

RESEARCH ARTICLE OPEN ACCESS

Solvability and Boundedness in Chemotaxis-Consumption Models With Oppositely Acting Nonlocal Sources

Rafael Díaz Fuentes  | Fatma Gamze Düzgün  | Silvia Frassu  | Giuseppe Viglialoro

Dipartimento di Matematica e Informatica, Università degli Studi di Cagliari, Cagliari, Italy

Correspondence: Silvia Frassu (silvia.frassu@unica.it)

Received: 10 November 2025 | **Revised:** 18 June 2026 | **Accepted:** 25 June 2026

Keywords: boundedness | chemotaxis | global existence | nonlocal growth terms

ABSTRACT

In this paper, we investigate a chemotaxis system featuring a class of external source terms that incorporate both local and nonlocal growth as well as dampening effects. The model describes the evolution of a cell density u , which migrates in response to a chemical signal v , within an impermeable habitat. Mathematically, this leads to the study of the following initial-boundary value problem

$$\begin{cases} u_t = \Delta u - \chi \nabla \cdot (u \nabla v) + f(u), & \text{in } \Omega \times (0, T_{\max}) \\ v_t = \Delta v - vu, & \text{in } \Omega \times (0, T_{\max}) \\ u_v = v_v = 0, & \text{on } \partial\Omega \times (0, T_{\max}) \\ u(x, 0) = u_0(x) \geq 0, \quad v(x, 0) = v_0(x) \geq 0, & \text{for } x \in \bar{\Omega}, \end{cases} \quad (\diamond)$$

where $x \in \Omega$, and Ω is a bounded and smooth domain of $\mathbb{R}^n$ ($n \in \mathbb{N}$) with boundary $\partial\Omega$, oriented by the outward unit normal vector ν , and $\chi > 0$. Moreover, for $a, b > 0$, $\alpha, \beta, \gamma \geq 1$, the function $f(u)$ is assumed to take one of the following forms:

$$\text{either } au^\alpha - bu^\beta \int_{\Omega} u^\gamma \text{ or } -au^\alpha + bu^\beta \int_{\Omega} u^\gamma.$$

Finally, $T_{\max} \in (0, \infty]$ indicates the maximal instant of time up to which the corresponding solutions exist. We provide conditions ensuring the absence of aggregation phenomena over time. Specifically, we show that the maximal existence time satisfies $T_{\max} = \infty$, and both u and v remain uniformly bounded for all time, in the following situations:

- ◊₁. For $f(u) = au^\alpha - bu^\beta \int_{\Omega} u^\gamma$, under the assumption $\beta \geq \alpha$, whenever either $1 \leq \alpha < 2$ and $\beta + \gamma > \frac{n}{2} + 2$, or $\alpha \geq 2$ and $\beta + \gamma > \frac{n}{2}(\alpha - 1) + \alpha$;

This is an open access article under the terms of the [Creative Commons Attribution](https://creativecommons.org/licenses/by/4.0/) License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited.

© 2026 The Author(s). *Mathematical Methods in the Applied Sciences* published by John Wiley & Sons Ltd.

$\diamond_2$. For $f(u) = -au^\alpha + bu^\beta \int_\Omega u^\gamma$, in these scenarios: $\alpha > 2$ and $\beta + \gamma < \alpha$; $\alpha > 2$, $\beta + \gamma = \alpha$ and $a > b|\Omega|$; and finally $\alpha = 2$, $\beta = \gamma = 1$ and $a > b|\Omega| + C_P \|\chi v_0\|_{L^\infty(\Omega)}$ for some $C_P = C_P(n) > 0$.

This work builds upon previous studies that established global existence and boundedness results for variants of system $(\diamond)$, in which the second equation takes the form $\tau v_t = \Delta v - v + u$ (cf. [7] for the case $\tau = 0$ and [11] for $\tau = 1$). In those settings, the source term f typically has a more ad hoc structure, namely $f(u) = au^\alpha - bu^\alpha \int_\Omega u^\beta$. Based on this framework, the present study focuses on distinguishing the different dynamical behaviors exhibited by the systems in terms of the reaction described in $(\diamond_1)$ and $(\diamond_2)$. In the first case, the model structure ensures the control of the total mass $\int_\Omega u$ over time, whereas in the second such control is no longer guaranteed, leading to substantially different qualitative dynamics.

MSC2020 Classification: 35A01, 35K55, 35Q92, 34B10, 92C17

1 | Introduction and Selected Known Results

1.1 | An Overview on Keller–Segel Models

The mathematical modeling of chemotaxis was first introduced by Keller and Segel in the 1970s [1, 2]. They proposed a system of partial differential equations governing the spatiotemporal evolution of a biological population $u = u(x, t)$ (which may represent cells, microorganisms, or bacteria) in response to a chemical substance $v = v(x, t)$, commonly termed a chemoattractant. The system is considered within a bounded and sufficiently smooth domain $\Omega \subset \mathbb{R}^n$ over a time interval $(0, T_{\max})$; naturally, $x \in \Omega$ and $t \in (0, T_{\max})$, where $T_{\max}$ represents the maximal existence time of the solution, that is, the supremum of all times for which u and v remain well defined.

More precisely, the most representative equation for the classical Keller–Segel model reads as follows

$$u_t = \Delta u - \chi \nabla \cdot (u \nabla v) \quad \text{in } \Omega \times (0, T_{\max}), \quad (1)$$

and it describes the population dynamics incorporating linear diffusion and directed movement driven by chemotaxis, with $\chi > 0$ denoting the chemotactic sensitivity. The term $-\chi \nabla \cdot (u \nabla v)$ represents the flux of the population moving up the gradient of the chemoattractant concentration.

This equation is coupled with one of the following models describing the chemoattractant dynamics: a production mechanism

$$v_t = \Delta v - v + u \quad \text{in } \Omega \times (0, T_{\max}), \quad (2)$$

or an absorption mechanism

$$v_t = \Delta v - vu \quad \text{in } \Omega \times (0, T_{\max}), \quad (3)$$

where the production/absorption terms are linear in u and idealize respectively the release or consumption of the chemical signal by the population. To ensure that the system of equations is well-posed, it is necessary to supplement it with appropriate boundary conditions and initial data for u and v . In particular, homogeneous Neumann boundary conditions indicate that the process takes place in an impermeable domain, while $u(x, 0) = u_0(x)$ and $v(x, 0) = v_0(x)$ specify the initial distributions of the cell density and the chemical signal, respectively.

In models incorporating chemoattractant production, the interplay between diffusive dispersal and chemotactic aggregation gives rise to complex spatiotemporal patterns. The outcome of this competition depends sensitively on the spatial dimension n , the initial total mass $m := \int_\Omega u_0(x) dx$, and the chemotactic sensitivity parameter χ . Under certain conditions, this interaction may lead to blow-up phenomena at finite time ($T_{\max}$ finite), characterized by the formation of singularities and unbounded population densities.

In the one-dimensional setting, global-in-time ($T_{\max} = \infty$) existence and uniform boundedness of solutions are well established. However, as the spatial dimension increases, the dynamics become substantially more intricate. Specifically, for

$n \geq 3$, solutions may blow up in finite time even when the initial mass is relatively small (see the survey [3, Section 3] and references therein).

The two-dimensional case has been extensively studied due to the emergence of the *critical mass phenomenon*, a threshold effect whereby the global behavior of solutions—whether they exist globally in time or exhibit finite-time blow-up—depends critically on the size of the initial total mass (see, e.g., [4, 5]).

1.2 | An Overview of Keller–Segel Models With Production and Logistic Effects

Incorporating logistic-type source terms into the classical Keller–Segel system governed by Equations (1) and (2) is known to enhance the prospects for global existence by introducing a damping effect.

1.2.1 | Classical Logistics

A typical example tied to what now said, is the modified equation

$$u_t = \Delta u - \chi \nabla \cdot (u \nabla v) + au - bu^\beta \quad \text{in } \Omega \times (0, T_{\max}), \text{ where } a, b > 0 \text{ and } \beta \geq 1,$$

which has been widely studied as a regularized variant of the original chemotaxis model. While the presence of the damping term indeed improves the regularity properties of solutions, global boundedness has been rigorously established only under specific conditions. In particular, for the case $\beta = 2$, global existence holds provided that the coefficient b is sufficiently large (see [6, 7]). However, blow-up may still occur when β is close to 1; see [8–10].

1.2.2 | Nonlocal Logistics

In parallel, attention has shifted to source terms involving nonlocal damping mechanisms of the form

$$f(u) = au^\alpha - bu^\alpha \int_{\Omega} u^\beta, \quad \text{with } a, b > 0, \alpha, \beta \geq 1. \quad (4)$$

Such expression models population growth regulated by a death rate that depends on the total biomass, and it can be viewed as nonlocal generalizations of logistic damping. The presence of the integral term introduces a spatially global feedback that may significantly alter the long-time dynamics of solutions.

A central analytical question is whether such nonlocal damping is strong enough to prevent finite-time blow-up and guarantee boundedness of solutions; that is, for the equation

$$u_t = \Delta u - \chi \nabla \cdot (u \nabla v) + au^\alpha - bu^\alpha \int_{\Omega} u^\beta,$$

can global boundedness arise solely from the nonlocal term, or are further constraints required? Studies incorporating nonlinear diffusion, chemotactic sensitivity, signal production, or multiple chemical signals establish that the nonlocal term can prevent blow-up under suitable conditions [11–14]. (Also in $\mathbb{R}^n$, nonlocal damping with Poisson-type chemoattractant interactions has a regularizing effect: [15–18].)

