

## 1 Rare charm and strange decays at LHCb

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6 Recent results by the LHCb collaboration in charm and strange-hadrons rare decays are summarised. A short summary of the related discussion at the ICHEP2024 conference is reported.

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## 1. Overview of the presented results

Since the discovery of charm quarks with  $K_L^0 \rightarrow \mu^+\mu^-$  decays [1], rare decays have always been an important tool to search for new physics effects. In particular, Flavour Changing Neutral Currents (FCNC) are a tool to probe the symmetries of the Standard Model (SM) and beyond. The LHCb collaboration has a large program of searches in beauty, charm and strange hadron decays. Here the latest results on charm and strange rare decays are briefly summarised and a glimpse of the related discussion at ICHEP2024 is given at the end.

### 1.1 Search for $D^0 \rightarrow \mu^+\mu^-$ decays

The  $D^0 \rightarrow \mu^+\mu^-$  decay is a very rare FCNC, additionally suppressed by helicity. In the SM this would have a short-distance branching fraction in the  $10^{-18}$  order, however, its amplitude is dominated by long-distance contributions from  $D^0 \rightarrow \gamma\gamma$ , resulting in a branching fraction of  $\mathcal{B}(D^0 \rightarrow \mu^+\mu^-) \simeq 2.7 \cdot 10^{-5} \cdot \mathcal{B}(D^0 \rightarrow \gamma\gamma)$ . Using the experimental limit on the two-photon decay [12, 13] this results in  $\mathcal{B}(D^0 \rightarrow \mu^+\mu^-) \lesssim 2.3 \cdot 10^{-11}$ . Hence, any branching fraction measurement in excess of this order of magnitude would be a unequivocal sign of new physics contributions. In fact this decay leads to the most-stringent constraints on pseudoscalar, scalar and axial vector contributions in the  $c\bar{u}\mu\bar{\mu}$  four-fermion coupling. In particular, many new physics models that allow to explain the so-called ‘‘b-anomalies’’ in the b-quark sector [19], would predict large (up to tree-level) contributions in the charm quark sector. This makes  $D^0 \rightarrow \mu^+\mu^-$  decays branching fraction constraints pre-requisites for model-building [2].

The latest measurement by the LHCb experiment uses the full Run 1 and 2 data sample [2], with a total integrated luminosity of  $9 \text{ fb}^{-1}$  of  $pp$  collision data at 7, 8 and 13 TeV. The search is performed using  $D^0$  from  $D^{*+} \rightarrow D^0\pi^+$  decays to reject a large fraction of the background. Run 2 conditions are considerably harsher in terms of multiplicity and thus combinatorial background with respect to Run 1. Therefore, compared to the first LHCb measurement [10] better multivariate methods against both combinatorial and misidentified background are employed. After full selection the residual background is mostly combinatorial and due to the misidentified  $D^{*+} \rightarrow D^0(\rightarrow \pi^+\pi^-)\pi^+$ : a component of irreducible  $\pi \rightarrow \mu$  decays in flight needs to be taken into account, but could be rejected with improved particle identification in future analyses. This residual background is precisely calibrated and constrained with control channels. The signal is normalised both to  $D^0 \rightarrow \pi^+\pi^-$  and  $D^0 \rightarrow K^-\pi^+$  decays, such that a check on the ratio of the normalisation channels branching fractions can be performed, and is consistent with PDG average.

A simultaneous two-dimensional unbinned fit is performed to search for the signal, in the dimuon invariant mass and in the  $\Delta m = m(\pi\mu\mu) - m(\mu\mu)$  variable. Six search regions are used, employing three intervals of the multivariate operator’s output, and analysing the two Runs in parallel. No evidence for an excess of events above the background expectation is seen, thus an upper limit of  $\mathcal{B}(D^0 \rightarrow \mu^+\mu^-) < 3.1(3.5) \times 10^{-9}$  at 90% (95%) CL is set, which is the most stringent published result on this channel.

During this conference, a new search for this decay has been presented by the CMS collaboration, improving this limit [11] by 35%. Additional searches are very welcome on this channel, as it keeps being the cornerstone of constraints on new physics models for what concerns up-sector FCNCs.

