

SYMMETRY AND REGULARITY OF AN OPTIMIZATION PROBLEM RELATED TO A NONLINEAR BVP

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ABSTRACT. We consider the functional

$$f \mapsto \int_{\Omega} \left(\frac{q+1}{2} |Du_f|^2 - u_f |u_f|^q f \right) dx,$$

where u_f is the unique nontrivial weak solution of the boundary-value problem

$$-\Delta u = f|u|^q \quad \text{in } \Omega, \quad u|_{\partial\Omega} = 0,$$

where $\Omega \subset \mathbb{R}^n$ is a bounded smooth domain. We prove a result of Steiner symmetry preservation and, if $n = 2$, we show the regularity of the level sets of minimizers.

1. INTRODUCTION

Let Ω be a bounded domain $\Omega \subset \mathbb{R}^n$ with smooth boundary. We consider the Dirichlet problem

$$\begin{aligned} -\Delta u &= f|u|^q \quad \text{in } \Omega, \\ u|_{\partial\Omega} &= 0, \end{aligned} \tag{1.1}$$

where $0 \leq q < \min\{1, 4/n\}$ and f is a nonnegative bounded function non identically zero. We consider nontrivial solutions of (1.1) in $H_0^1(\Omega)$. The equation (1.1) is the Euler-Lagrange equation of the integral functional

$$v \mapsto \int_{\Omega} \left(\frac{q+1}{2} |Dv|^2 - v|v|^q f \right) dx, \quad v \in H_0^1(\Omega).$$

By using a standard compactness argument, it can be proved that there exists a nontrivial minimizer of the above functional. This minimizer is a nontrivial solution of (1.1).

From the maximum principle, every nontrivial solution of (1.1) is positive. Then, by [5, Theorem 3.2] the uniqueness of problem (1.1) follows. To underscore the dependence on f of the solution of (1.1), we denote it by u_f . Moreover, $u_f \in W^{2,2}(\Omega) \cap C^{1,\alpha}(\bar{\Omega})$ for all α , $0 < \alpha < 1$ (see [8, 10]).

Let f_0 be a fixed bounded nonnegative function. We study the problem

$$\inf_{v \in H_0^1(\Omega), f \in \mathfrak{F}(f_0)} \int_{\Omega} \left(\frac{q+1}{2} |Dv|^2 - v|v|^q f \right) dx, \tag{1.2}$$

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where, denoting by $|A|$ the Lebesgue measure of a set A ,

$$\mathfrak{F}(f_0) = \{f \in L^\infty(\Omega) : |\{f \geq c\}| = |\{f_0 \geq c\}| \forall c \in \mathbb{R}\}; \quad (1.3)$$

here $\mathfrak{F}(f_0)$ is called class of rearrangements of f_0 (see [9]).

Problems of this kind are not new; see for example [1, 2, 3, 5]. From the results in [5] it follows that (1.2) has a minimum and a representation formula. Let

$$E(f) = \inf_{v \in H_0^1(\Omega)} \int_{\Omega} \left(\frac{q+1}{2} |Dv|^2 - v|v|^q f \right) dx. \quad (1.4)$$

Renaming q' the constant q and putting $p = 2$ and $q' = q + 1$ in [5] we have

$$E(f) = \frac{q' - 2}{2} I(f),$$

where

$$I(f) = \sup_{H_0^1(\Omega)} \frac{q'}{2 - q'} \int_{\Omega} \left(\frac{2}{q'} f |v|^{q'} - |Dv|^2 \right) dx$$

is defined in the same paper. By [5, Theorem 2.2] it follows that there exist minimizers of $E(f)$ and that, if \bar{f} is a minimizer, there exists an increasing function ϕ such that

$$\bar{f} = \phi(u_{\bar{f}}). \quad (1.5)$$

We denote by $\text{supp } f$ the support of f , and we call a level set of f the set $\{x \in \Omega : f(x) > c\}$, for some constant c .

In Section 2, we consider a Steiner symmetric domain Ω and f_0 bounded and nonnegative, such that $|\text{supp } f_0| < |\Omega|$. Under these assumptions, we prove that the level sets of the minimizer \bar{f} are Steiner symmetric with respect to the same hyperplane of Ω . As a consequence, we have exactly one optimizer when Ω is a ball.

Chanillo, Kenig and To [4] studied the regularity of the minimizers to the problem

$$\lambda(\alpha, A) = \inf_{u \in H_0^1(\Omega), \|u\|_2, |D| = A} \int_{\Omega} |Du|^2 dx + \alpha \int_D u^2 dx,$$

where $\Omega \subset \mathbb{R}^2$ is a bounded domain, $0 < A < |\Omega|$ and $\alpha > 0$. In particular they prove that, if D is a minimizer, then ∂D is analytic.

