar{\mathsf{m}},\mathsf{k}} - \hat{f}_{ij}^{\mathsf{k}}(t_{1}) \right|$$

$$= \left(1 - \frac{t_{1} - t_{\bar{\mathsf{m}}-1}^{\mathsf{k}}}{\Delta t_{\mathsf{k}}} \right) \left| f_{ij}^{\bar{\mathsf{m}},\mathsf{k}} - f_{ij}^{\bar{\mathsf{m}}-1,\mathsf{k}} \right| \leq 2(1 + \bar{\eta})(t_{\bar{\mathsf{m}}}^{\mathsf{k}} - t_{1}).$$

Finally, by plugging (5.12)–(5.14) into (5.11), we get

$$\left|\hat{f}_{ij}^{\mathbf{k}}(t_2) - \hat{f}_{ij}^{\mathbf{k}}(t_1)\right| \leq 2(1+\bar{\eta})(t_{\bar{\mathbf{m}}}^{\mathbf{k}} - t_1 + t_2 - t_{\bar{\mathbf{n}}-1}^{\mathbf{k}} + t_{\bar{\mathbf{n}}-1}^{\mathbf{k}} - t_{\bar{\mathbf{m}}}^{\mathbf{k}}) = 2(1+\bar{\eta})(t_2 - t_1),$$

whence

$$\|\hat{\mathbf{f}}^{\mathsf{k}}(t_2) - \hat{\mathbf{f}}^{\mathsf{k}}(t_1)\|_1 \le 2mn(1+\bar{\eta})(t_2-t_1).$$

This implies the desired equicontinuity of the family $\{\hat{\mathbf{f}}^k\}_{k\geq 0}$. In fact, for every $\epsilon>0$ it is sufficient to take $|t_2-t_1|<\frac{\epsilon}{2mn(1+\bar{\eta})}$, independently of k, in order to have $\|\hat{\mathbf{f}}^k(t_2)-\hat{\mathbf{f}}^k(t_1)\|_1<\epsilon$.

Step 3. Now we are in a position to state that a mild solution to problem (5.1) does indeed exist.

THEOREM 5.5 (existence). Let Assumptions 4(b) and 5 hold. Then $\mathbf{f} \in \mathcal{B}$ found in Lemma 5.4 is a mild solution to problem (5.1) in the sense of Definition 5.1.

Proof. Owing to the convergence result stated in Lemma 5.4, we can guess that, for a fixed k, the function $\hat{\mathbf{f}}^k$ defined componentwise by (5.10) is a sort of approximation of a possible mild solution to problem (5.1). Therefore, if we plug \hat{f}_{ij}^k into (5.2), we expect a reminder e_{ij}^k to appear:

$$\hat{f}_{ij}^{k}(t) - f_{ij}^{0} + \int_{0}^{t} \left\{ v_{j}(\Phi_{i,i+1}[\hat{\mathbf{f}}^{k}](s)\hat{f}_{ij}^{k}(s) - \Phi_{i-1,i}[\hat{\mathbf{f}}^{k}](s)\hat{f}_{i-1,j}^{k}(s)) - G_{ij}[\hat{\mathbf{f}}^{k}, \hat{\mathbf{f}}^{k}](s) + \hat{f}_{ij}^{k}(s)L_{ij}[\hat{\mathbf{f}}^{k}](s) \right\} ds = \int_{0}^{t} e_{ij}^{k}(s) ds.$$

At this point we would like to take the limit $k \to \infty$ using Lemma 5.4. This will tell us that the distribution function f found in the latter is indeed a mild solution to problem (5.1), provided that the following hold:

- (i) the left-hand side of (5.15) tends to the corresponding expression evaluated for $\hat{\mathbf{f}}$ rather than for $\hat{\mathbf{f}}^{\mathbf{k}}$;
- (ii) the right-hand side of (5.15) tends to zero.

In order to prove (i) it is sufficient to observe that all terms appearing in the integral are bounded by an integrable constant; thus by dominated convergence it is possible to commute the limit in k with the integral in t. Using Lemma 5.4 and taking into account (5.4)–(5.6) with $\mathbf{g} = \hat{\mathbf{f}}^k$ then yields

$$\lim_{k \to \infty} \left| \hat{f}_{ij}^{k}(t) - f_{ij}^{0} + \int_{0}^{t} \left\{ v_{j}(\Phi_{i,i+1}[\hat{\mathbf{f}}^{k}](s)\hat{f}_{ij}^{k}(s) - \Phi_{i-1,i}[\hat{\mathbf{f}}^{k}](s)\hat{f}_{i-1,j}^{k}(s) \right\} \right.$$

$$\left. - G_{ij}[\hat{\mathbf{f}}^{k}, \hat{\mathbf{f}}^{k}](s) + \hat{f}_{ij}^{k}(s)L_{ij}[\hat{\mathbf{f}}^{k}](s) \right\} ds \right|$$

$$= \left| f_{ij}(t) - f_{ij}^{0} + \int_{0}^{t} \left\{ v_{j}(\Phi_{i,i+1}[\mathbf{f}](s)f_{ij}(s) - \Phi_{i-1,i}[\mathbf{f}](s)f_{i-1,j}(s)) - G_{ij}[\mathbf{f}, \mathbf{f}](s) + f_{ij}(s)L_{ij}[\mathbf{f}](s) \right\} ds \right|,$$

as desired.