## 48 1.2 Search for $D^{*0}(2007) \rightarrow \mu^+\mu^-$ decays

49 The  $D^{*0}(2007) \rightarrow \mu^+\mu^-$  decay is very similar to  $D^0 \rightarrow \mu^+\mu^-$ , with the exception of the  
 50 spin of the initial meson. This removes the helicity suppression in the decay, making it sensitive  
 51 to vector couplings rather than axial-vector ones. However, given that the  $D^*$  (here and in the  
 52 following  $D^*$  represents the  $D^{*0}(2007)$ ), can decay via strong and electromagnetic interactions,  
 53 the branching fraction of any weak decay will be very suppressed. The SM branching fraction  
 54 for  $D^{*0}(2007) \rightarrow \mu^+\mu^-$  decays is tiny, at the  $10^{-19}$  level [7], however its experimental value is  
 55 unconstrained. A limit on the corresponding electron channel at  $\mathcal{B}(D^* \rightarrow e^+e^-) < 1.7 \times 10^{-6}$  [9]  
 56 has been set by the CMD-3 collaboration, but no limit on the muonic channel is present.

57 The presented search at LHCb [3] is based on data corresponding to  $9\text{fb}^{-1}$  of integrated  
 58 luminosity, as for the  $D^0 \rightarrow \mu^+\mu^-$  analysis. The search is performed with the method outlined in  
 59 Ref. [8], i.e. searching for the cascade  $B^+ \rightarrow D^*(\rightarrow \mu^+\mu^-)\pi$  decays. In practice the signature  
 60 would correspond to a  $D^*$  peak in the dimuon invariant mass of  $B^+ \rightarrow \pi^+\mu^-\mu^+$  decays. The decay  
 61 branching fraction is normalised to  $B^+ \rightarrow J/\psi K^+$  decays, where the  $J/\psi$  decays into muons.

62 After a standard selection and a multivariate selection, the main background left is combina-  
 63 torial. In addition, non-resonant  $B^+ \rightarrow \pi^+\mu^-\mu^+$  decays are present. The  $B^+ \rightarrow D^*(\rightarrow \mu^+\mu^-)\pi^+$   
 64 signal yield is obtained from a two-dimensional unbinned fit to the dimuon invariant mass and the  
 65  $\pi^+\mu^+\mu^-$  invariant mass. A yield consistent with zero is measured, consequently an upper limit on  
 66 the signal branching fraction is set at  $\mathcal{B}(D^{*0}(2007) \rightarrow \mu^+\mu^-) < 2.6(3.4) \times 10^{-8}$  at 90 (95)% CL,  
 67 which is the world first result on this channel.

## 68 1.3 Search for $\Lambda_c^+ \rightarrow p\mu^+\mu^-$ decays

69 The non-resonant  $\Lambda_c^+ \rightarrow p\mu^+\mu^-$  decays are very rare in the SM, with a predicted branching  
 70 fraction around  $10^{-8}$  for short-distance contributions, enhanced up to  $10^{-6}$  from long-distance  
 71 contributions due to intermediate vector resonances [12, 13]. In particular,  $\Lambda_c^+ \rightarrow pV$  with  
 72  $V = \phi, \rho, \eta$  and  $\omega$ , decaying to two muons, can contribute. The analysis presented here [4]  
 73 therefore divides the dimuon invariant mass in 6 regions, of which only the low invariant mass  
 74 one (below  $m_{\mu\mu} < 507.86 \text{ MeV}/c^2$ ) and the high-mass one ( $m_{\mu\mu} > 1059.46 \text{ MeV}/c^2$ ) are used to  
 75 search for the non-resonant  $\Lambda_c^+ \rightarrow p\mu^+\mu^-$  decay. This analysis uses  $pp$  collision data at 13 TeV  
 76 corresponding to  $5.4\text{fb}^{-1}$  of integrated luminosity acquired in Run 2. As often for rare decays,  
 77 a standard selection plus a multivariate classifier are employed to reject most of the background.  
 78 After that, the residual background is mostly combinatorial plus a small irreducible component due  
 79 to misidentified  $\Lambda_c^+ \rightarrow p\pi^+\pi^-$  decays.