1.2.3 | Gradient-Dependent Logistics

Furthermore, finite-time blow-up can be prevented by incorporating a logistic-type source $h = h(u, |\nabla u|)$ into the population cells' equation $u_t = \Delta u - \chi \nabla \cdot (u \nabla v)$. Typical forms include $h(u, |\nabla u|) = au^\alpha - bu^\beta - c|\nabla u|^\gamma$, or alternatively, $h(u, |\nabla u|) = au^\alpha - bu^\alpha \int_{\Omega} u^\beta - c|\nabla u|^\gamma$, with positive constants a, b, c and exponents $\alpha, \beta, \gamma \geq 1$. While gradient-dependent nonlinearities arising from external sources often interpreted biologically as mechanisms of *accidental death* have been extensively studied within single-equation models (see, e.g., [19], and references therein), their influence within chemotactic systems has been lately garnering significant attention. Notably, recent investigations such as [20–23] have begun to elucidate their role in biologically and ecologically relevant contexts.

Remark 1 (Consumption models with classical and gradient-dependent logistics). Above, we have summarized the results we consider most relevant to the present study, focusing on Keller–Segel type models in which the chemical signal(s) are produced by the same cell population. In parallel, for consumption-type systems, namely those in which the signal equation is of the form (3), there also exist investigations incorporating classical logistic terms (cf. [24–26]), as well as studies addressing gradient-dependent logistic mechanisms [27]. However, to the best of our knowledge, no study has yet addressed consumption models involving nonlocal logistic terms. As we shall emphasize later, this gap in the literature represents one of the main motivations of the present work.

2 | Presentation of the Model and Motivations; Claim of the Main Result

In accordance with Remark 1, and as it will be discussed in detail below, our analysis is devoted to consumption-type Keller–Segel systems endowed with nonlocal source terms. Furthermore, we shall consider logistic components in which the nonlocality constitutes, on the one hand, a natural generalization of that appearing in (4), and, on the other hand, may influence both proliferative and dissipative dynamics within the population.

2.1 | Basic Description of the Model

In this paper, we study the initial-boundary value problem

$$\begin{cases} u_t = \Delta u - \chi \nabla \cdot (u \nabla v) + f(u) & \text{in } \Omega \times (0, T_{\max}), \\ v_t = \Delta v - vu & \text{in } \Omega \times (0, T_{\max}), \\ u_v = v_v = 0 & \text{on } \partial\Omega \times (0, T_{\max}), \\ u(x, 0) = u_0(x), \quad v(x, 0) = v_0(x) & x \in \bar{\Omega}, \end{cases} \quad (5)$$

where $x \in \Omega \subset \mathbb{R}^n$ ($n \in \mathbb{N}$), Ω is a bounded domain with smooth boundary $\partial\Omega$, and the parameter χ is positive. The initial data $u_0(x)$, $v_0(x)$ are assumed to be nonnegative and sufficiently regular and for $a, b > 0$, $\alpha, \beta, \gamma \geq 1$ the function $f(u)$ may take one of the following forms:

$$\text{either } au^\alpha - bu^\beta \int_{\Omega} u^\gamma \text{ or } -au^\alpha + bu^\beta \int_{\Omega} u^\gamma. \quad (6)$$

As anticipated, and reiterated here for the sake of clarity, the notation $(\cdot)_v$ denotes the outward normal derivative on $\partial\Omega$, while $T_{\max}$ represents the maximal existence time of the solution.

In line with the previous sections, the first equation models the population density u via diffusion, chemotaxis, and local growth or decay, while the second governs the chemoattractant v through diffusion and consumption. Homogeneous Neumann conditions enforce no flux on $\partial\Omega$.

2.1.1 | Interpretation of $f(u)$ and Its Role in the Investigation

The expressions of $f(u)$ in (6) describe nonlocal damping or growth mechanisms, where the evolution of $u(x, t)$ depends not only on the local value but also on the global quantity $\int_{\Omega} u^\gamma$, representing the total population mass raised to a power and introducing spatially averaged feedback.

Precisely, in the first scenario the term au^α models local population growth, typically driven by reproduction or proliferation processes. The accompanying nonlocal inhibitory term $-bu^\beta \int_{\Omega} u^\gamma$ serves as a saturation mechanism, reflecting effects such as overcrowding or competition for limited resources that depend on the overall population size. Exactly in the opposite direction, the second formulation features a leading local decay term $-au^\alpha$, which may be interpreted as natural mortality or dissipation, while the nonlocal term $bu^\beta \int_{\Omega} u^\gamma$ can induce growth once the total population surpasses a certain threshold.

Despite differences in the source term, both forms fit a unified framework for analyzing boundedness and global existence. However, from a mathematical point of view, the choice of logistic term in (6) is crucial and determines the technical approach. In the first case, uniform-in-time boundedness of the mass, $\int_{\Omega} u$, holds for all $t \in (0, T_{\max})$, while in the second

case this property is lost. Accordingly, the analysis must be adapted depending on the form of $f(u)$ (see Remark 4 and Section 5 for details).

2.2 | Formulation of the Main Theorem

In order to present our result, here we assume

$$\begin{cases} \text{for some } \sigma \in (0, 1) \text{ and } n \in \mathbb{N}, \Omega \subset \mathbb{R}^n \text{ is a bounded domain of class } C^{2+\sigma}, \text{ with boundary } \partial\Omega, \\ u_0, v_0 : \bar{\Omega} \rightarrow \mathbb{R}_0^+, \text{ with } u_0, v_0 \in C^{2+\sigma}(\bar{\Omega}) \text{ such that } u_{0\nu} = 0 \text{ and } v_{0\nu} = 0 \text{ on } \partial\Omega. \end{cases} \quad (7)$$

Subsequently, we will prove the following

Theorem 2.1. *Let the hypotheses in (7) be fulfilled and $\chi, a, b > 0$ and $\alpha, \beta, \gamma \geq 1$. Then, problem (5) admits a unique solution*

$$(u, v) \in C^{2+\sigma, 1+\frac{\sigma}{2}}(\bar{\Omega} \times [0, \infty)) \times C^{2+\sigma, 1+\frac{\sigma}{2}}(\bar{\Omega} \times [0, \infty))$$

with $0 \leq u, v \in L^\infty(\Omega \times (0, \infty))$, in these circumstances:

1. For $f(u) = au^\alpha - bu^\beta \int_\Omega u^\gamma$, with any $a > 0$ and $\beta \geq \alpha$ whenever
 - (a) $\alpha \geq 2$ and $\beta + \gamma > \frac{n}{2}(\alpha - 1) + \alpha$;
 - (b) $1 \leq \alpha < 2$ and $\beta + \gamma > \frac{n}{2} + 2$;
2. For $f(u) = -au^\alpha + bu^\beta \int_\Omega u^\gamma$, whenever $\beta \geq 1$ and
 - (a) $\alpha > 2$, $\beta + \gamma < \alpha$, with any $a > 0$;
 - (b) $\alpha > 2$, $\beta + \gamma = \alpha$, with $a > b|\Omega|$;
 - (c) $\alpha = 2$, $\beta = \gamma = 1$, with $a > b|\Omega| + \mathcal{K}\|\chi v_0\|_{L^\infty(\Omega)}$ for some $\mathcal{K} = \mathcal{K}(n) > 0$.

Remark 2. In the absence of a logistic term, that is, when $f(u) = 0$, our model (5) reduces to pure consumption mechanisms, which remain an interesting subject of study, particularly regarding the existence of blow-up solutions. In this context, in [28] and [29] the authors obtained boundedness under suitable smallness assumptions on $\chi\|v_0\|_{L^\infty(\Omega)}$, while in [26] it is established that, in the presence of a classical logistic term of the form $bu - au^2$, boundedness is guaranteed provided $a > \varphi(\chi\|v_0\|_{L^\infty(\Omega)}, n)$, for some continuous expression of $\varphi(\xi, n)$, $\xi > 0$. Notably, the limit scenario obtained for $b = 0$ in case 2(c) of Theorem 2.1 qualitatively recovers [26, Theorem 1.1].

3 | Preliminary Results

The upcoming auxiliary results will be used in the proof of our main estimates. Moreover, we shall adopt the following conventions:

- ♦ all constants c_i , $i = 1, 2, \dots$, appearing below are tacitly assumed to be positive and to depend on the parameters of the problem as well as on certain norms of the initial data. For the sake of notational simplicity, this dependence will be omitted; however, any dependence on auxiliary parameters, or other constants that will appear in the claim of the theorems, introduced in the course of the analysis will be stated explicitly;
- ♦ the symbol ϵ will denote an arbitrary positive real number; multiplication by another constant or addition of another arbitrary constant will not be explicitly performed, and for convenience, the final result will also be denoted by ϵ .