In Section 3, following the ideas in [4], we give our main result. We restrict our attention to $\Omega \subset \mathbb{R}^2$. Let $b_1, \dots, b_m > 0$ and $0 < a_1 < \dots < a_m < |\Omega|$, $m \geq 2$, be fixed. Consider $f_0 = b_1 \chi_{G_1} + \dots + b_m \chi_{G_m}$, where $|G_i| = a_i$ for all i , $G_i \subset G_{i+1}$, $i = 1, \dots, m-1$.

We call η the minimum in (1.2); i.e.,

$$\eta = \int_{\Omega} \left(\frac{q+1}{2} |Du_{\bar{f}}|^2 - u_{\bar{f}} |u_{\bar{f}}|^q \bar{f} \right) dx,$$

where \bar{f} and $u_{\bar{f}}$ are, respectively, the minimizing function and the corresponding solution of (1.1).

In this case (1.5) becomes

$$\bar{f} = \sum_{i=1}^m b_i \chi_{D_i},$$

where

$$D_1 = \{u_{\bar{f}} > c_1\}, \quad D_2 = \{u_{\bar{f}} > c_2\}, \quad \dots, \quad D_m = \{u_{\bar{f}} > c_m\},$$

for suitable constants $c_1 > c_2 > \dots > c_m > 0$.

We show regularity of ∂D_i for each i proving that $|Du_{\bar{f}}| > 0$ in ∂D_i . Following the method used in [4], we consider

$$E(s, t) = \int_{\Omega} \left(\frac{q+1}{2} |Du_{\bar{f}} + sDv|^2 - (u_{\bar{f}} + sv)|u_{\bar{f}} + sv|^q \bar{f}_t \right) dx - \eta, \tag{1.6}$$

where $v \in H_0^1(\Omega)$, \bar{f}_t is a family of functions such that $\bar{f}_t \in \mathfrak{F}(f_0)$ with $\bar{f}_0 = \bar{f}$, and $s \in \mathbb{R}$. We have

$$E(s, t) \geq E(0, 0) = 0 \quad \forall s, t.$$

Therefore, $(s, t) = (0, 0)$ is a minimum point; it follows that

$$\begin{vmatrix} \frac{\partial^2 E}{\partial s^2}(0, 0) & \frac{\partial^2 E}{\partial s \partial t}(0, 0) \\ \frac{\partial^2 E}{\partial t \partial s}(0, 0) & \frac{\partial^2 E}{\partial t^2}(0, 0) \end{vmatrix} \geq 0. \tag{1.7}$$

Expanding (1.7) in detail and using some lemmas from [4] we prove that the boundaries of level sets of \bar{f} are regular.

Theorem 1.1. *Let $\Omega \subset \mathbb{R}^2$, $f_0 = \sum_{i=1}^m b_i \chi_{G_i}$ with $m \geq 2$, and $\bar{f} = \sum_{i=1}^m b_i \chi_{D_i}$ a minimizer of (1.2). Then $|Du_{\bar{f}}| > 0$ on ∂D_i , $i = 1, \dots, m$.*

2. SYMMETRY

In this section we consider Steiner symmetric domains. We prove that, under suitable conditions on f_0 in (1.3), minimizers inherit Steiner symmetry.

Definition 2.1. Let $P \subset \mathbb{R}^n$ be a hyperplane. We say that a set $A \subset \mathbb{R}^n$ is *Steiner symmetric* relative to the hyperplane P if for every straight line L perpendicular to P , the set $A \cap L$ is either empty or a symmetric segment with respect to P .

To prove the symmetry, we need [6, Theorem 3.6 and Corollary 3.9], that, for more convenience for the reader, we state here. These results are related to the classical paper [7].

Theorem 2.2. *Let $\Omega \subset \mathbb{R}^n$ be bounded, connected and Steiner symmetric relative to the hyperplane P . Assume that $u : \bar{\Omega} \rightarrow \mathbb{R}$ has the following properties:*

- $u \in C(\bar{\Omega}) \cap C^1(\Omega)$, $u > 0$ in Ω , $u|_{\partial\Omega} = 0$;
- for all $\phi \in C_0^\infty(\Omega)$,

$$\int_{\Omega} Du \cdot D\phi \, dx = \int_{\Omega} \phi F(u) \, dx,$$

where F has a decomposition $F = F_1 + F_2$ such that $F_1 : [0, \infty) \rightarrow \mathbb{R}$ is locally Lipschitz continuous, while $F_2 : [0, \infty) \rightarrow \mathbb{R}$ is non-decreasing and identically 0 on $[0, \epsilon]$ for some $\epsilon > 0$.