It remains to prove (ii), for which we need to find an expression of the term e_{ij}^{k} . Using (5.7), (5.10), we write

$$\bullet \frac{d\hat{f}_{ij}^{k}}{dt} = \sum_{n=1}^{N_{k}} \left\{ v_{j} \left(\Phi_{i-1,i}[\mathbf{f}^{n-1,k}] f_{i-1,j}^{n-1,k} - \Phi_{i,i+1}[\mathbf{f}^{n-1,k}] f_{ij}^{n-1,k} \right) \right. \\ \left. + G_{ij}[\mathbf{f}^{n-1,k}, \mathbf{f}^{n-1,k}] - f_{ij}^{n-1,k} L_{ij}[\mathbf{f}^{n-1,k}] \right\};$$

$$\bullet v_{j} \left(\Phi_{i,i+1}[\hat{\mathbf{f}}^{k}] \hat{f}_{ij}^{k} - \Phi_{i-1,i}[\hat{\mathbf{f}}^{k}] \hat{f}_{i-1,j}^{k} \right) \\ = v_{j} \sum_{n=1}^{N_{k}} \left\{ \Phi_{i,i+1}[\hat{\mathbf{f}}^{k}] f_{ij}^{n-1,k} - \Phi_{i-1,i}[\mathbf{f}^{n,k}] f_{i-1,j}^{n-1,k} \right. \\ \left. + \frac{t - t_{n-1}^{k}}{\Delta t_{k}} \Phi_{i,i+1}[\hat{\mathbf{f}}^{k}] \left(f_{ij}^{n,k} - f_{ij}^{n-1,k} \right) \right. \\ \left. + \frac{t - t_{n-1}^{k}}{\Delta t_{k}} \Phi_{i-1,i}[\hat{\mathbf{f}}^{k}] \left(f_{ij}^{n-1,k} - f_{i-1,j}^{n-1,k} \right) \right\} \mathbb{1}_{[t_{n-1}^{k}, t_{n}^{k}]};$$

$$\bullet G_{ij}[\hat{\mathbf{f}}^{k}, \hat{\mathbf{f}}^{k}] = \sum_{n,k=1}^{n} \sum_{n=1}^{N_{k}} \eta[\hat{\mathbf{f}}^{k}] A_{hk}^{j}[\hat{\mathbf{f}}^{k}] \\ \cdot \left\{ f_{ih}^{n-1,k} f_{ik}^{n-1,k} + \frac{t - t_{n-1}^{k}}{\Delta t_{k}} f_{ih}^{n-1,k} \left(f_{ih}^{n,k} - f_{ih}^{n-1,k} \right) \right. \\ \left. + \frac{t - t_{n-1}^{k}}{\Delta t_{k}} \hat{f}_{ik}^{n-1,k} \left(f_{ih}^{n,k} - f_{ih}^{n-1,k} \right) + \left(\frac{t - t_{n-1}^{k}}{\Delta t_{k}} \right)^{2} \left(f_{ih}^{n,k} - f_{ih}^{n-1,k} \right) \left(f_{ik}^{n,k} - f_{ik}^{n-1,k} \right) \right\} \mathbb{1}_{[t_{n-1}^{k}, t_{n}^{k}]};$$

$$\hat{f}_{ij}^{k} L_{ij}[\hat{\mathbf{f}}^{k}] = \sum_{k=1}^{n} \sum_{n=1}^{N_{k}} \eta[\hat{\mathbf{f}}^{k}]$$

$$\cdot \left\{ f_{ij}^{n-1,k} f_{ik}^{n-1,k} + \frac{t - t_{n-1}^{k}}{\Delta t_{k}} f_{ij}^{n-1,k} \left(f_{ik}^{n,k} - f_{ik}^{n-1,k} \right) \right.$$

$$+ \frac{t - t_{n-1}^{k}}{\Delta t_{k}} f_{ik}^{n-1,k} \left(f_{ij}^{n,k} - f_{ij}^{n-1,k} \right)$$

$$+ \left(\frac{t - t_{n-1}^{k}}{\Delta t_{k}} \right)^{2} \left(f_{ij}^{n,k} - f_{ij}^{n-1,k} \right) \left(f_{ik}^{n,k} - f_{ik}^{n-1,k} \right) \right\} \mathbb{1}_{[t_{n-1}^{k}, t_{n}^{k}]}.$$