Lemma 3.1. *Let Ω comply with the hypotheses in (7), and for $n \geq 3$ let*

$$p := \frac{2n}{n-2}.$$

Additionally, let q, r satisfy $1 \leq r < q < p$ and $\frac{q}{r} < \frac{2}{r} + \frac{2}{n}$. Then, for all $\epsilon > 0$ there exists $C_0 = C_0(\epsilon) > 0$ such that for all $\varphi \in H^1(\Omega) \cap L^r(\Omega)$,

$$\|\varphi\|_{L^q(\Omega)}^q \leq C_0 \|\varphi\|_{L^r(\Omega)}^\delta + \epsilon \|\nabla \varphi\|_{L^2(\Omega)}^2 + \|\varphi\|_{L^2(\Omega)}^2, \quad (8)$$

where

$$\delta := \frac{2(1-\lambda)q}{2-\lambda q} \quad \text{with} \quad \lambda := \frac{\frac{1}{r} - \frac{1}{q}}{\frac{1}{r} - \frac{1}{p}} \in (0, 1).$$

The same conclusion holds for $n \in \{1, 2\}$ whenever q, r fulfill $1 \leq r < q$ and $\frac{q}{r} < \frac{2}{r} + \frac{2}{n}$.

Proof. The proof can be found in [30, Lemma 3.1]. □

Lemma 3.2. Let $T > 0$ and $\phi : (0, T) \times \mathbb{R}_0^+ \rightarrow \mathbb{R}$. If $0 \leq y \in C^0([0, T)) \cap C^1((0, T))$ is such that

$$y' \leq \phi(t, y) \quad \text{for all } t \in (0, T),$$

and there is $y_1 > 0$ with the property that whenever $y > y_1$ for some $t \in (0, T)$ one has that $\phi(t, y) \leq 0$, then

$$y \leq \max\{y_1, y(0)\} \quad \text{on } (0, T).$$

Proof. For the proof we refer the reader to [30, Lemma 3.3]. □

4 | Local Existence, Parabolic Regularity Results, and Globality

Let us establish existence and uniqueness of local classical solutions to problem (5). We refrain from presenting the full proof, as it is routine in this framework, and instead we emphasize the main steps.

Lemma 4.1. Let the hypotheses in (7) be fulfilled. Additionally, let $\chi, a, b > 0$ and $\alpha, \beta, \gamma \geq 1$ and $f(u)$ as in (6). Then, there exist $T_{\max} \in (0, \infty]$ and a unique couple of functions (u, v) , with

$$(u, v) \in C^{2+\sigma, 1+\frac{\sigma}{2}}(\bar{\Omega} \times [0, T_{\max})) \times C^{2+\sigma, 1+\frac{\sigma}{2}}(\bar{\Omega} \times [0, T_{\max})),$$

solving problem (5), and complying with

$$u \geq 0 \quad \text{and} \quad 0 \leq v \leq \|v_0\|_{L^\infty(\Omega)} \quad \text{in } \bar{\Omega} \times [0, T_{\max}). \quad (9)$$

Finally,

$$\text{if } T_{\max} < \infty \quad \text{then} \quad \limsup_{t \rightarrow T_{\max}} \|u(\cdot, t)\|_{L^\infty(\Omega)} = \infty. \quad (10)$$

Proof. Classical fixed-point arguments in Hölder spaces ensure the local well-posedness of the problem, establishing the existence and uniqueness of a classical solution $(u, v) \in C^{2+\sigma, 1+\frac{\sigma}{2}}(\bar{\Omega} \times [0, T]) \times C^{2+\sigma, 1+\frac{\sigma}{2}}(\bar{\Omega} \times [0, T])$ for some sufficiently small $T = T(u_0, v_0) > 0$. By virtue of the maximum principle, the structural condition $f(0) = 0$, together with the homogeneous Neumann boundary condition and the assumption $u_0 \geq 0$ on $\bar{\Omega}$, guarantees the nonnegativity of u for all $t \in [0, T]$; consequently, the same principle yields the second relation in (9). Taking T as initial time and $u(\cdot, T)$ as the initial condition, the above reasoning would provide a solution $(\hat{u}, \hat{v})$ defined on $\bar{\Omega} \times [T, \hat{T}]$, for some $\hat{T} > 0$, which by uniqueness would be the prolongation of (u, v) (exactly, from $\bar{\Omega} \times [0, T]$ to $\bar{\Omega} \times [0, \hat{T}]$). This procedure can be repeated up to construct a maximal interval time $[0, T_{\max})$ of existence with $T_{\max} \in (0, \infty]$; furthermore, if $T_{\max} < \infty$, this time corresponds to a finite-time blow-up, in the sense of (10). For more details we indicate [11, Proposition 4]. □

From this point onward (u, v) will denote the unique local solution to problem (5), defined on $\Omega \times (0, T_{\max})$ and guaranteed by Lemma 4.1.

Lemma 4.2. For every $q > 1$ there exist $C_P = C_P(n, q) > 0$ and $c_1 = c_1(C_P)$ for which v is such that

$$\int_0^t e^s \int_{\Omega} |\Delta v(\cdot, s)|^q ds \leq C_P 2^{q-1} \|v_0\|_{L^\infty(\Omega)}^q \int_0^t e^s \int_{\Omega} u(\cdot, s)^q ds + c_1 e^t \quad \text{for all } t \in (0, T_{\max}).$$

Additionally, if $u \in L^\infty((0, T_{\max}); L^k(\Omega))$ for some $k > n$, then

$$v \in L^\infty((0, T_{\max}); W^{1,\infty}(\Omega)). \quad (11)$$

Proof. Since v solves $v_t = \Delta v - vu$ in $\Omega \times (0, T_{\max})$, it equivalently solves $v_t = \Delta v - v + v(1 - u)$ in $\Omega \times (0, T_{\max})$. By invoking maximal Sobolev regularity results for this latter equation (see [31] or [32, Theorem 2.3]) and following the methodology outlined in [33–35], one obtains that there exists $C_P = C_P(n, q)$ such that

$$\int_0^t e^s \int_{\Omega} |\Delta v(\cdot, s)|^q ds \leq C_P \left(\|v_0\|_{W^{2,q}(\Omega)} + \int_0^t e^s \int_{\Omega} |v(1 - u(\cdot, s))|^q ds \right), \quad \text{for all } t \in (0, T_{\max}). \quad (12)$$

Exploiting the upper bound in (9), and $(A + B)^q \leq 2^{q-1}(A^q + B^q)$, with $A, B \geq 0$ and $q \geq 1$ (and such an inequality might be employed in other places of the paper without explicit mention), we have that $|v(1 - u(\cdot, s))|^q \leq 2^{q-1} \|v_0\|_{L^\infty(\Omega)}^q (1 + u(\cdot, s)^q)$, and by substituting this into (12) yields the desired conclusion.

As for the final claim, we first note that for $k > n$ it holds that $-\frac{1}{2} - \frac{n}{2k} > -1$. Making use of the standard L^k - L^∞ estimate for the Neumann heat semigroup (see, for instance, [36, Lemma 1.3]), we infer that

$$\begin{aligned} \|\nabla v(\cdot, t)\|_{L^\infty(\Omega)} &\leq \|\nabla e^{t\Delta} v(\cdot, 0)\|_{L^\infty(\Omega)} + \int_0^t \|\nabla e^{(t-s)\Delta} (u(\cdot, s)v(\cdot, s))\|_{L^\infty(\Omega)} ds \\ &\leq \|\nabla v_0\|_{L^\infty(\Omega)} + c_2 \int_0^t \left(1 + (t-s)^{-\frac{1}{2} - \frac{n}{2k}}\right) e^{-\mu_1(t-s)} \|u(\cdot, s)v(\cdot, s)\|_{L^k(\Omega)} ds, \end{aligned}$$

for all $t \in (0, T_{\max})$, where μ_1 denotes the first nonzero eigenvalue of $-\Delta$ in Ω under homogeneous Neumann boundary conditions. Since, again, v is bounded, the assumption on u ensures that $\|u(\cdot, s)v(\cdot, s)\|_{L^k(\Omega)}$ remains uniformly controlled for $s \in (0, T_{\max})$, we may conclude that $\|\nabla v(\cdot, t)\|_{L^\infty(\Omega)} \leq c_3$ for all $t \in (0, T_{\max})$ and some $c_3 = c_3(k, \mu_1)$, and thus the claim follows. $\square$

4.1 | A Boundedness Criterion

As specified, when relation (10) holds the solution is said to exhibit blow-up at some finite time $T_{\max}$ in the $L^\infty(\Omega)$ -norm. The next result rules out this blow-up scenario and ensures globality of solutions.