Then u is symmetric with respect to P and $\frac{\partial u}{\partial \mathbf{v}}(x) < 0$, where \mathbf{v} is a unit vector orthogonal to P and x belongs to the part of Ω that lies in the halfspace (with origin in P) in which \mathbf{v} points.

Theorem 2.3. *Let Ω be Steiner symmetric and f_0 a bounded nonnegative function. If $|\text{supp } f_0| < |\Omega|$ and $\bar{f} \in \mathfrak{F}(f_0)$ is a minimizer of (1.2), then the level sets of \bar{f} are Steiner symmetric with respect to the same hyperplane of Ω .*

Proof. Let $u = u_{\bar{f}}$ be the solution of (1.2). Then $u \in C^0(\bar{\Omega}) \cap C^1(\Omega)$ and satisfies

$$\int_{\Omega} Du \cdot D\psi \, dx = \int_{\Omega} \psi f u^q \, dx \quad \forall \psi \in C_0^\infty(\Omega).$$

Since $u > 0$ and since (from (1.5)) $\bar{f} = \phi(u)$ with ϕ increasing function, it follows that $\phi(u) \equiv 0$ on $\{x \in \Omega : u(x) < d\}$ for some positive constant d . Then we have $u^q \bar{f} = F_1(u) + F_2(u)$ with $F_1(u) \equiv 0$ and $F_2(u) = \phi(u)u^q$. From Theorem 2.1 and $\bar{f} = \phi(u)$ we have the assertion. \square

Remark 2.4. By this theorem, if Ω is an open ball and $|\text{supp } f_0| < |\Omega|$, then \bar{f} is radially symmetric and decreasing.

3. REGULARITY OF THE FREE BOUNDARIES

In this section we prove the following result.

Theorem 3.1. *Let $\Omega \subset \mathbb{R}^2$, $f_0 = \sum_{i=1}^m b_i \chi_{G_i}$ with $m \geq 2$, and $\bar{f} = \sum_{i=1}^m b_i \chi_{D_i}$ a minimizer of (1.2). Then $|Du_{\bar{f}}| > 0$ on ∂D_i , $i = 1, \dots, m$.*

We use the notation introduced in Section 1. Without loss of generality we can assume $m = 3$; the general case easily follows. Let $\bar{f} = b_1 \chi_{D_1} + b_2 \chi_{D_2} + b_3 \chi_{D_3}$. We will prove $|Du_{\bar{f}}| > 0$ in ∂D_2 ; we omit the proof for D_1 and D_3 because it is similar. We define the family \bar{f}_t by replacing only the set D_2 by a family of domains $D_2(t)$.

First of all, we explain how to define the family $D_2(t)$. In the sequel we use the notation introduced in [4], reorganized according to our needs.

We call a curve $\gamma : [a, b] \rightarrow \mathbb{R}$, $-\infty < a < b < \infty$, regular if:

- (i) it is simple, that is: if $a \leq x < y \leq b$ and $x \neq a$ or $y \neq b$, then $\gamma(x) \neq \gamma(y)$;
- (ii) $\|\gamma\|_{C^2(a,b)}$ is finite;
- (iii) $|\gamma'|$ is uniformly bounded away from zero.

If, in addition, $\gamma(a) = \gamma(b)$, we say that the curve is closed and regular. If the domain of γ is (a, b) we say that γ is regular (respectively, closed and regular) if the continuous extension of γ to $[a, b]$ is regular (respectively, closed and regular).

Now, we introduce the notation

$$\mathcal{F} := \partial D_2; \quad \mathcal{F}^* := \mathcal{F} \cap \{|Du_{\bar{f}}| > 0\}.$$

Let $J = \cup_{k=1}^p J_k$ be a finite union of open bounded intervals $J_k \subset \mathbb{R}$, $\gamma = (\gamma_1, \gamma_2) : J \rightarrow \mathcal{F}^*$ a simple curve which is regular on each interval J_k and $\gamma(J) \subset \mathcal{F}^*$. We suppose that $\text{dist}(\gamma(J_k), \gamma(J_h)) > 0$ for $1 \leq h \neq k \leq p$. Assume also that $|\gamma'| \geq \theta$ on J . For each $\xi \in J$, we denote by $\mathbf{N}(\xi) = (N_1(\xi), N_2(\xi))$ the outward unit normal with respect to D_2 at $\gamma(\xi)$. We also define the tangent vector to γ $\mathbf{N}^\perp(\xi) = (-N_2(\xi), N_1(\xi))$, and \mathbf{N}' the first derivative of \mathbf{N} .