80 The search is then performed through an unbinned fit to the  $p\mu^+\mu^-$  in regions of the dimuon  
 81 invariant mass. No amplitude analysis is performed, hence interferences between resonances or  
 82 with the non-resonant part are not considered. Significant signals are obtained in the  $\rho$ ,  $\omega$  and  $\phi$   
 83 regions. An evidence (significance of  $3\sigma$ ) is also obtained in the  $\eta$  region. The corresponding  
 84 branching fractions are determined to be

$$\begin{aligned}\mathcal{B}(\Lambda_c^+ \rightarrow p\omega) &= (9.82 \pm 1.23 \text{ (stat.)} \pm 0.73 \text{ (syst.)} \pm 2.79 \text{ (ext.)}) \times 10^{-4}, \\ \mathcal{B}(\Lambda_c^+ \rightarrow p\rho) &= (1.52 \pm 0.34 \text{ (stat.)} \pm 0.14 \text{ (syst.)} \pm 0.24 \text{ (ext.)}) \times 10^{-3}, \\ \mathcal{B}(\Lambda_c^+ \rightarrow p\eta) &= (1.67 \pm 0.69 \text{ (stat.)} \pm 0.23 \text{ (syst.)} \pm 0.34 \text{ (ext.)}) \times 10^{-3}.\end{aligned}$$

85 No evidence for  $\Lambda_c^+ \rightarrow p\mu^+\mu^-$  decays is instead obtained in the non-resonant regions, and an  
 86 upper limit of  $\frac{\mathcal{B}(\Lambda_c^+ \rightarrow p\mu^+\mu^-)}{\mathcal{B}(\Lambda_c^+ \rightarrow p\phi)\mathcal{B}(\phi \rightarrow \mu^+\mu^-)} < 0.09$  (0.10) at 90% (95%) CL is obtained for a non-resonant  
 87  $\Lambda_c^+ \rightarrow p\mu^+\mu^-$  extrapolated to the full dimuon invariant mass with a phase space model. This, using  
 88 the values of the branching fractions for  $\Lambda_c^+ \rightarrow p\phi$  and  $\phi \rightarrow \mu^+\mu^-$  decays from PDG, corresponds  
 89 to  $\mathcal{B}(\Lambda_c^+ \rightarrow p\mu^+\mu^-) < 2.9$  (3.2)  $\times 10^{-8}$  at 90% (95%) CL which is the most stringent to date.

#### 90 1.4 First observation of $\Sigma^+ \rightarrow p\mu^+\mu^-$ decays

91 The LHCb collaboration has also an expanding program of rare strange-hadrons decays. The  
 92  $\Sigma^+ \rightarrow p\mu^+\mu^-$  is a very rare FCNC: contrary to the  $\Lambda_c^+ \rightarrow p\mu^+\mu^-$  just described, there cannot be  
 93 resonances in the  $\Sigma^+ \rightarrow p\mu^+\mu^-$  due to the small dimuon invariant mass, hence this decay has only  
 94 non-resonant (rare) contributions. In particular, the short-distance SM prediction for its branching  
 95 fraction is around  $10^{-12}$ , while the long-distance one, dominating the decay, leads to a branching  
 96 fraction of  $1.2 \cdot 10^{-8} < \mathcal{B}(\Sigma^+ \rightarrow p\mu^+\mu^-) < 10.2 \cdot 10^{-8}$  [14, 15, 18]. This decay gained attention  
 97 when the HyperCP experiment saw an evidence for it with three background-free events [16], for  
 98 a branching fraction  $\mathcal{B}(\Sigma^+ \rightarrow p\mu^+\mu^-) = (8.6_{-5.4}^{+6.6} \pm 5.5) \cdot 10^{-8}$ . While the integrated branching  
 99 fraction was compatible with the SM, the dimuon invariant mass of the three events was very similar,  
 100 rather than distributed along the phase space. This pointed towards a possible  $\Sigma^+ \rightarrow pX^0(\rightarrow \mu\mu)$   
 101 decay with  $X^0$  being a particle beyond the SM, with a mass  $m_X^0 = 214.3 \pm 0.5$  MeV/ $c^2$ . Many  
 102 new physics explanations were proposed for this possible particle and searches were conducted in  
 103 different final states including a dimuon pair. While no other evidence was found, before LHCb no  
 104 other experiment could search for  $\Sigma^+ \rightarrow p\mu^+\mu^-$  decays. At LHC  $pp$  collision energies naturally  
 105 a large production of hyperons is present, with about one per every minimum bias collision. The  
 106 long hyperon lifetime, while reducing considerably the prompt background, together with the small  
 107  $p_t$  reduces also the efficiency considerably. A search for  $\Sigma^+ \rightarrow p\mu^+\mu^-$  decays was performed  
 108 already with Run 1 data [5], which produced an evidence for this decay with about ten events and  
 109 a branching fraction compatible with the SM. No structure in the dimuon invariant mass was seen  
 110 excluding the central value of the HyperCP branching fraction under the  $X^0$  hypothesis.