Lemma 4.3 (A boundedness criterion). *There exists $k_0 = k_0(n, \alpha, \beta, \gamma) \geq n + 2$ such that whenever $u \in L^\infty((0, T_{\max}); L^k(\Omega))$ for some $k > k_0$ then we have $T_{\max} = \infty$ and in particular $u \in L^\infty((0, \infty); L^\infty(\Omega))$.*

Proof Once it has been established that for some $k > k_0$ one has the inclusion $u \nabla v \in L^\infty((0, T_{\max}); L^{q_1}(\Omega))$ (see (11)), with $q_1 > n + 2$, the uniform-in-time boundedness of u in $L^\infty(\Omega)$ for all $t \in (0, T_{\max})$ follows by appropriately adjusting k_0 also in term of α, β, γ so as to satisfy all the structural assumptions of [37, Lemma A.1]. Subsequently, by invoking the dichotomy criterion in (10) of Lemma 4.1, we necessarily have $T_{\max} = \infty$. $\square$

Remark 3 (On the role of the value k_0 in Lemma 4.3). Within the framework of our model, the expressions of $f(u)$ given in (6) effectively play the role of the source term $g(x, t)$ appearing in [37, Lemma A.1]. Building upon this correspondence, we emphasize the following considerations, which will be relevant for the subsequent analysis:

- 1 The parameter k_0 should be regarded merely as a lower bound which, for certain combinations of α, β , and γ , may suffice to guarantee the boundedness of the solution, although in other parameter regimes it may need to be enlarged. In this sense, k_0 is to be understood as a quantity which, while in principle computable, is neither optimal nor uniquely determined, since its precise value depends on the interplay among n, α, β , and γ . Accordingly, when $f(u)$ is taken in the forms of (6), no specific restrictions on α, β , and γ will be required, and the value of k_0 may be tacitly assumed to increase whenever necessary, so to ensures the mathematical consistency of the subsequent arguments.
- 2 In contrast, when $f(u)$ is considered in the second form of (6) and the corresponding analysis is confined to the case $\alpha = 2$ and $\beta = \gamma = 1$, we can fix $k_0 = n + 2$; indeed, from $u \in L^\infty((0, T_{\max}), L^k(\Omega))$, with $k > k_0$, we would have:
 - $\int_{\Omega} u(x, t) dx \leq L$, for all $t \in (0, T_{\max})$, and henceforth $g(x, t) = -au(x, t)^2 + bu(x, t) \int_{\Omega} u(x, t) dx \leq -au(x, t)^2 + bLu(x, t)$, that is, $g(x, t) \leq \frac{b^2 L^2}{4a}$ on $(0, T_{\max})$, and the second hypothesis in [37, Lemma A.1, (A.6)] is accomplished;
 - $\nabla v \in L^\infty((0, T_{\max}), L^\infty(\Omega))$ (see (11)), and subsequently

$$f(x, t) = -\chi u(x, t) \nabla v(x, t) \in L^\infty((0, T_{\max}), L^k(\Omega)),$$

so that also the first hypothesis in [37, Lemma A.1, (A.6)] is fulfilled.

5 | A Priori Estimates

In this section, we derive a priori estimates for the local solution (u, v) to model (5). These estimates yield bounds on the solution that depend essentially only on the prescribed initial data, and they represent a fundamental tool in establishing the uniform-in-time boundedness of the solution itself.

Remark 4. We remark that, in the analysis that follows, a certain property of the solution is guaranteed in one case but not in the other, namely the uniform-in-time boundedness of the mass $\int_{\Omega} u$. Indeed, when $f(u)$ takes the first form in (6), the mass can be controlled; in the second case, this is no longer possible. As a result, the mathematical approach differs: in the first case, we will make use of the Gagliardo–Nirenberg inequality and the mentioned boundedness of the mass to handle, by means of the nonpositive nonlocal term of the source, contributions from the cross diffusion term and the nonnegative local part of the source itself. Conversely, in the second case, the nonnegative nonlocal term and the cross-diffusion contribution are controlled by the nonpositive local term of the logistic component, without it being previously required to ensure mass boundedness.

In the next steps we will dedicate to derive uniform-in-time bounds of u in some $L^k(\Omega)$ -spaces, for $k > 1$. More specifically, let k_0 be the value whose existence is ensured by Lemma 4.3 for the purpose of establishing a boundedness criterion. Since no upper bound restriction on k is needed in our analysis, throughout the remainder of the paper the notation $k > k_0$ will be understood to denote an arbitrarily large value of k .

5.1 | The Case $f(u) = au^{\alpha} - bu^{\beta} \int_{\Omega} u^{\gamma}$

Lemma 5.1 (Boundedness of the mass). *Let $\beta \geq \alpha \geq 1$ and $\gamma \geq 1$. Then, there exists $m_0 > 0$ such that*

$$\int_{\Omega} u(x, t) \, dx \leq m_0 \quad \text{for all } t \in (0, T_{\max}). \quad (13)$$

Proof For $\beta \geq \alpha$, we integrate over Ω the first equation of problem (5) so that by Hölder's inequalities,

$$\begin{aligned} y'(t) &:= \frac{d}{dt} \int_{\Omega} u \leq a|\Omega|^{\frac{\beta-\alpha}{\beta}} \left(\int_{\Omega} u^{\beta} \right)^{\frac{\alpha}{\beta}} - b|\Omega|^{(1-\gamma)} \int_{\Omega} u^{\beta} \left(\int_{\Omega} u \right)^{\gamma} \\ &= \left(\int_{\Omega} u^{\beta} \right)^{\frac{\alpha}{\beta}} \left[a|\Omega|^{\frac{\beta-\alpha}{\beta}} - b|\Omega|^{(1-\gamma)} \left(\int_{\Omega} u^{\beta} \right)^{\frac{\beta-\alpha}{\beta}} \left(\int_{\Omega} u \right)^{\gamma} \right] \\ &\leq \left(\int_{\Omega} u^{\beta} \right)^{\frac{\alpha}{\beta}} \left[a|\Omega|^{\frac{\beta-\alpha}{\beta}} - b|\Omega|^{(1-\gamma) - \frac{(\beta-1)(\beta-\alpha)}{\beta}} \left(\int_{\Omega} u \right)^{\beta-\alpha+\gamma} \right] \quad \text{for all } t \in (0, T_{\max}). \end{aligned}$$

Now we apply Lemma 3.2 with $T = T_{\max}$, $\rho(t) := \left(\int_{\Omega} u^{\beta} \right)^{\frac{\alpha}{\beta}} \geq 0$ on $(0, T_{\max})$,

$$\phi(t, y) = \rho(t) \left[a|\Omega|^{\frac{\beta-\alpha}{\beta}} - b|\Omega|^{(1-\gamma) - \frac{(\beta-1)(\beta-\alpha)}{\beta}} (y(t))^{\beta-\alpha+\gamma} \right],$$

$$y_0 = y(0) = \int_{\Omega} u_0 \text{ and } y_1 := \left(\frac{a|\Omega|^{\frac{\beta-\alpha}{\beta}}}{b|\Omega|^{(1-\gamma) - \frac{(\beta-1)(\beta-\alpha)}{\beta}}} \right)^{\frac{1}{\beta-\alpha+\gamma}}, \text{ so concluding with } m_0 = \max\{y_0, y_1\}. \quad \square$$

Let us start by considering the case $\alpha \geq 2$.

Lemma 5.2. *Assume that $\alpha, \beta, \gamma \geq 1$ satisfy that*

$$\beta \geq \alpha \geq 2 \quad \text{and} \quad \beta + \gamma > \frac{n(\alpha-1)}{2} + \alpha. \quad (14)$$

Then $u \in L^{\infty}((0, T_{\max}); L^k(\Omega))$, for all $k > k_0$.

Proof. From the first equation in (5), an integration by parts yields for all $k > k_0$

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u^k &= -k(k-1) \int_{\Omega} u^{k-2} |\nabla u|^2 + k(k-1) \chi \int_{\Omega} u^{k-1} \nabla u \cdot \nabla v + ka \int_{\Omega} u^{k+\alpha-1} \\ &\quad - kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) \\ &= -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 - (k-1) \chi \int_{\Omega} u^k \Delta v + ka \int_{\Omega} u^{k+\alpha-1} \\ &\quad - kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) \quad \text{on } (0, T_{\max}). \end{aligned} \quad (15)$$

Since $\alpha > 1$, the Young inequality gives for some $c_4 = c_4(k)$

$$-(k-1) \chi \int_{\Omega} u^k \Delta v \leq \int_{\Omega} u^{k+\alpha-1} + c_4 \int_{\Omega} |\Delta v|^{\frac{k+\alpha-1}{\alpha-1}} \quad \text{for all } t \in (0, T_{\max}). \quad (16)$$

We now add to both sides of (15) the term $\int_{\Omega} u^k$ and then we multiply by e^t . Since $e^t \frac{d}{dt} \int_{\Omega} u^k + e^t \int_{\Omega} u^k = \frac{d}{dt} (e^t \int_{\Omega} u^k)$ and taking into account bound (16), an integration over $(0, t)$ provides for some $c_5 = c_5(k)$ on $(0, T_{\max})$

$$\begin{aligned} e^t \int_{\Omega} u^k - \int_{\Omega} u_0^k + kb \int_0^t e^s \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) ds &\leq -\frac{4(k-1)}{k} \int_0^t e^s \left(\int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 \right) ds \\ &\quad + \int_0^t e^s \left(\int_{\Omega} u^k \right) ds + c_5 \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds + c_4 \int_0^t e^s \left(\int_{\Omega} |\Delta v|^{\frac{k+\alpha-1}{\alpha-1}} \right) ds. \end{aligned} \quad (17)$$

Let us now address the term $c_4 \int_0^t e^s \left(\int_{\Omega} |\Delta v|^{\frac{k+\alpha-1}{\alpha-1}} \right) ds$. By applying Lemma 4.2 with $q = \frac{k+\alpha-1}{\alpha-1}$, we have for all $t \in (0, T_{\max})$ and some $c_6 = c_6(k, C_P)$

$$c_4 \int_0^t e^s \left(\int_{\Omega} |\Delta v(\cdot, s)|^{\frac{k+\alpha-1}{\alpha-1}} \right) ds \leq c_4 C_P 2^{\frac{k}{\alpha-1}} \|v_0\|_{L^{\infty}(\Omega)}^{\frac{k+\alpha-1}{\alpha-1}} \int_0^t e^s \left(\int_{\Omega} (u(\cdot, s))^{\frac{k+\alpha-1}{\alpha-1}} \right) ds + c_6 e^t.$$