Reversing the direction of γ if necessary, we will assume, without loss of generality, that γ' and \mathbf{N}^\perp have the same direction; i.e., $\angle \gamma', \mathbf{N}^\perp = |\gamma'|$. We observe that, because γ is C^2 and simple on \bar{J}_k , for each k there exists $\beta_k > 0$ such that the function

$$\phi_k : J_k \times [-\beta_k, \beta_k] \rightarrow \mathbb{R}^2, \quad (\xi, \beta) \mapsto (x_1, x_2) = \phi_k(\xi, \beta) = \gamma(\xi) + \beta \mathbf{N}(\xi)$$

is injective.

Because $\text{dist}(\gamma(J_k), \gamma(J_h)) > 0$ for all $h \neq k$, we can find a number $\beta_0 > 0$ and we can paste together the functions ϕ_k to obtain a function ϕ injective on $J \times [-\beta_0, \beta_0]$. Choose β_0 such that $\text{dist}(\phi(J \times [-\beta_0, \beta_0]), \partial D_1) > 0$ and $\text{dist}(\phi(J \times [-\beta_0, \beta_0]), \partial D_3) > 0$.

Now, we define

$$K = D_2 \setminus \phi(J \times (-\beta_0, 0]);$$

for $t \in (-t_0, t_0)$ we define

$$D_2(t) = K \cup \{\phi(\xi, \beta) : \xi \in J, \beta < g(\xi, t)\}, \tag{3.1}$$

where $g : J \times (-t_0, t_0) \rightarrow \mathbb{R}$, $t_0 > 0$, is a function such that

$$g(\xi, t), g_t(\xi, t), g_{tt}(\xi, t) \in C(\bar{J}) \quad \forall t \in (-t_0, t_0) \tag{3.2}$$

and

$$g(\xi, 0) \equiv 0 \quad \forall \xi \in J. \tag{3.3}$$

We observe that $D_2(0) = D_2$. Next we compute the measure of $D_2(t)$. Put $A(t) = |D_2(t)|$ and $A = |D_2(0)| = |D_2|$; we have

$$A(t) = |D_2| + \int_J \int_0^{g(\xi, t)} J(\xi, \beta) d\beta d\xi,$$

where

$$\begin{aligned} J(\xi, \beta) &= \frac{\partial(x_1, x_2)}{\partial(\xi, \beta)} = \begin{vmatrix} \gamma'_1 + \beta N'_1 & N_1 \\ \gamma'_2 + \beta N'_2 & N_2 \end{vmatrix} \\ &= |-\langle \gamma', \mathbf{N}^\perp \rangle - \beta \langle \mathbf{N}', \mathbf{N}^\perp \rangle| = ||\gamma'| + \beta \langle \mathbf{N}', \mathbf{N}^\perp \rangle|. \end{aligned}$$

We show that $|\gamma'| + \beta \langle \mathbf{N}', \mathbf{N}^\perp \rangle \geq 0$. Indeed, from the fact that $\|\gamma\|_{C^2(J)} < \infty$, we have $\|\langle \mathbf{N}', \mathbf{N}^\perp \rangle\|_{L^\infty(J)} < \infty$. Substituting t_0 by a smaller positive number if necessary, we can assume that

$$\|g\|_{L^\infty(J \times (-t_0, t_0))} < \beta_0$$

and

$$\|\langle \mathbf{N}', \mathbf{N}^\perp \rangle\|_{L^\infty(J)} \|g\|_{L^\infty(J \times (-t_0, t_0))} < \theta.$$

Note that the first of these assumptions guarantees that $\partial D_2(t)$ has positive distance from ∂D_1 and ∂D_3 . We have

$$|\beta| |\langle \mathbf{N}', \mathbf{N}^\perp \rangle| \leq \|g\|_{L^\infty(J \times (-t_0, t_0))} \|\langle \mathbf{N}', \mathbf{N}^\perp \rangle\|_{L^\infty(J)} \leq \theta \leq |\gamma'|$$

for all $\xi \in J$ and $|\beta| \leq \|g\|_{L^\infty(J \times (-t_0, t_0))}$. Thus, $J(\xi, \beta) = |\gamma'| + \beta \langle \mathbf{N}', \mathbf{N}^\perp \rangle$. Substituting into the formula for $A(t)$ we have

$$\begin{aligned} A(t) &= A + \int_J \int_0^{g(\xi, t)} (|\gamma'| + \beta \langle \mathbf{N}', \mathbf{N}^\perp \rangle) d\beta d\xi \\ &= A + \int_J \left(g(\xi, t) |\gamma'| + \frac{1}{2} (g(\xi, t))^2 \langle \mathbf{N}', \mathbf{N}^\perp \rangle \right) d\xi. \end{aligned}$$