Then we take the time derivative of both sides of (5.15) and insert these expressions into the left-hand side. As a result, we get an expression for e_{ij}^{k} in which terms featuring the difference $f_{ij}^{\mathsf{n},\mathsf{k}} - f_{ij}^{\mathsf{n}-1,\mathsf{k}}$ appear along with others involving differences between pairs of flux limiters, interaction rates, and tables of games evaluated at $\mathbf{f}^{\mathsf{n}-1,\mathsf{k}}$ and $\hat{\mathbf{f}}^{\mathsf{k}}$, respectively. Applying estimate (5.12) to the first ones and invoking the Lipschitz continuity of the second ones, we finally obtain

$$\left| e_{ij}^{\mathsf{k}}(t) \right| \leq \mathcal{C} \Delta t_{\mathsf{k}}$$

for a suitable constant C > 0 independent of k, and hence

$$\int_0^t \left| e_{ij}^{\mathsf{k}}(s) \right| \, ds \le \mathcal{C} \Delta t_{\mathsf{k}} T \xrightarrow{\mathsf{k} \to \infty} 0,$$

which completes the proof.

COROLLARY 5.6 (improved regularity). The mild solution to problem (5.1) is of class C^1 in [0, T]; hence it is actually a classical solution.

Proof. The continuity of **f** implies that the right-hand side of (5.1) is continuous in t for each i = 1, ..., m and j = 1, ..., n. Hence $\frac{df_{ij}}{dt}$ is continuous.

REFERENCES

- A. AW AND M. RASCLE, Resurrection of "second order" models of traffic flow, SIAM J. Appl. Math., 60 (2000), pp. 916-938.
- [2] N. Bellomo and C. Dogbe, On the modelling of traffic and crowds: A survey of models, speculations, and perspectives, SIAM Rev., 53 (2011), pp. 409-463.
- [3] A. Bellouquid, E. De Angelis, and L. Fermo, Towards the modeling of vehicular traffic as a complex system: A kinetic theory approach, Math. Models Methods Appl. Sci., 22 (2012), 1140003.
- [4] F. Berthelin, P. Degond, M. Delitala, and M. Rascle, A model for the formation and evolution of traffic jams, Arch. Ration. Mech. Anal., 187 (2008), pp. 185–220.
- [5] I. Bonzani and L. Mussone, From experiments to hydrodynamic traffic flow models. I. Modelling and parameter identification, Math. Comput. Modelling, 37 (2003), pp. 1435–1442.
- [6] R. M. COLOMBO, Hyperbolic phase transitions in traffic flow, SIAM J. Appl. Math., 63 (2002), pp. 708–721.
- V. COSCIA, M. DELITALA, AND P. FRASCA, On the mathematical theory of vehicular traffic flow.
 II. Discrete velocity kinetic models, Internat. J. Non-Linear Mech., 42 (2007), pp. 411–421.
- [8] C. F. DAGANZO, The cell transmission model: A dynamic representation of highway traffic consistent with the hydrodynamic theory, Transp. Res. B-Meth., 28B (1994), pp. 269–287.
- [9] C. F. DAGANZO, In traffic flow, cellular automata = kinematic waves, Transp. Res. B-Meth., 40 (2006), pp. 396–403.
- [10] E. DE ANGELIS, Nonlinear hydrodynamic models of traffic flow modelling and mathematical problems, Math. Comput. Modelling, 29 (1999), pp. 83–95.
- [11] M. Delitala and A. Tosin, Mathematical modeling of vehicular traffic: A discrete kinetic theory approach, Math. Models Methods Appl. Sci., 17 (2007), pp. 901–932.

- [12] A. Deutsch and S. Dormann, Cellular Automaton Modeling of Biological Pattern Formation,
- Model. Simul. Sci. Eng. Technol., Birkhäuser, Boston, 2005.
 [13] D. C. GAZIS, R. HERMAN, AND R. W. ROTHERY, Nonlinear follow-the-leader models of traffic flow, Oper. Res., 9 (1961), pp. 545-567.
- [14] D. Helbing, Traffic and related self-driven many-particle systems, Rev. Modern Phys., 73 (2001), pp. 1067–1141.
- [15] B. S. Kerner, The Physics of Traffic, Springer, Berlin, 2004.
- [16] M. J. LIGHTHILL AND G. B. WHITHAM, On kinematic waves. II. A theory of traffic flow on long crowded roads, Proc. Roy. Soc. London Ser. A., 229 (1955), pp. 317–345.
- $[17]\,$ P. I. RICHARDS, Shock waves on the highway, Oper. Res., 4 (1956), pp. 42–51.
- [18] A. Tosin, From generalized kinetic theory to discrete velocity modeling of vehicular traffic. A stochastic game approach, Appl. Math. Lett., 22 (2009), pp. 1122-1125.
- [19] M. Treiber, A. Kesting, and D. Helbing, Delays, inaccuracies and anticipation in microscopic traffic models, Phys. A, 360 (2006), pp. 71-88.