From the condition $\alpha \geq 2$ we have that $\frac{k+\alpha-1}{\alpha-1} \leq k + \alpha - 1$, the Young inequality leads for some $c_7 = c_7(k, C_P)$ to

$$c_4 C_P 2^{\frac{k}{\alpha-1}} \|v_0\|_{L^{\infty}(\Omega)}^{\frac{k+\alpha-1}{\alpha-1}} \int_{\Omega} u^{\frac{k+\alpha-1}{\alpha-1}} \leq c_4 C_P 2^{\frac{k}{\alpha-1}} \|v_0\|_{L^{\infty}(\Omega)}^{\frac{k+\alpha-1}{\alpha-1}} \int_{\Omega} u^{k+\alpha-1} + c_7 \quad \text{on } (0, T_{\max}).$$

(For the limiting case $\alpha = 2$, the constant c_7 may naturally be set to 0.) Combining the previous estimate with (17) then yields for some $c_8 = c_8(k, C_P)$

$$\begin{aligned} e^t \int_{\Omega} u^k - \int_{\Omega} u_0^k + kb \int_0^t e^s \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) ds &\leq -\frac{4(k-1)}{k} \int_0^t e^s \left(\int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 \right) ds \\ &\quad + \int_0^t e^s \left(\int_{\Omega} u^k \right) ds + c_8 \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds + c_7 e^t \quad \text{for all } t \in (0, T_{\max}). \end{aligned} \quad (18)$$

Here, we turn our attention to the third integral on the right-hand side of (18). Observing that $\int_{\Omega} u^{k+\alpha-1} = \|u^{\frac{k}{2}}\|_{L^{\frac{2(k+\alpha-1)}{k}}(\Omega)}^{\frac{2(k+\alpha-1)}{k}}$, we aim to exploit Lemma 3.1 with $\varphi := u^{\frac{k}{2}}$ and appropriately chosen q and r . From the definition of k_0 and condition (14), for any $k > k_0$ it is feasible to set

$$k' := \frac{k + \beta + \gamma - 1}{2}, \quad (19)$$

which fulfills

$$\max \left\{ \gamma, \frac{k}{2}, \frac{n(\alpha-1)}{2} \right\} < k' < k + \alpha - 1. \quad (20)$$

In this way, for $n \geq 3$ and for

$$p := \frac{2n}{n-2}, \quad q := \frac{2(k+\alpha-1)}{k}, \quad r := \frac{2k'}{k}$$

straightforward calculations reveal $1 \leq r < q < p$ and $\frac{q}{r} < \frac{2}{r} + \frac{2}{n}$. Therefore, we infer from (8) that for all $t \in (0, T_{\max})$ and some $c_9 = c_9(k, C_P, C_0)$

$$c_8 \int_{\Omega} u^{k+\alpha-1} = c_8 \|u^{\frac{k}{2}}\|_{L^{\frac{2(k+\alpha-1)}{k}}(\Omega)}^{\frac{2(k+\alpha-1)}{k}} \leq \frac{2(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + c_8 \int_{\Omega} u^k + c_9 \left(\int_{\Omega} u^{k'} \right)^{\frac{\delta}{r}}. \quad (21)$$

Here, the interpolation inequality (see [38, page 93]) leads on $(0, T_{\max})$ to

$$\begin{aligned} \left(\int_{\Omega} u^{k'} \right)^{\frac{\delta}{r}} &= \|u\|_{L^{k'}(\Omega)}^{b_1} \leq \|u\|_{L^{\gamma}(\Omega)}^{a_1 b_1} \|u\|_{L^{k+\beta-1}(\Omega)}^{(1-a_1)b_1} \\ &= \left(\|u\|_{L^{\gamma}(\Omega)}^{\gamma} \|u\|_{L^{k+\beta-1}(\Omega)}^{k+\beta-1} \right)^{\frac{a_1 b_1}{\gamma}} \|u\|_{L^{k+\beta-1}(\Omega)}^{\left[1-a_1-\frac{a_1(k+\beta-1)}{\gamma}\right] b_1}, \end{aligned} \quad (22)$$

where

$$b_1 = b_1(q) := \frac{k'\delta(q)}{r} = \frac{k'\delta}{r}, \quad a_1 := \frac{\frac{1}{k'} - \frac{1}{k+\beta-1}}{\frac{1}{\gamma} - \frac{1}{k+\beta-1}} \in (0, 1). \quad (23)$$

Recalling the expression of k' in (19) and the range of γ in (14), some calculations provide

$$\left[1 - a_1 - \frac{a_1(k+\beta-1)}{\gamma} \right] b_1 = 0 \quad \text{and} \quad \frac{a_1 b_1}{\gamma} < 1.$$

Consequently, applying Young's inequality, relation (22) reads for some $c_{10} = c_{10}(k, C_P, C_0)$

$$c_9 \left(\int_{\Omega} u^{k'} \right)^{\frac{\delta}{r}} \leq c_9 \left(\|u\|_{L^{\gamma}(\Omega)}^{\gamma} \|u\|_{L^{k+\beta-1}(\Omega)}^{k+\beta-1} \right)^{\frac{a_1 b_1}{\gamma}} \leq kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) + c_{10} \quad \text{for all } t \in (0, T_{\max}),$$

which in conjunction with (21) implies for all $t \in (0, T_{\max})$,

$$c_8 \int_{\Omega} u^{k+\alpha-1} \leq \frac{2(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + c_8 \int_{\Omega} u^k + kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) + c_{10}. \quad (24)$$

For $n \in \{1, 2\}$ the arguments are similar once relation (20) is replaced for $n = 1$ by

$$\max \left\{ \gamma, \frac{k}{2}, \frac{(\alpha-1)}{2} \right\} < k' < k + \alpha - 1$$

and for $n = 2$ by

$$\max \left\{ \gamma, \frac{k}{2}, (\alpha-1) \right\} < k' < k + \alpha - 1.$$

With this range of k' , some computations imply that q, r fulfill $1 \leq r < q$ and $\frac{q}{r} < \frac{2}{r} + \frac{2}{n}$, so we can apply also in these cases Lemma 3.1, leading to inequality (24) up to some constants.

Now we focus on the second integral at the right-hand side of estimate (24): the Gagliardo–Nirenberg inequality and (13) produce for

$$\theta_1 := \frac{\frac{k}{2} - \frac{1}{2}}{\frac{k}{2} + \frac{1}{n} - \frac{1}{2}} \in (0, 1)$$

and all $\hat{c} > 0$ and some $c_{11} = c_{11}(k, \hat{c})$, $c_{12} = c_{12}(k, \hat{c})$, this bound on $(0, T_{\max})$:

$$\hat{c} \int_{\Omega} u^k = \hat{c} \|u^{\frac{k}{2}}\|_{L^2(\Omega)}^2 \leq c_{11} \|\nabla u^{\frac{k}{2}}\|_{L^2(\Omega)}^{2\theta_1} \|u^{\frac{k}{2}}\|_{L^{\frac{2}{k}}(\Omega)}^{2(1-\theta_1)} + c_{11} \|u^{\frac{k}{2}}\|_{L^{\frac{2}{k}}(\Omega)}^2 \leq c_{12} \left(\int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 \right)^{\theta_1} + c_{12}.$$

In turn, we have from the Young inequality that for some $c_{13} = c_{13}(k, \hat{c})$

$$\hat{c} \int_{\Omega} u^k \leq \frac{2(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + c_{13} \quad \text{for all } t \in (0, T_{\max}). \quad (25)$$

By inserting estimate (24) into (18) and taking into account bound (25), we arrive for some $c_{14} = c_{14}(k)$ and $c_{15} = c_{15}(k, C_P, C_0, \hat{e})$ at

$$e^t \int_{\Omega} u^k \leq c_{15} e^t + c_{14} \quad \text{on } (0, T_{\max}),$$

so the claim is proved. $\square$

Let us turn now our attention to the case $1 \leq \alpha < 2$.

Lemma 5.3. Assume that $\alpha, \beta, \gamma \geq 1$ satisfy

$$1 \leq \alpha < 2, \quad \beta \geq \alpha \quad \text{and} \quad \beta + \gamma > \frac{n}{2} + 2. \quad (26)$$

Then $u \in L^{\infty}((0, T_{\max}); L^k(\Omega))$, for all $k > k_0$.