To obtain $|D_2(t)| = |D_2|$ for all $t \in (-t_0, t_0)$, we find the further constraint on g :

$$\int_J \left(g(\xi, t) |\gamma'| + \frac{1}{2} (g(\xi, t))^2 \langle \mathbf{N}', \mathbf{N}^\perp \rangle \right) d\xi = 0 \quad \forall t \in (-t_0, t_0). \tag{3.4}$$

Moreover, we calculate the derivatives of $A(t)$, that we will use later.

$$\begin{aligned} A'(t) &= \int_J (g_t(\xi, t)|\gamma'(\xi)| + g(\xi, t)g_t(\xi, t)\langle \mathbf{N}', \mathbf{N}^\perp \rangle) d\xi = 0; \\ A''(t) &= \int_J (g_{tt}(\xi, t)|\gamma'(\xi)| + (g(\xi, t)g_{tt}(\xi, t) + g_t^2(\xi, t))\langle \mathbf{N}', \mathbf{N}^\perp \rangle) d\xi = 0. \end{aligned} \quad (3.5)$$

Once we have defined the family $D_2(t)$, we can go back to the functional (1.6). The following lemma describes (1.7) with $\bar{f}_t = b_1\chi_{D_1} + b_2\chi_{D_2(t)} + b_3\chi_{D_3}$. We find an inequality corresponding to [4, (2.3) of Lemma 2.1].

Lemma 3.2. *Let $\bar{f}_t = b_1\chi_{D_1} + b_2\chi_{D_2(t)} + b_3\chi_{D_3}$, where the variation of domain $D_2(t)$ is described by (3.1) and $g : J \times (-t_0, t_0) \rightarrow \mathbb{R}$, $t_0 > 0$, satisfies (3.2), (3.3) and (3.4). Then, for all $v \in H_0^1(\Omega)$, the conditions (1.7) becomes*

$$\begin{aligned} &\int_\Omega (|Dv|^2 - qu_{\bar{f}} v^2 |u_{\bar{f}}|^{q-2} \bar{f}) dx \cdot \int_\gamma g_t^2(\gamma^{-1}, 0) |Du_{\bar{f}}| d\sigma \\ &\geq b_2 c_2^q \left(\int_\gamma g_t(\gamma^{-1}, 0) v d\sigma \right)^2. \end{aligned} \quad (3.6)$$

Proof. We calculate the second derivative of the functional (1.6), with respect to s . We have

$$\frac{\partial E}{\partial s} = (q+1) \int_\Omega (\langle Du_{\bar{f}} + sDv, Dv \rangle - v|u_{\bar{f}} + sv|^q \bar{f}_t) dx$$

and

$$\frac{\partial^2 E}{\partial s^2}(0, 0) = (q+1) \int_\Omega (|Dv|^2 - qu_{\bar{f}} v^2 |u_{\bar{f}}|^{q-2} \bar{f}) dx. \quad (3.7)$$

Before calculating the second derivative of E with respect to t , we rewrite (1.6) in the form

$$\begin{aligned} E(s, t) &= \int_\Omega \frac{q+1}{2} |Du_{\bar{f}} + sDv|^2 dx - b_1 \int_{D_1} (u_{\bar{f}} + sv)|u_{\bar{f}} + sv|^q dx \\ &\quad - b_2 \int_{D_2(t)} (u_{\bar{f}} + sv)|u_{\bar{f}} + sv|^q dx - b_3 \int_{D_3} (u_{\bar{f}} + sv)|u_{\bar{f}} + sv|^q dx - \eta. \end{aligned}$$

We observe that, if $F : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a continuous function, then

$$\int_{D_2(t)} F - \int_{D_2} F = \int_J \int_0^{g(\xi, t)} F(\phi(\xi, g(\xi, \beta))) J(\xi, \beta) d\beta d\xi;$$

whence, from the Fundamental Theorem of Calculus,

$$\frac{\partial}{\partial t} \int_{D_2(t)} F = \int_J g_t(\xi, t) F(\phi(\xi, g(\xi, t))) J(\xi, g(\xi, t)) d\xi.$$