111 With Run 2 data, a significantly improved trigger with dedicated lines for soft dimuons and  
 112 rare strange decays was deployed [17], bringing a factor 10 larger trigger efficiency. Adding to the  
 113 increased cross-section and integrated luminosity this brings almost two orders of magnitude of  
 114 improvement. The Run 2 analysis [6] is based on  $5.4\text{fb}^{-1}$  of  $pp$  collisions at 13 TeV. The search  
 115 is based on a standard selection and a multivariate classifier to reject most of the combinatorial  
 116 background, which is the main background source. A second background is due to  $\Lambda \rightarrow p\pi^-$   
 117 decays where the pion is misidentified as a muon, and the pair is associated with an unrelated muon.  
 118 This background is rejected through a veto on the  $p\mu^-$  invariant mass in the  $p\pi^-$  hypothesis around  
 119 the  $\Lambda$  known mass, and stringent particle identification criteria. No other background is present,  
 120 due to the tight phase space and small mass of the  $\Sigma^+$  hyperon. As an example the  $\Sigma^+ \rightarrow p\pi^+\pi^-$   
 121 decay is kinematically forbidden.

122 The signal yield is obtained, after full selection, with an unbinned fit to the  $p\mu^+\mu^-$  invariant  
 123 mass. A clear peak is observed with very large significance and  $N(\Sigma^+ \rightarrow p\mu^+\mu^-) = 279 \pm 19$  signal  
 124 events are found. This constitutes the first observation of this channel, which is the rarest hyperon  
 125 and possibly the rarest baryon decay ever observed. The dimuon invariant mass distribution is also  
 126 examined: no evidence of any structure is present and the distribution is found very compatible with

127 the SM prediction. An article is in preparation where the signal yield will be also normalised to  
128  $\Sigma^+ \rightarrow p\pi^0$  decays in order to measure its branching fraction. Beyond the integrated and differential  
129 rate, this signal yield will allow for a measurement of the violation of Charge-Parity symmetry and  
130 possibly of other observables such as the forward-backward asymmetry.

## 131 2. Discussion and outlook

132 There was some discussion at the conference, around the physics interest on the  $D^{*0}(2007) \rightarrow$   
133  $\mu^+\mu^-$  decay. It was pointed out, that given the large amplitude of the strong and electromagnetic  
134 decays of the  $D^*$ , any other decay would have an insignificantly low branching fraction, hence inval-  
135 idating the presented search. This argument seems to miss that what is relevant is the modification of  
136 the amplitude, by possible new physics terms, and not the absolute value of the branching fraction,  
137 which is scaled by the large intrinsic width of the  $D^*$ . Clearly measuring the  $D^{*0}(2007) \rightarrow \mu^+\mu^-$   
138 branching fraction, especially at the SM level, will be challenging, as many rare decay searches are,  
139 however assuming to reach that experimental sensitivity, the physics interest does not depend on the  
140 branching fraction per se, but on its modification w.r.t the SM prediction. Furthermore, assuming  
141 a branching fraction enhanced by NP, the experimental observation would be very clean, with a  
142 dimuon peak in an already background-free  $B$  decay, and an easy normalisation to a decay with the  
143 same signature ( $B^+ \rightarrow J/\psi K^+$ ). In any case, the interplay between several rare charm decays with  
144 complementary couplings could corner also small new physics signals, in the same way as current  
145 sensitivity on  $B$  FCNC is cornering the different couplings involved in the “b-anomalies”.

146 All the mentioned rare decays measurements presented here are limited mostly by statistical  
147 uncertainties, hence the data currently being acquired with the LHCb upgraded detector in Run  
148 3 will allow to improve all these searches and measurements, particularly thanks to the improved  
149 trigger. The removal of the hardware trigger, which was limiting the efficiency of the charm and  
150 strange physics programs, opens room for large improvements.

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