Proof. By following the same argument of Lemma 5.2, we derive for all $k > k_0$ and some $c_{16} = c_{16}(k)$ that

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u^k &= -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 - (k-1) \chi \int_{\Omega} u^k \Delta v + ka \int_{\Omega} u^{k+\alpha-1} - kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) \\ &\leq -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + \int_{\Omega} u^{k+1} + c_{16} \int_{\Omega} |\Delta v|^{k+1} + ka \int_{\Omega} u^{k+\alpha-1} \\ &\quad - kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) \quad \text{for all } t \in (0, T_{\max}), \end{aligned} \quad (27)$$

where we used Young's inequality thanks to $\alpha \geq 1$. From the condition $\alpha < 2$, an application of Young's inequality to the fourth integral at the right-hand side of (27) gives for some $c_{17} = c_{17}(k)$

$$ka \int_{\Omega} u^{k+\alpha-1} \leq \int_{\Omega} u^{k+1} + c_{17} \quad \text{on } (0, T_{\max}). \quad (28)$$

Combining estimates (28) with bound (27), we have for all $t \in (0, T_{\max})$,

$$\frac{d}{dt} \int_{\Omega} u^k + kb \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) \leq -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + 2 \int_{\Omega} u^{k+1} + c_{16} \int_{\Omega} |\Delta v|^{k+1} + c_{17}. \quad (29)$$

As to the term $\int_{\Omega} |\Delta v|^{k+1}$ in expression (29), by exploiting in this circumstance Lemma 4.2 with $q = k+1$, we obtain for all $t \in (0, T_{\max})$ and some $c_{18} = c_{18}(k, C_P)$

$$c_{16} \int_0^t e^s \left(\int_{\Omega} |\Delta v(\cdot, s)|^{k+1} \right) ds \leq c_{16} C_P 2^k \|v_0\|_{L^{\infty}(\Omega)}^{k+1} \int_0^t e^s \left(\int_{\Omega} (u(\cdot, s))^{k+1} \right) ds + c_{18} e^t. \quad (30)$$

On the other hand, by adding $\int_{\Omega} u^k$ at both sides of estimate (29), by multiplying what obtained by e^t and considering bound (30), a subsequent integration over $(0, t)$ yields on $(0, T_{\max})$ for some $c_{19} = c_{19}(k, C_P)$ and $c_{20} = c_{20}(k, C_P)$

$$\begin{aligned} e^t \int_{\Omega} u^k - \int_{\Omega} u_0^k + kb \int_0^t e^s \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^{\gamma} \right) ds \\ \leq -\frac{4(k-1)}{k} \int_0^t e^s \left(\int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 \right) ds + c_{19} \int_0^t e^s \left(\int_{\Omega} u^{k+1} \right) ds + \int_0^t e^s \left(\int_{\Omega} u^k \right) ds + c_{20} e^t. \end{aligned} \quad (31)$$

Now let us focus on the second integral on the right-hand side of (31). Since $\int_{\Omega} u^{k+1} = \|u^{\frac{k}{2}}\|_{L^{\frac{2(k+1)}{k}}(\Omega)}^{\frac{2(k+1)}{k}}$, we proceed to apply Lemma 3.1 with $\varphi := u^{\frac{k}{2}}$ and suitable q and r . In the specific, for any $k > k_0$ and k' given in (19), it is possible to check that

$$\max \left\{ \gamma, \frac{k}{2}, \frac{n}{2} \right\} < k' < k+1. \quad (32)$$

In this way, for $n \geq 3$ letting

$$p := \frac{2n}{n-2}, \quad q := \frac{2(k+1)}{k}, \quad r := \frac{2k'}{k}$$

we can establish that $1 \leq r < q < p$ and $\frac{q}{r} < \frac{2}{r} + \frac{2}{n}$. Consequently, we deduce from (8) that for some $c_{21} = c_{21}(k, C_P, C_0)$

$$c_{19} \|u^{\frac{k}{2}}\|_{L^{\frac{2(k+1)}{k}}(\Omega)}^{\frac{\delta}{r}} \leq \frac{2(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + c_{19} \int_{\Omega} u^k + c_{21} \left(\int_{\Omega} u^{k'} \right)^{\frac{\delta}{r}} \quad \text{for all } t \in (0, T_{\max}). \quad (33)$$

Now an application of the interpolation inequality yields for all $t \in (0, T_{\max})$,

$$\begin{aligned} \left(\int_{\Omega} u^{k'} \right)^{\frac{\delta}{r}} &= \|u\|_{L^{k'}(\Omega)}^{b_2} \leq \|u\|_{L^\gamma(\Omega)}^{a_2 b_2} \|u\|_{L^{k+\beta-1}(\Omega)}^{(1-a_2)b_2} \\ &= \left(\|u\|_{L^\gamma(\Omega)}^\gamma \|u\|_{L^{k+\beta-1}(\Omega)}^{k+\beta-1} \right)^{\frac{a_2 b_2}{\gamma}} \|u\|_{L^{k+\beta-1}(\Omega)}^{\left[1-a_2-\frac{a_2(k+\beta-1)}{\gamma}\right] b_2}, \end{aligned}$$

where

$$b_2 = b_2(q) := \frac{k' \delta(q)}{r} = \frac{k' \delta}{r}, \quad a_2 := \frac{\frac{1}{k'} - \frac{1}{k+\beta-1}}{\frac{1}{\gamma} - \frac{1}{k+\beta-1}} \in (0, 1).$$

(A comparison between the couple (a_2, b_2) above and (a_1, b_1) in (23) shows that $a_1 = a_2$, whereas b_i , $i = 1, 2$ depends on q .) From straightforward calculations and the condition (26), we note that

$$\left[1 - a_2 - \frac{a_2(k+\beta-1)}{\gamma} \right] b_2 = 0 \quad \text{and} \quad \frac{a_2 b_2}{\gamma} < 1.$$

Subsequently, we can exploit the Young inequality entailing for some $c_{22} = c_{22}(k, C_P, C_0)$

$$c_{21} \left(\int_{\Omega} u^{k'} \right)^{\frac{\delta}{r}} \leq c_{21} \left(\|u\|_{L^\gamma(\Omega)}^\gamma \|u\|_{L^{k+\beta-1}(\Omega)}^{k+\beta-1} \right)^{\frac{a_2 b_2}{\gamma}} \leq k b \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^\gamma \right) + c_{22} \quad \text{on } (0, T_{\max}).$$

This in conjunction with (33) implies that $t \in (0, T_{\max})$,

$$c_{19} \int_{\Omega} u^{k+1} \leq \frac{2(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 + c_{19} \int_{\Omega} u^k + k b \left(\int_{\Omega} u^{k+\beta-1} \right) \left(\int_{\Omega} u^\gamma \right) + c_{22}. \quad (34)$$

For $n \in \{1, 2\}$ the argument proceeds similarly, provided that condition (32) is replaced by

$$\max \left\{ \gamma, \frac{k}{2} \right\} < k' < k + \alpha - 1.$$

For k' in this range, some computations show that q, r satisfy $1 \leq r < q$ and $\frac{q}{r} < \frac{2}{r} + \frac{2}{n}$. Therefore, Lemma 3.1 can also be applied in these cases, yielding to inequality (34) up to some constants.

By rearranging bound (31) by virtue of estimates (34) and (25), it is provided for some $c_{23} = c_{23}(k)$ and $c_{24} = c_{24}(k, C_P, C_0, \hat{c})$

$$\int_{\Omega} u^k \leq c_{23} + c_{24} e^{-t} \quad \text{on } (0, T_{\max}),$$

which gives the statement. $\square$

5.2 | The Case $f(u) = -au^\alpha + bu^\beta \int_{\Omega} u^\gamma$

Lemma 5.4. Assume $a, b > 0$ and $\alpha, \beta, \gamma \geq 1$ satisfy either

$$\alpha > 2, \quad \beta + \gamma < \alpha, \quad (35)$$

or

$$\alpha > 2, \quad \beta + \gamma = \alpha \quad \text{and} \quad a > b|\Omega|. \quad (36)$$

Then $u \in L^\infty((0, T_{\max}); L^k(\Omega))$ for all $k > k_0$.

Proof. By reasoning as in the proof of Lemma 5.2, an application of the Young inequality (recall $\alpha > 1$) yields for all $\epsilon > 0$, $k > k_0$ and some $c_{25} = c_{25}(k)$

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u^k &= -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{\frac{k}{2}}|^2 - (k-1)\chi \int_{\Omega} u^k \Delta v + kb \int_{\Omega} u^{k+\beta-1} \int_{\Omega} u^\gamma - ka \int_{\Omega} u^{k+\alpha-1} \\ &\leq (\epsilon - ka) \int_{\Omega} u^{k+\alpha-1} + c_{25} \int_{\Omega} |\Delta v|^{\frac{k+\alpha-1}{\alpha-1}} + kb \int_{\Omega} u^{k+\beta-1} \int_{\Omega} u^\gamma \quad \text{for all } t \in (0, T_{\max}). \end{aligned} \quad (37)$$

(Throughout this proof we adopt the convention concerning the role of ϵ stated at the beginning of Section 3.)