Using the above relation with $F = (u + sv)|u + sv|^q$, we have

$$\frac{\partial E}{\partial t} = -b_2 \int_J g_t(\xi, t) (u_{\bar{f}} + sv)|u_{\bar{f}} + sv|^q J(\xi, g(\xi, t)) d\xi,$$

where, for simplicity of notation, we set $u_{\bar{f}}(\phi(\xi, g(\xi, t))) = u_{\bar{f}}$ and $v(\phi(\xi, g(\xi, t))) = v$. Moreover

$$\begin{aligned} \frac{\partial^2 E}{\partial t^2} &= -b_2 \int_J |u_{\bar{f}} + sv|^q \left\{ [g_{tt}(\xi, t)(u_{\bar{f}} + sv) + (q+1)g_t^2(\xi, t)\langle Du_{\bar{f}} + sDv, \mathbf{N} \rangle] \right. \\ &\quad \left. \times J(\xi, g(\xi, t)) + g_t^2(\xi, t)(u_{\bar{f}} + sv)\langle \mathbf{N}', \mathbf{N}^\perp \rangle \right\} d\xi, \end{aligned}$$

where we have used that

$$\begin{aligned} \frac{\partial}{\partial t} u_{\bar{f}}(\phi(\xi, g(\xi, t))) &= \langle Du_{\bar{f}}(\phi(\xi, g(\xi, t))), \mathbf{N} \rangle g_t(\xi, t), \\ \frac{\partial}{\partial t} J(\xi, g(\xi, t)) &= g_t \langle \mathbf{N}', \mathbf{N}^\perp \rangle. \end{aligned}$$

We note that, when $t = 0$,

$$u_{\bar{f}}(\phi(\xi, g(\xi, t))) = u_{\bar{f}}(\gamma(\xi)) = c_2,$$

$Du_{\bar{f}}(\phi(\xi, g(\xi, 0))) = -|Du_{\bar{f}}(\gamma(\xi))|\mathbf{N}(\xi)$ and $J(\xi, g(\xi, 0)) = J(\xi, 0) = |\gamma'(\xi)|$. Evaluating the above expression in $(0, 0)$, we find

$$\begin{aligned} \frac{\partial^2 E}{\partial t^2}(0, 0) &= -b_2 c_2^{q+1} \int_J [g_{tt}(\xi, 0)|\gamma'(\xi)| + g_t^2(\xi, 0)\langle \mathbf{N}', \mathbf{N}^\perp \rangle] d\xi \\ &\quad + b_2 c_2^q (q + 1) \int_J g_t^2(\xi, 0) |Du_{\bar{f}}(\gamma(\xi))| |\gamma'(\xi)| d\xi. \end{aligned}$$

By using (3.5) with $t = 0$ we find

$$\begin{aligned} \frac{\partial^2 E}{\partial t^2}(0, 0) &= b_2 c_2^q (q + 1) \int_J g_t^2(\xi, 0) |Du_{\bar{f}}(\gamma(\xi))| |\gamma'(\xi)| d\xi \\ &= b_2 c_2^q (q + 1) \int_\gamma g_t^2(\gamma^{-1}, 0) |Du_{\bar{f}}| d\sigma. \end{aligned} \tag{3.8}$$

We also have

$$\frac{\partial^2 E}{\partial s \partial t} = -b_2 (q + 1) \int_J g_t(\xi, t) v |u_{\bar{f}} + sv|^q J(\xi, g(\xi, t)) d\xi;$$

that is,

$$\begin{aligned} \frac{\partial^2 E}{\partial s \partial t}(0, 0) &= -b_2 c_2^q (q + 1) \int_J g_t(\xi, 0) v(\gamma(\xi)) |\gamma'(\xi)| d\xi \\ &= -b_2 c_2^q (q + 1) \int_\gamma g_t(\gamma^{-1}, 0) v d\sigma. \end{aligned} \tag{3.9}$$

Using (1.7) in the form

$$\frac{\partial^2 E}{\partial s^2}(0, 0) \frac{\partial^2 E}{\partial t^2}(0, 0) \geq \left(\frac{\partial^2 E}{\partial s \partial t}(0, 0) \right)^2,$$

and using (3.7), (3.8) and (3.9) in this inequality, we obtain (3.6). □

Note that in inequality (3.6) only $g(\gamma^{-1}, 0)$ appears. Moreover, $g(\gamma^{-1}, 0)$ has null integral on γ . Indeed, differentiating (3.4) with respect to t and putting $t = 0$, we obtain

$$\int_J g(\xi, 0) |\gamma'| d\xi = 0.$$

Now a natural question arises: does inequality (3.6) hold for any function h with null integral on γ ? The answer is contained in the following result.

Lemma 3.3. *Let J and γ be the same as described. Let $h : \gamma \rightarrow \mathbb{R}$ bounded, continuous and such that $\int_\gamma h d\sigma = 0$. Then, for all $v \in H_0^1(\Omega)$ and for all $a \in \mathbb{R}$ we have*

$$\int_\Omega (|Dv|^2 - qu_{\bar{f}} v^2 |u_{\bar{f}}|^{q-2} \bar{f}) dx \cdot \int_\gamma h^2 |Du_{\bar{f}}| d\sigma \geq b_2 c_2^q \left(\int_\gamma h(v - a) d\sigma \right)^2. \tag{3.10}$$

The proof of the above lemma is similar to that of [4, Lemma 2.2]; we omit it. The following lemma is an analogue to [4, Lemma 3.1].

Lemma 3.4. *Let P be a point on $\mathcal{F} = \partial\{u_{\bar{f}} > c_2\}$. Suppose that for all $k \in \mathbb{Z}^+$ there exist a positive number r_k , a bounded open interval J_k and a regular curve $\gamma_k : J_k \rightarrow \mathcal{F}^*$ such that $r_1 > r_2 > \dots \rightarrow 0$, $\overline{\gamma_k(J_k)} \subset \mathcal{F}^* \cap B_{r_k}(P) \setminus \overline{B_{r_{k+1}}(P)}$. Then we must have*

$$\sum_{k=1}^{\infty} \int_{\gamma(J_k)} \frac{1}{|Du_{\bar{f}}|} d\sigma < \infty.$$

Proof. Without loss of generality, we assume that P is the origin. We suppose also that $J_k \cap J_h = \emptyset$ for all $k \neq h$, and denote all γ_k with γ . We define

$$J_{k,m} = \begin{cases} J_k \cup J_{k+1} \cup \dots \cup J_m & \text{if } m \geq k, \\ \emptyset & \text{otherwise.} \end{cases}$$

We suppose by contradiction that

$$\sum_{k=1}^{\infty} \int_{\gamma(J_k)} \frac{1}{|Du_{\bar{f}}|} d\sigma = \infty. \quad (3.11)$$

Let V be a smooth radial function in \mathbb{R}^2 , decreasing in $|x|$, defined by

$$\begin{cases} V(x) = 2, & |x| = 0 \\ 1 < V(x) < 2, & 0 < |x| < 1/2 \\ 0 < V(x) < 1, & 1/2 < |x| < 1 \\ V(x) = 0, & |x| \geq 1. \end{cases}$$

For all $k \in \mathbb{Z}^+$ we define $v_k(x) = V(\frac{x}{r_k})$. Consider k large enough such that $\text{supp } v_k \subset \Omega$. Now we fix k ; we have

$$\begin{cases} v_k(x) - 1 = 1, & |x| = 0 \\ 0 < v_k(x) - 1 < 1, & 0 < |x| < r_k/2 \\ -1 < v_k(x) - 1 < 0, & r_k/2 < |x| < r_k \\ v_k(x) - 1 = -1, & |x| \geq r_k. \end{cases}$$

Since J_k and $|\gamma'|$ are bounded, $\gamma(J_k)$ is of finite length. Moreover, $|Du_{\bar{f}}|$ is uniformly bounded away from 0 on $\gamma(J_k)$ since $\overline{\gamma(J_k)} \subset \mathcal{F}^*$. Together with the fact that $\gamma(J_{1,k-1}) \subset (B_{r_k})^C$, we have

$$-\infty < \int_{\gamma(J_{1,k-1})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma = - \int_{\gamma(J_{1,k-1})} \frac{1}{|Du_{\bar{f}}|} d\sigma < 0.$$

Choose m such that $r_m < r_k/2$. From the facts that $v_k(x) - 1 > 0$ in B_{r_m} , $\gamma(J_l) \subset B_{r_m}$ for all $l \geq m$ and $v_k(x) - 1 \rightarrow 1$ as $x \rightarrow 0$ and (3.11), we have

$$\int_{\gamma(J_{m,l})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma \rightarrow \infty \quad \text{for } l \rightarrow \infty.$$

Consequently, there must be a number $l \geq m$ such that

$$\int_{\gamma(J_{m,l-1})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma \leq - \int_{\gamma(J_{1,k-1})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma < \int_{\gamma(J_{m,l})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma.$$

Choose a subinterval $J'_l \subset J_l$ such that

$$\int_{\gamma(J_{m,l-1})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma + \int_{\gamma(J'_l)} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma = - \int_{\gamma(J_{1,k-1})} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma.$$

Then we have

$$\int_{\gamma(J^k)} \frac{v_k - 1}{|Du_{\bar{f}}|} d\sigma = 0,$$

where $J^k = J_{1,k-1} \cup J_{m,l-1} \cup J'_l$.