Adding $\int_{\Omega} u^k$ to both sides of estimate (37) and multiplying the resulting expression by e^t , integration over $(0, t)$ then yields, for all $t \in (0, T_{\max})$,

$$\begin{aligned} e^t \int_{\Omega} u^k - \int_{\Omega} u_0^k &\leq (\epsilon - ka) \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds + \int_0^t e^s \left(\int_{\Omega} u^k \right) ds \\ &\quad + c_{25} \int_0^t e^s \left(\int_{\Omega} |\Delta v|^{\frac{k+\alpha-1}{\alpha-1}} \right) ds + kb \int_0^t e^s \left(\int_{\Omega} u^{k+\beta-1} \int_{\Omega} u^\gamma \right) ds. \end{aligned} \quad (38)$$

By applying Lemma 4.2 with $q = \frac{k+\alpha-1}{\alpha-1}$ and Young's inequality since $\alpha > 2$, we obtain for some $c_{26} = c_{26}(k, C_P)$ and $c_{27} = c_{27}(k, C_P)$

$$\begin{aligned} c_{25} \int_0^t e^s \left(\int_{\Omega} |\Delta v|^{\frac{k+\alpha-1}{\alpha-1}} \right) ds &\leq c_{26} e^t + c_{25} C_P 2^{\frac{k}{\alpha-1}} \|v_0\|_{L^\infty(\Omega)}^{\frac{k+\alpha-1}{\alpha-1}} \int_0^t e^s \left(\int_{\Omega} u^{\frac{k+\alpha-1}{\alpha-1}} \right) ds \\ &\leq c_{27} e^t + \epsilon \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds \quad \text{on } (0, T_{\max}). \end{aligned} \quad (39)$$

A further application of Young's inequality yields for some $c_{28} = c_{28}(k)$

$$\int_0^t e^s \left(\int_{\Omega} u^k \right) ds \leq \epsilon \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds + c_{28} e^t \quad \text{for all } t \in (0, T_{\max}). \quad (40)$$

By virtue of $\beta < \beta + \gamma < \alpha$ and $k > k_0$, a combination of Hölder's and Young's inequalities leads for all $t \in (0, T_{\max})$ and some $c_{29} = c_{29}(k)$ to

$$kb \int_{\Omega} u^{k+\beta-1} \int_{\Omega} u^\gamma \leq \begin{cases} kb|\Omega|^{\frac{k+2\alpha-\beta-\gamma-1}{k+\alpha-1}} \left(\int_{\Omega} u^{k+\alpha-1} \right)^{\frac{k+\beta+\gamma-1}{k+\alpha-1}} \leq \epsilon \int_{\Omega} u^{k+\alpha-1} + c_{29} & \text{if } \beta + \gamma < \alpha, \\ kb|\Omega|^{\frac{k+2\alpha-\beta-\gamma-1}{k+\alpha-1}} \left(\int_{\Omega} u^{k+\alpha-1} \right)^{\frac{k+\beta+\gamma-1}{k+\alpha-1}} = kb|\Omega| \int_{\Omega} u^{k+\alpha-1} & \text{if } \beta + \gamma = \alpha. \end{cases} \quad (41)$$

By plugging estimates (39), (40) and (41) into (38), we derive for some $c_{30} = c_{30}(k)$ and $c_{31} = c_{31}(k, C_P)$ on $(0, T_{\max})$

$$e^t \int_{\Omega} u^k \leq \begin{cases} c_{30} + c_{31} e^t + (\epsilon - ka) \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds & \text{if } \beta + \gamma < \alpha \\ c_{30} + c_{31} e^t + (\epsilon + kb|\Omega| - ka) \int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds & \text{if } \beta + \gamma = \alpha, \end{cases}$$

which implies, by assumption (35) and choosing $0 < \epsilon \leq ka$ (for instance exactly the limit value ka), that

$$e^t \int_{\Omega} u^k \leq c_{30} + c_{31} e^t \quad \text{on } (0, T_{\max}).$$

On the other hand, if $a > b|\Omega|$ as in hypothesis (36) and by virtue of the arbitrariness of ϵ , we can make $\epsilon + kb|\Omega| - ka \leq 0$ and consequently neglect the term associated to $\int_0^t e^s \left(\int_{\Omega} u^{k+\alpha-1} \right) ds$, leading to the above inequality up to a constant. Henceforth, in both cases we conclude that $u \in L^\infty((0, T_{\max}); L^k(\Omega))$ for all $k > k_0$. $\square$

In the following lemma, we set $\alpha = 2$ and $\beta + \gamma = \alpha$. Under the conditions $\beta \geq 1$ and $\gamma \geq 1$, this uniquely determines $\beta = \gamma = 1$.

Lemma 5.5. Assume $\alpha = 2$, and $\beta = \gamma = 1$, $k > k_0$. Moreover, let $C_P = C_P(n, k + 1)$ being C_P the constant in Lemma 4.2. If

$$a > b|\Omega| + \frac{(k-1)}{k}(C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)}, \quad (42)$$

then the same statement of Lemma 5.4 holds.

Proof. By taking $\alpha = 2$ and $\beta = \gamma = 1$ in bounds (37) and (41), we obtain for all $t \in (0, T_{\max})$

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u^k &\leq \left(\frac{k(k-1)}{(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} - ka \right) \int_{\Omega} u^{k+1} \\ &\quad + \frac{(k-1)\chi}{k+1} (C_P 2^k)^{-\frac{k}{k+1}} \|v_0\|_{L^\infty(\Omega)}^{-k} \int_{\Omega} |\Delta v|^{k+1} + kb \int_{\Omega} u^k \int_{\Omega} u \\ &\leq \left(\frac{k(k-1)}{(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} + kb|\Omega| - ka \right) \int_{\Omega} u^{k+1} \\ &\quad + \frac{(k-1)\chi}{k+1} (C_P 2^k)^{-\frac{k}{k+1}} \|v_0\|_{L^\infty(\Omega)}^{-k} \int_{\Omega} |\Delta v|^{k+1}. \end{aligned}$$

Following the argument used in the proof of Lemma 5.4, we obtain for all $t \in (0, T_{\max})$

$$\begin{aligned} e^t \int_{\Omega} u^k - \int_{\Omega} u_0^k &\leq \left(\frac{k(k-1)}{(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} + kb|\Omega| - ka \right) \int_0^t e^s \left(\int_{\Omega} u^{k+1} \right) ds \\ &\quad + \frac{(k-1)\chi}{k+1} (C_P 2^k)^{-\frac{k}{k+1}} \|v_0\|_{L^\infty(\Omega)}^{-k} \int_0^t e^s \left(\int_{\Omega} |\Delta v|^{k+1} \right) ds + \int_0^t e^s \left(\int_{\Omega} u^k \right) ds. \end{aligned} \quad (43)$$

Now, an application of Lemma 4.2 with $q = k + 1$ entails on $(0, T_{\max})$

$$\begin{aligned} &\frac{(k-1)\chi}{k+1} (C_P 2^k)^{-\frac{k}{k+1}} \|v_0\|_{L^\infty(\Omega)}^{-k} \int_0^t e^s \left(\int_{\Omega} |\Delta v|^{k+1} \right) ds \\ &\leq c_{26} e^t + \frac{(k-1)}{k+1} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} \int_0^t e^s \left(\int_{\Omega} u^{k+1} \right) ds. \end{aligned} \quad (44)$$

By plugging estimates (44) and (40) into (43), we have on $(0, T_{\max})$

$$\begin{aligned} e^t \int_{\Omega} u^k &\leq c_{30} + c_{31} e^t + k \left[\frac{(k-1)}{(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} + b|\Omega| + \epsilon \right. \\ &\quad \left. + \frac{(k-1)}{k(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} - a \right] \int_0^t e^s \left(\int_{\Omega} u^{k+1} \right) ds. \end{aligned}$$

Finally, the assumption (42) made on a and the arbitrariness of ϵ allow us to make

$$\frac{(k-1)}{(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} + b|\Omega| + \epsilon + \frac{(k-1)}{k(k+1)} (C_P 2^k)^{\frac{1}{k+1}} \|\chi v_0\|_{L^\infty(\Omega)} - a \leq 0$$

and consequently neglect the term associated to $\int_0^t e^s \left(\int_{\Omega} u^{k+1} \right) ds$, achieving the claim. $\square$

6 | Proof of Theorem 2.1

The claims concerning 1(a) and 1(b) for the case where $f(u) = au^\alpha - bu^\beta \int_{\Omega} u^\gamma$, are consequence of Lemmas 5.2 and 5.3, ensuring the uniform-in-time estimate of u in $L^k(\Omega)$ for sufficiently large k , combined with Lemma 4.3, ruling out the possibility that $T_{\max}$ is finite.

For $f(u) = -au^\alpha + bu^\beta \int_{\Omega} u^r$ the cases 2(a) and 2(b) follow by reasoning as before, invoking Lemmas 5.4 and 4.3. Finally, let $\mathcal{K} = \mathcal{K}(n) = \frac{n+1}{n+2}(C_P(n, n+3) 2^{n+2})^{\frac{1}{n+3}}$, being C_P introduced in Lemma 4.2. Under the assumption on a in case 2(c), by continuity arguments also involving the constant C_P itself, there is $k > n + 2$ for which (42) holds, so that by exploiting Lemma 5.5 we conclude.

Author Contributions

Rafael Díaz Fuentes: investigation, writing – original draft. **Fatma Gamze Düzgün:** Silvia Frassu: investigation, writing – original draft. **Giuseppe Viglialoro:** investigation, writing – original draft.

Acknowledgments

The authors are members of the Gruppo Nazionale per l'Analisi Matematica, la Probabilità e le loro Applicazioni (GNAMPA) of the Istituto Nazionale di Alta Matematica (INdAM), and are partially supported by the research project *Partial Differential Equations and their role in understanding natural phenomena* (2023, CUP F23C25000080007), funded by [Fondazione di Sardegna](#). R.D.F. acknowledges financial support by PNRR e.INS Ecosystem of Innovation for Next Generation Sardinia (CUP F53C22000430001, codice MUR ECS0000038). S.F. and G.V. are also supported by MIUR (Italian Ministry of Education, University and Research) Prin 2022 *Non-linear differential problems with applications to real phenomena* (Grant Number: 2022ZXZTN2). R.D.F., S.F. and G.V. acknowledge GNAMPA-INdAM Project *Modelli di reazione-diffusione-trasporto: dall'analisi alle applicazioni* (CUP E53C25002010001). F.G.D. is partially supported by GNAMPA-INdAM Project *Classificazione delle soluzioni di equazioni evolutive non locali* (CUP E53C25002010001). Open access publishing facilitated by Università degli Studi di Cagliari, as part of the Wiley - CRUI-CARE agreement.

Funding

This work was supported by the Gruppo Nazionale per l'Analisi Matematica, la Probabilità e le loro Applicazioni (Grant No. CUP E53C25002010001), the Fondazione di Sardegna (Grant No. CUP F23C25000080007), the Ministero dell'Università e della Ricerca (Grant No. 2022ZXZTN2), and the PNRR e.INS Ecosystem of Innovation for Next Generation Sardinia (Grant No. CUP F53C22000430001).

Conflicts of Interest

The authors declare no conflicts of interest.

Data Availability Statement

This manuscript did not involve the generation or analysis of data requiring deposition or sharing.

References

1. E. F. Keller and L. A. Segel, "Initiation of Slime Mold Aggregation Viewed as an Instability," *Journal of Theoretical Biology* 26, no. 3 (1970): 399–415.
2. E. F. Keller and L. A. Segel, "Model for Chemotaxis," *Journal of Theoretical Biology* 30, no. 2 (1971): 225–234.
3. N. Bellomo, A. Bellouquid, Y. Tao, and M. Winkler, "Toward a Mathematical Theory of Keller–Segel Models of Pattern Formation in Biological Tissues," *Mathematical Models & Methods in Applied Sciences* 25, no. 9 (2015): 1663–1763.
4. W. Jäger and S. Luckhaus, "On Explosions of Solutions to a System of Partial Differential Equations Modelling Chemotaxis," *Transactions of the American Mathematical Society* 329, no. 2 (1992): 819–824.
5. T. Nagai, "Blowup of Nonradial Solutions to Parabolic–Elliptic Systems Modeling Chemotaxis in Two-Dimensional Domains," *Journal of Inequalities and Applications* 6, no. 1 (2001): 37–55.
6. J. I. Tello and M. Winkler, "A Chemotaxis System With Logistic Source," *Communications in Partial Differential Equations* 32, no. 6 (2007): 849–877.
7. M. Winkler, "Boundedness in the Higher-Dimensional Parabolic-Parabolic Chemotaxis System With Logistic Source," *Communications in Partial Differential Equations* 35, no. 8 (2010): 1516–1537.
8. M. Fuest, "Approaching Optimality in Blow-Up Results for Keller–Segel Systems With Logistic-Type Dampening," *NoDEA Nonlinear Differential Equations and Applications* 28, no. 2 (2021): 16.
9. M. Winkler, "Blow-Up in a Higher-Dimensional Chemotaxis System Despite Logistic Growth Restriction," *Journal of Mathematical Analysis and Applications* 384, no. 2 (2011): 261–272.
10. M. Winkler, "Finite-Time Blow-Up in Low-Dimensional Keller–Segel Systems With Logistic-Type Degradation," *Zeitschrift für Angewandte Mathematik und Physik* 69, no. 2 (2018): 40.

11. S. Bian, L. Chen, and E. A. Latos, "Nonlocal Nonlinear Reaction Preventing Blow-Up in Supercritical Case of Chemotaxis System," *Nonlinear Analysis* 176 (2018): 178–191.
12. M. Negreanu and J. I. Tello, "On a Competitive System Under Chemotactic Effects With Non-Local Terms," *Nonlinearity* 26, no. 4 (2013): 1083–1103.
13. M. Negreanu, J. I. Tello, and A. Vargas, "On a Fully Parabolic Chemotaxis System With Nonlocal Growth Term," *Nonlinear Analysis* 213 (2021): 112518.
14. X. Tao and Z. B. Fang, "Global Boundedness of Solutions to a Quasilinear Chemotaxis System With Nonlocal Nonlinear Reaction," *Applied Mathematics and Optimization* 87, no. 2 (2023): 20.
15. S. Bian, L. Chen, and E. A. Latos, "Chemotaxis Model With Nonlocal Nonlinear Reaction in the Whole Space," *Discrete and Continuous Dynamical Systems* 38, no. 10 (2018): 5067–5083.
16. J. A. Carrillo and J. Wang, "Uniform in Time L^∞ -Estimates for Nonlinear Aggregation-Diffusion Equations," *Acta Applicandae Mathematicae* 164 (2019): 1–19.
17. L. Chen and J. Wang, "Exact Criterion for Global Existence and Blow up to a Degenerate Keller–Segel System," *Documenta Mathematica* 19 (2014): 103–120.
18. T. Li and G. Viglialoro, "Boundedness for a Nonlocal Reaction Chemotaxis Model Even in the Attraction-Dominated Regime," *Differential and Integral Equations* 34, no. 5–6 (2021): 315–336.
19. P. Quittner and P. Souplet, *Superlinear Parabolic Problems. Blow-Up, Global Existence and Steady States*. Birkhäuser Advanced Texts: Basler Lehrbücher (Birkhäuser/Springer, 2019).
20. K. Baghaei, S. Frassu, Y. Tanaka, and G. Viglialoro, "To What Extent Does the Consideration of Positive Total Flux Influence the Dynamics of Keller–Segel-Type Models?," *Journal of Differential Equations* 452 (2026): 113808.
21. R. Díaz Fuentes, S. Frassu, and G. Viglialoro, "Dissipation Through Combinations of Nonlocal and Gradient Nonlinearities in Chemotaxis Models," *Acta Applicandae Mathematicae* 195 (2025): 10.
22. S. Ishida, J. Lankeit, and G. Viglialoro, "A Keller–Segel Type Taxis Model With Ecological Interpretation and Boundedness due to Gradient Nonlinearities," *Discrete and Continuous Dynamical Systems - B* 29, no. 9 (2024): 3955–3969.
23. T. Li, D. Acosta-Soba, A. Columbu, and G. Viglialoro, "Dissipative Gradient Nonlinearities Prevent -Formations in Local and Non-local Attraction–Repulsion Chemotaxis Models," *Studies in Applied Mathematics* 154, no. 2 (2025): e70018.
24. K. Baghaei, "Boundedness of Classical Solutions to a Chemotaxis Consumption System With Signal Dependent Motility and Logistic Source," *Comptes Rendus Mathematique* 361 (2023): 1641–1652.
25. K. Baghaei and A. Khelghati, "Global Existence and Boundedness of Classical Solutions for a Chemotaxis Model With Consumption of Chemoattractant and Logistic Source," *Mathematical Methods in the Applied Sciences* 40, no. 10 (2017): 3799–3807.
26. J. Lankeit and Y. Wang, "Global Existence, Boundedness and Stabilization in a High-Dimensional Chemotaxis System With Consumption," *Discrete and Continuous Dynamical Systems* 37, no. 12 (2017): 6099–6121.
27. A. Columbu, "Boundedness Criteria for a Chemotaxis Consumption Model With Gradient Nonlinearities," *Journal of Mathematical Analysis and Applications* 546, no. 1 (2025): 129226.
28. Y. Tao, "Boundedness in a Chemotaxis Model With Oxygen Consumption by Bacteria," *Journal of Mathematical Analysis and Applications* 381, no. 2 (2011): 521–529.
29. K. Baghaei and A. Khelghati, "Boundedness of Classical Solutions for a Chemotaxis Model With Consumption of Chemoattractant," *Comptes Rendus Mathematique* 355, no. 6 (2017): 633–639.
30. Y. Chiyo, F. G. Düzgün, S. Frassu, and G. Viglialoro, "Boundedness Through Nonlocal Dampening Effects in a Fully Parabolic Chemotaxis Model With Sub and Superquadratic Growth," *Applied Mathematics and Optimization* 89, no. 1 (2024): 1–30.
31. M. Hieber and J. Prüss, "Heat Kernels and Maximal L^p - L^q Estimates for Parabolic Evolution Equations," *Communications in Partial Differential Equations* 22, no. 9–10 (1997): 1647–1669.
32. Y. Giga and H. Sohr, "Abstract L^p Estimates for the Cauchy Problem With Applications to the Navier–Stokes Equations in Exterior Domains," *Journal of Functional Analysis* 102, no. 1 (1991): 72–94.
33. S. Ishida and T. Yokota, "Global Existence of Weak Solutions to Quasilinear Degenerate Keller–Segel Systems of Parabolic-Parabolic Type With Small Data," *Journal of Differential Equations* 252, no. 3 (2012): 2469–2491.
34. S. Ishida and T. Yokota, "Boundedness in a Quasilinear Fully Parabolic Keller-Segel System via Maximal Sobolev Regularity," *Discrete and Continuous Dynamical Systems - Series S* 13, no. 2 (2020): 212–232.
35. T. Senba and T. Suzuki, "A Quasi-Linear Parabolic System of Chemotaxis," *Abstract and Applied Analysis* 2006 (2006): 21.
36. M. Winkler, "Aggregation vs. Global Diffusive Behavior in the Higher-Dimensional Keller–Segel Model," *Journal of Differential Equations* 248, no. 12 (2010): 2889–2905.

37. Y. Tao and M. Winkler, “Boundedness in a Quasilinear Parabolic-Parabolic Keller–Segel System With Subcritical Sensitivity,” *Journal of Differential Equations* 252, no. 1 (2012): 692–715.
38. H. Brezis, *Functional Analysis, Sobolev Spaces and Partial Differential Equations*, Universitext (Springer, 2011).