Now we can apply Lemma 3.3 to J^k , γ , v_k , $a = 1$ and $h = \frac{v_k - 1}{|Du_{\bar{f}}|}$ and, after rearranging, obtain

$$\int_{\Omega} \left(|Dv_k|^2 - qu_{\bar{f}} v_k^2 |u_{\bar{f}}|^{q-2} \bar{f} \right) dx \geq b_2 c_2^q \int_{\gamma(J^k)} \frac{(v_k - 1)^2}{|Du_{\bar{f}}|} d\sigma.$$

We find that

$$\int_{\Omega} \left(|Dv_k|^2 - qu_{\bar{f}} v_k^2 |u_{\bar{f}}|^{q-2} \bar{f} \right) dx \leq \int_{B_1(0)} |DV|^2 dx.$$

By the above estimate, for a suitable constant C , we have

$$\begin{aligned} C \int_{B_1(0)} |DV|^2 dx &\geq \int_{\gamma(J^k)} \frac{(v_k - 1)^2}{|Du_{\bar{f}}|} d\sigma \\ &\geq \int_{\gamma(J_{1,k-1})} \frac{(v_k - 1)^2}{|Du_{\bar{f}}|} d\sigma \\ &= \sum_{h=1}^{k-1} \int_{\gamma(J_h)} \frac{(v_k - 1)^2}{|Du_{\bar{f}}|} d\sigma. \end{aligned}$$

Then, when $k \rightarrow \infty$, we have

$$C \int_{B_1(0)} |DV|^2 dx \geq \sum_{h=1}^{\infty} \int_{\gamma(J_h)} \frac{(v_k - 1)^2}{|Du_{\bar{f}}|} d\sigma = +\infty,$$

which is a contradiction. So we must have

$$\sum_{k=1}^{\infty} \int_{\gamma(J_k)} \frac{d\sigma}{|Du_{\bar{f}}|} < \infty,$$

as desired. □

Lemma 3.5. *Let P be a point on $\mathcal{F} = \partial\{u_{\bar{f}} > c_2\}$. Suppose that there are numbers $K \in \mathbb{Z}$ and $\bar{\sigma} > 0$ such that, for each $k \geq K$, there exists a regular curve $\gamma_k : J_k \rightarrow \mathcal{F}^*$ with the following two properties:*

$$\begin{aligned} \overline{\gamma(J_k)} &\subset \mathcal{F}^* \cap B_{2^{-k}}(P) \setminus \overline{B_{2^{-(k+1)}}(P)}, \\ \mathcal{H}^1(\gamma_k(J_k)) &= \int_{J_k} |\gamma'(\xi)| d\xi > \bar{\sigma} 2^{-k}. \end{aligned}$$

Then $|Du_{\bar{f}}(P)| > 0$.

For a proof of the above lemma, see [4, Lemma 3.2]. From an intuitive point of view, this lemma says that, if the set $\partial\{u_{\bar{f}} > c_2\} \cap \{|Du_{\bar{f}}| > 0\}$ is big enough around a point of $\partial\{u_{\bar{f}} > c_2\}$, then $|Du_{\bar{f}}| > 0$ at this point.

Now, we are able to prove our main theorem.

Proof of Theorem 3.1. By using the previous Lemmas and superharmonicity of $u_{\bar{f}}$ the Theorem follows from the results of sections 5 and 6 in [4]. \square

Open problems. The method used in this paper to prove regularity does not work when the number of level sets of \bar{f} is infinite. Therefore it remains to study the boundaries of level sets of \bar{f} in the case of the rearrangement class $\mathfrak{F}(f_0)$ of a general function f_0 .

We can obtain an analogous result to Lemma 3.4 for the p -Laplacian operator, but we cannot go further because we lack a suitable regularity theory for the p -Laplacian operator and its solutions. We think that it is reasonable to guess that a regularity result of the type that we have proven in this work will hold for the situation with the p -Laplacian when $p < 2$.

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